PART - A

UNIT – 1 7 Hours

Transistors, UJTs, and Thyristors:
Operating Point, Common-Emitter Configuration, Thermal Runaway, Transistor Switch, Unijunction Transistors, SCR.

UNIT – 2 6 Hours

Field Effect Transistors:
Bipolar Junction Transistors versus Field Effect Transistors, Junction Field Effect Transistors, Metal Oxide Field Effect Transistors, Differences between JFETs and MOSFETs, Handling MOSFETs, Biasing MOSFETs, FET Applications, CMOS Devices, Insulated Gate Bipolar Transistors (IGBTs)

UNIT – 3 6 Hours

Optoelectronic Devices:
Introduction, Photosensors, Photoconductors, Photodiodes, Phototransistors, Light-Emitting Diodes, Liquid Crystal Displays, Cathode Ray Tube Displays, Emerging Display Technologies, Optocouplers

UNIT – 4 7 Hours

Small Signal Analysis of Amplifiers:
Amplifier Bandwidth: General Frequency Considerations, Hybrid h-Parameter Model for an Amplifier, Transistor Hybrid Model, Analysis of a Transistor Amplifier using complete h-Parameter Model, Analysis of a Transistor Amplifier Configurations using Simplified h-Parameter Model (CE configuration only), Small-Signal Analysis of FET Amplifiers, Cascading Amplifiers, Darlington Amplifier, Low-Frequency Response of Amplifiers (BJT amplifiers only).

PART – B

UNIT - 5 6 Hours

Large Signal Amplifiers, Feedback Amplifier:
UNIT - 6
**Sinusoidal Oscillators, Wave-Shaping Circuits:**

UNIT - 7
**Linear Power Supplies, Switched mode Power Supplies:**

UNIT - 8
**Operational Amplifiers:** Ideal Opamp versus Practical Opamp, Performance Parameters, Some Applications: Peak Detector Circuit, Absolute Value Circuit, Comparator, Active Filters, Phase Shifters, Instrumentation Amplifier, Non-Linear Amplifier, Relaxation Oscillator, Current-To-Voltage Converter, Voltage-To-Current Converter, Sine Wave Oscillators.

Text Book:
   (4.1, 4.2, 4.7, 4.8, 5.1 to 5.3, 5.5, 5.6, 5.8, 5.9, 5.13, 5.14, 6.1, 6.3, 7.1 to 7.5, 7.10 to 7.14, Listed topics only from 8, 10.1, 11, 12.1, 12.2, 12.3, 12.5, 13.1 to 13.6, 13.9, 13.10, 14.1, 14.2, 14.6, 14.7, 15.1, 15.5 to 15.7, 16.3, 16.4, 17.12 to 17.22)

Reference Books:
2. R. D. Sudhaker Samuel: Electronic Circuits, Sanguine-Pearson
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Unit – 1: Transistor, UJT’s, and Thyristors

In the Diode tutorials we saw that simple diodes are made up from two pieces of semiconductor material, either silicon or germanium to form a simple PN-junction and we also learnt about their properties and characteristics. If we now join together two individual signal diodes back-to-back, this will give us two PN-junctions connected together in series that share a common P or N terminal. The fusion of these two diodes produces a three layer, two junctions, and three terminal devices forming the basis of a Bipolar Junction Transistor, or BJT for short.

1.1 Operating Point

Operating Regions

The pink shaded area at the bottom of the curves represents the "Cut-off" region while the blue area to the left represents the "Saturation" region of the transistor. Both these transistor regions are defined as:

1. Cut-off Region

Here the operating conditions of the transistor are zero input base current ($I_B$), zero output collector current ($I_C$) and maximum collector voltage ($V_{CE}$) which results in a large depletion layer and no current flowing through the device. Therefore the transistor is switched "Fully-Off".
Cut-off Characteristics

The input and Base are grounded (0v)
- Base-Emitter voltage $V_{BE} < 0.7V$
- Base-Emitter junction is reverse biased
- Base-Collector junction is reverse biased
- Transistor is "fully-OFF" (Cut-off region)
  - No Collector current flows ($I_C = 0$)
  - $V_{OUT} = V_{CE} = V_{CC} = "1"$
  - Transistor operates as an "open switch"

Then we can define the "cut-off region" or "OFF mode" when using a bipolar transistor as a switch as being, both junctions reverse biased, $I_B < 0.7V$ and $I_C = 0$. For a PNP transistor, the Emitter potential must be negative with respect to the Base.

2. Saturation Region

Here the transistor will be biased so that the maximum amount of base current is applied, resulting in maximum collector current resulting in the minimum collector emitter voltage drop which results in the depletion layer being as small as possible and maximum current flowing through the transistor. Therefore the transistor is switched "Fully-ON".

Saturation Characteristics

The input and Base are connected to $V_{CC}$
- Base-Emitter voltage $V_{BE} > 0.7V$
- Base-Emitter junction is forward biased
- Base-Collector junction is forward biased
- Transistor is "fully-ON" (saturation region)
  - Max Collector current flows ($I_C = V_{CC}/R_L$)
  - $V_{CE} = 0$ (ideal saturation)
  - $V_{OUT} = V_{CE} = "0"$
  - Transistor operates as a "closed switch"
Then we can define the "saturation region" or "ON mode" when using a bipolar transistor as a switch as being, both junctions forward biased, $I_B > 0.7V$ and $I_C = \text{Maximum}$. For a PNP transistor, the Emitter potential must be positive with respect to the Base.

Then the transistor operates as a "single-pole single-throw" (SPST) solid state switch. With a zero signal applied to the Base of the transistor it turns "OFF" acting like an open switch and zero collector current flows. With a positive signal applied to the Base of the transistor it turns "ON" acting like a closed switch and maximum circuit current flows through the device.

An example of an NPN Transistor as a switch being used to operate a relay is given below. With inductive loads such as relays or solenoids a flywheel diode is placed across the load to dissipate the back EMF generated by the inductive load when the transistor switches "OFF" and so protect the transistor from damage. If the load is of a very high current or voltage nature, such as motors, heaters etc, then the load current can be controlled via a suitable relay as shown.

Transistor

Transistors are three terminal active devices made from different semiconductor materials that can act as either an insulator or a conductor by the application of a small signal voltage. The transistor's ability to change between these two states enables it to have two basic functions: "switching" (digital electronics) or "amplification" (analogue electronics). Then bipolar transistors have the ability to operate within three different regions:

1. Active Region - the transistor operates as an amplifier and $I_c = \beta . I_B$
2. Saturation - the transistor is "fully-ON" operating as a switch and $I_c = I(\text{saturation})$
3. Cut-off - the transistor is "fully-OFF" operating as a switch and $I_c = 0$

Typical Bipolar Transistor
The word Transistor is an acronym, and is a combination of the words Transfer Varistor used to describe their mode of operation way back in their early days of development. There are two basic types of bipolar transistor construction, PNP and NPN, which basically describes the physical arrangement of the P-type and N-type semiconductor materials from which they are made.

The **Bipolar Transistor** basic construction consists of two PN-junctions producing three connecting terminals with each terminal being given a name to identify it from the other two. These three terminals are known and labeled as the Emitter (E), the Base (B) and the Collector (C) respectively.

Bipolar Transistors are current regulating devices that control the amount of current flowing through them in proportion to the amount of biasing voltage applied to their base terminal acting like a current-controlled switch. The principle of operation of the two transistor types PNP and NPN, is exactly the same the only difference being in their biasing and the polarity of the power supply for each type.

**Bipolar Transistor Construction**

![Bipolar Transistor Construction Diagram](image-url)

The construction and circuit symbols for both the PNP and NPN bipolar transistor are given above with the arrow in the circuit symbol always showing the direction of "conventional current..."
flow" between the base terminal and its emitter terminal. The direction of the arrow always points from the positive P-type region to the negative N-type region for both transistor types, exactly the same as for the standard diode symbol.

Bipolar Transistor Configurations

As the Bipolar Transistor is a three terminal device, there are basically three possible ways to connect it within an electronic circuit with one terminal being common to both the input and output. Each method of connection responding differently to its input signal within a circuit as the static characteristics of the transistor vary with each circuit arrangement.

1. Common Base Configuration   - has Voltage Gain but no Current Gain.
2. Common Emitter Configuration - has both Current and Voltage Gain.
3. Common Collector Configuration - has Current Gain but no Voltage Gain.

The Common Base (CB) Configuration

As its name suggests, in the Common Base or grounded base configuration, the BASE connection is common to both the input signal AND the output signal with the input signal being applied between the base and the emitter terminals. The corresponding output signal is taken from between the base and the collector terminals as shown with the base terminal grounded or connected to a fixed reference voltage point. The input current flowing into the emitter is quite large as its the sum of both the base current and collector current respectively therefore, the collector current output is less than the emitter current input resulting in a current gain for this type of circuit of "1" (unity) or less, in other words the common base configuration "attenuates" the input signal.

This type of amplifier configuration is a non-inverting voltage amplifier circuit, in that the signal voltages Vin and Vout are in-phase. This type of transistor arrangement is not very common due to its unusually high voltage gain characteristics. Its output characteristics represent that of a forward biased diode while the input characteristics represent that of an illuminated photo-diode. Also this type of bipolar transistor configuration has a high ratio of output to input resistance or
more importantly "load" resistance (RL) to "input" resistance (Rin) giving it a value of "Resistance Gain". Then the voltage gain (Av) for a common base configuration is therefore given as:

**Common Base Voltage Gain**

\[
A_v = \frac{V_{out}}{V_{in}} = \frac{I_C \times R_L}{I_E \times R_{IN}}
\]

Where: Ic/Ie is the current gain, alpha (\(\alpha\)) and RL/Rin is the resistance gain.

The common base circuit is generally only used in single stage amplifier circuits such as microphone pre-amplifier or radio frequency (Rf) amplifiers due to its very good high frequency response.

### 1.2 The Common Emitter (CE) Configuration

In the **Common Emitter** or grounded emitter configuration, the input signal is applied between the base, while the output is taken from between the collector and the emitter as shown. This type of configuration is the most commonly used circuit for transistor based amplifiers and which represents the "normal" method of bipolar transistor connection. The common emitter amplifier configuration produces the highest current and power gain of all the three bipolar transistor configurations. This is mainly because the input impedance is LOW as it is connected to a forward-biased PN-junction, while the output impedance is HIGH as it is taken from a reverse-biased PN-junction.

**The Common Emitter Amplifier Circuit**

![Common Emitter Amplifier Circuit Diagram](www.vtucs.com)

In this type of configuration, the current flowing out of the transistor must be equal to the currents flowing into the transistor as the emitter current is given as \(I_E = I_C + I_B\). Also, as the load resistance (RL) is connected in series with the collector, the current gain of the common emitter transistor configuration is quite large as it is the ratio of \(I_C/I_B\) and is given the Greek symbol of Beta, (\(\beta\)). As the emitter current for a common emitter configuration is defined as \(I_E = I_C + I_B\),
the ratio of Ic/Ie is called Alpha, given the Greek symbol of \( \alpha \). Note: that the value of Alpha will always be less than unity.

Since the electrical relationship between these three currents, Ib, Ic and Ie is determined by the physical construction of the transistor itself, any small change in the base current (Ib), will result in a much larger change in the collector current (Ic). Then, small changes in current flowing in the base will thus control the current in the emitter-collector circuit. Typically, Beta has a value between 20 and 200 for most general purpose transistors.

By combining the expressions for both Alpha, \( \alpha \) and Beta, \( \beta \) the mathematical relationship between these parameters and therefore the current gain of the transistor can be given as:

\[
\text{Alpha} (\alpha) = \frac{I_C}{I_E} \quad \text{and} \quad \text{Beta} (\beta) = \frac{I_C}{I_B}
\]

\[\therefore I_C = \alpha I_E = \beta I_B\]

\[\text{as: } \alpha = \frac{\beta}{\beta + 1} \quad \beta = \frac{\alpha}{1 - \alpha}\]

\[I_E = I_C + I_B\]

Where: "Ic" is the current flowing into the collector terminal, "Ib" is the current flowing into the base terminal and "Ie" is the current flowing out of the emitter terminal.

Then to summarize, this type of bipolar transistor configuration has a greater input impedance, current and power gain than that of the common base configuration but its voltage gain is much lower. The common emitter configuration is an inverting amplifier circuit resulting in the output signal being 180° out-of-phase with the input voltage signal.

The Common Collector (CC) Configuration

In the Common Collector or grounded collector configuration, the collector is now common through the supply. The input signal is connected directly to the base, while the output is taken from the emitter load as shown. This type of configuration is commonly known as a Voltage Follower or Emitter Follower circuit. The emitter follower configuration is very useful for impedance matching applications because of the very high input impedance, in the region of hundreds of thousands of Ohms while having a relatively low output impedance.
The common emitter configuration has a current gain approximately equal to the $\beta$ value of the transistor itself. In the common collector configuration the load resistance is situated in series with the emitter so its current is equal to that of the emitter current. As the emitter current is the combination of the collector AND the base current combined, the load resistance in this type of transistor configuration also has both the collector current and the input current of the base flowing through it. Then the current gain of the circuit is given as:

**The Common Collector Current Gain**

$$I_E = I_C + I_B$$

$$A_C = \frac{I_E}{I_B} = \frac{I_C + I_B}{I_B}$$

$$A_i = \frac{I_C}{I_B} + 1$$

$$A_i = \beta + 1$$

This type of bipolar transistor configuration is a non-inverting circuit in that the signal voltages of Vin and Vout are in-phase. It has a voltage gain that is always less than "1" (unity). The load resistance of the common collector transistor receives both the base and collector currents giving a large current gain (as with the common emitter configuration) therefore, providing good current amplification with very little voltage gain.

**Bipolar Transistor Summary**
Then to summarize, the behavior of the bipolar transistor in each one of the above circuit configurations is very different and produces different circuit characteristics with regards to input impedance, output impedance and gain whether this is voltage gain, current gain or power gain and this is summarized in the table below.

**Bipolar Transistor Characteristics**

The static characteristics for a **Bipolar Transistor** can be divided into the following three main groups.

<table>
<thead>
<tr>
<th>Input Characteristics:</th>
<th>Common Base</th>
<th>$\Delta V_{EB} / \Delta I_E$</th>
<th>Common Emitter</th>
<th>$\Delta V_{BE} / \Delta I_B$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Output Characteristics:</td>
<td>Common Base</td>
<td>$\Delta V_C / \Delta I_C$</td>
<td>Common Emitter</td>
<td>$\Delta V_C / \Delta I_C$</td>
</tr>
<tr>
<td>Transfer Characteristics:</td>
<td>Common Base</td>
<td>$\Delta I_C / \Delta I_E$</td>
<td>Common Emitter</td>
<td>$\Delta I_C / \Delta I_B$</td>
</tr>
</tbody>
</table>

With the characteristics of the different transistor configurations given in the following table:

<table>
<thead>
<tr>
<th>Characteristic</th>
<th>Common Base</th>
<th>Common Emitter</th>
<th>Common Collector</th>
</tr>
</thead>
<tbody>
<tr>
<td>Input Impedance</td>
<td>Low</td>
<td>Medium</td>
<td>High</td>
</tr>
<tr>
<td>Output Impedance</td>
<td>Very High</td>
<td>High</td>
<td>Low</td>
</tr>
<tr>
<td>Phase Angle</td>
<td>0°</td>
<td>180°</td>
<td>0°</td>
</tr>
<tr>
<td>Voltage Gain</td>
<td>High</td>
<td>Medium</td>
<td>Low</td>
</tr>
<tr>
<td>Current Gain</td>
<td>Low</td>
<td>Medium</td>
<td>High</td>
</tr>
<tr>
<td>Power Gain</td>
<td>Low</td>
<td>Very High</td>
<td>Medium</td>
</tr>
</tbody>
</table>

In the next tutorial about **Bipolar Transistors**, we will look at the NPN Transistor in more detail when used in the common emitter configuration as an amplifier as this is the most widely used configuration due to its flexibility and high gain. We will also plot the output characteristics curves commonly associated with amplifier circuits as a function of the collector current to the base current.
The NPN Transistor

In the previous tutorial we saw that the standard **Bipolar Transistor** or BJT, comes in two basic forms. An NPN (Negative-Positive-Negative) type and a PNP (Positive-Negative-Positive) type, with the most commonly used transistor type being the **NPN Transistor**. We also learnt that the transistor junctions can be biased in one of three different ways - **Common Base**, **Common Emitter** and **Common Collector**. In this tutorial we will look more closely at the "Common Emitter" configuration using **NPN Transistors** with an example of the construction of a NPN transistor along with the transistors current flow characteristics is given below.

### An NPN Transistor Configuration

![NPN Transistor Configuration Diagram](image)

(Note: Arrow defines the emitter and conventional current flow, "out" for an NPN transistor.)

The construction and terminal voltages for an NPN transistor are shown above. The voltage between the Base and Emitter (V$_{BE}$), is positive at the Base and negative at the Emitter because for an NPN transistor, the Base terminal is always positive with respect to the Emitter. Also the Collector supply voltage is positive with respect to the Emitter (V$_{CE}$). So for an NPN transistor to conduct the Collector is always more positive with respect to both the Base and the Emitter.

![NPN Transistor Connections Diagram](image)

**NPN Transistor Connections**

Then the voltage sources are connected to an NPN transistor as shown. The Collector is connected to the supply voltage V$_{CC}$ via the load resistor, RL which also acts to limit the maximum current flowing through the device. The Base supply voltage V$_{B}$ is connected to the Base resistor R$_{B}$, which again is used to limit the maximum Base current.
We know that the transistor is a "current" operated device (Beta model) and that a large current (Ic) flows freely through the device between the collector and the emitter terminals when the transistor is switched "fully-ON". However, this only happens when a small biasing current (Ib) is flowing into the base terminal of the transistor at the same time thus allowing the Base to act as a sort of current control input.

The transistor current in an NPN transistor is the ratio of these two currents (Ic/Ib), called the DC Current Gain of the device and is given the symbol of hfe or nowadays Beta (β). The value of β can be large up to 200 for standard transistors, and it is this large ratio between Ic and Ib that makes the NPN transistor a useful amplifying device when used in its active region as Ib provides the input and Ic provides the output. Note that Beta has no units as it is a ratio.

Also, the current gain of the transistor from the Collector terminal to the Emitter terminal, Ic/Ie, is called Alpha (α), and is a function of the transistor itself (electrons diffusing across the junction). As the emitter current Ie is the product of a very small base current plus a very large collector current, the value of alpha α, is very close to unity, and for a typical low-power signal transistor this value ranges from about 0.950 to 0.999

\[ \alpha \text{ and } \beta \text{ Relationship in a NPN Transistor} \]

\[
\text{DC Current Gain} = \frac{\text{Output Current}}{\text{Input Current}} = \frac{I_c}{I_b}
\]

\[
I_E = I_B + I_C \quad \text{(KCL)} \quad \text{and} \quad \frac{I_C}{I_E} = \alpha
\]

Thus: \[ I_E = I_B + I_C = \frac{I_C}{\alpha} \]

and \[ I_B = I_C \left( 1 - \frac{1}{\alpha} \right) \]

\[ \therefore \quad \beta = \frac{I_C}{I_B} = \frac{1}{1 - \frac{1}{\alpha}} = \frac{\alpha}{1 - \alpha} \]

By combining the two parameters α and β we can produce two mathematical expressions that gives the relationship between the different currents flowing in the transistor.

\[ \beta = \frac{\alpha}{1 - \alpha} \quad \text{and} \quad \alpha = \frac{\beta}{\beta + 1} \]

If \[ \alpha = 0.99 \quad \beta = \frac{0.99}{0.01} = 99 \]
The values of Beta vary from about 20 for high current power transistors to well over 1000 for high frequency low power type bipolar transistors. The value of Beta for most standard NPN transistors can be found in the manufactures datasheets but generally range between 50 - 200.

The equation above for Beta can also be re-arranged to make Ic as the subject, and with a zero base current ( $I_b = 0$ ) the resultant collector current Ic will also be zero, ( $\beta \times 0$ ). Also when the base current is high the corresponding collector current will also be high resulting in the base current controlling the collector current. One of the most important properties of the Bipolar Junction Transistor is that a small base current can control a much larger collector current. Consider the following example.

**Example No1**

An NPN Transistor has a DC current gain, (Beta) value of 200. Calculate the base current Ib required to switch a resistive load of 4mA.

\[ I_B = \frac{I_C}{\beta} = \frac{4 \times 10^{-3}}{200} = 20\mu A \]

Therefore, $\beta = 200$, $I_c = 4mA$ and $I_b = 20\mu A$.

One other point to remember about NPN Transistors. The collector voltage, ( $V_c$ ) must be greater and positive with respect to the emitter voltage, ( $V_e$ ) to allow current to flow through the transistor between the collector-emitter junctions. Also, there is a voltage drop between the Base and the Emitter terminal of about 0.7v (one diode volt drop) for silicon devices as the input characteristics of an NPN Transistor are of a forward biased diode. Then the base voltage, ( $V_{be}$ ) of a NPN transistor must be greater than this 0.7V otherwise the transistor will not conduct with the base current given as.

\[ I_B = \frac{V_B - V_{BE}}{R_B} \]

Where: $I_b$ is the base current, $V_b$ is the base bias voltage, $V_{be}$ is the base-emitter volt drop (0.7v) and $R_b$ is the base input resistor. Increasing $I_b$, $V_{be}$ slowly increases to 0.7V but $I_c$ rises exponentially.

**Example No2**

An NPN Transistor has a DC base bias voltage, $V_b$ of 10v and an input base resistor, $R_b$ of 100kΩ. What will be the value of the base current into the transistor.

\[ I_B = \frac{V_B - V_{BE}}{R_B} = \frac{10 - 0.7}{100k\Omega} = 93\mu A \]
Therefore, \( I_b = 93\mu A \).

**The Common Emitter Configuration**

As well as being used as a semiconductor switch to turn load currents "ON" or "OFF" by controlling the Base signal to the transistor in either its saturation or cut-off regions, **NPN Transistors** can also be used in its active region to produce a circuit which will amplify any small AC signal applied to its Base terminal with the Emitter grounded. If a suitable DC "biasing" voltage is firstly applied to the transistors Base terminal thus allowing it to always operate within its linear active region, an inverting amplifier circuit called a single stage common emitter amplifier is produced.

One such Common Emitter Amplifier configuration of an NPN transistor is called a **Class A Amplifier**. A "Class A Amplifier" operation is one where the transistors Base terminal is biased in such a way as to forward bias the Base-emitter junction. The result is that the transistor is always operating halfway between its cut-off and saturation regions, thereby allowing the transistor amplifier to accurately reproduce the positive and negative halves of any AC input signal superimposed upon this DC biasing voltage. Without this "Bias Voltage" only one half of the input waveform would be amplified. This common emitter amplifier configuration using an NPN transistor has many applications but is commonly used in audio circuits such as pre-amplifier and power amplifier stages.

With reference to the common emitter configuration shown below, a family of curves known as the **Output Characteristics Curves**, relates the output collector current, \( I_c \) to the collector voltage, \( V_{ce} \) when different values of Base current, \( I_b \) are applied to the transistor for transistors with the same \( \beta \) value. A DC "Load Line" can also be drawn onto the output characteristics curves to show all the possible operating points when different values of base current are applied. It is necessary to set the initial value of \( V_{ce} \) correctly to allow the output voltage to vary both up and down when amplifying AC input signals and this is called setting the operating point or Quiescent Point, \( Q \)-point for short and this is shown below.

**Single Stage Common Emitter Amplifier Circuit**

![Output Characteristics Curves of a Typical Bipolar Transistor](image-url)
The most important factor to notice is the effect of $V_{ce}$ upon the collector current $I_c$ when $V_{ce}$ is greater than about 1.0 volts. We can see that $I_c$ is largely unaffected by changes in $V_{ce}$ above this value and instead it is almost entirely controlled by the base current, $I_b$. When this happens we can say then that the output circuit represents that of a "Constant Current Source". It can also be seen from the common emitter circuit above that the emitter current $I_e$ is the sum of the collector current, $I_c$ and the base current, $I_b$, added together so we can also say that $I_e = I_c + I_b$ for the common emitter (CE) configuration.

By using the output characteristics curves in our example above and also Ohm’s Law, the current flowing through the load resistor, $(RL)$, is equal to the collector current, $I_c$ entering the transistor which in turn corresponds to the supply voltage, $(V_{cc})$ minus the voltage drop between the collector and the emitter terminals, $(V_{ce})$ and is given as:

$$I_c = \frac{V_{cc} - V_{ce}}{R_L}$$

Also, a straight line representing the **Dynamic Load Line** of the transistor can be drawn directly onto the graph of curves above from the point of "Saturation" (A) when $V_{ce} = 0$ to the point of "Cut-off" (B) when $I_c = 0$ thus giving us the "Operating" or **Q-point** of the transistor. These two points are joined together by a straight line and any position along this straight line represents the "Active Region" of the transistor. The actual position of the load line on the characteristics curves can be calculated as follows:
Then, the collector or output characteristics curves for **Common Emitter NPN Transistors** can be used to predict the Collector current, $I_C$, when given $V_{ce}$ and the Base current, $I_b$. A Load Line can also be constructed onto the curves to determine a suitable Operating or **Q-point** which can be set by adjustment of the base current. The slope of this load line is equal to the reciprocal of the load resistance which is given as: $-1/R_L$

Then we can define a **NPN Transistor** as being normally "OFF" but a small input current and a small positive voltage at its Base (B) relative to its Emitter (E) will turn it "ON" allowing a much large Collector-Emitter current to flow. NPN transistors conduct when $V_c$ is much greater than $V_e$.

In the next tutorial about **Bipolar Transistors**, we will look at the opposite or complementary form of the NPN Transistor called the **PNP Transistor** and show that the PNP Transistor has very similar characteristics to their NPN transistor except that the polarities (or biasing) of the current and voltage directions are reversed.

**The PNP Transistor**

The **PNP Transistor** is the exact opposite to the **NPN Transistor** device we looked at in the previous tutorial. Basically, in this type of transistor construction the two diodes are reversed with respect to the NPN type giving a Positive-Negative-Positive configuration, with the arrow which also defines the Emitter terminal this time pointing inwards in the transistor symbol.

Also, all the polarities for a PNP transistor are reversed which means that it "sinks" current as opposed to the NPN transistor which "sources" current. The main difference between the two types of transistors is that holes are the more important carriers for PNP transistors, whereas electrons are the important carriers for NPN transistors. Then, PNP transistors use a small output base current and a negative base voltage to control a much larger emitter-collector current. The construction of a PNP transistor consists of two P-type semiconductor materials either side of the N-type material as shown below.

**A PNP Transistor Configuration**
The construction and terminal voltages for an NPN transistor are shown above. The **PNP Transistor** has very similar characteristics to their NPN bipolar cousins, except that the polarities (or biasing) of the current and voltage directions are reversed for any one of the possible three configurations looked at in the first tutorial, Common Base, Common Emitter and Common Collector.

![PNP Transistor Connections](image)

**PNP Transistor Connections**

The voltage between the Base and Emitter ($V_{BE}$), is now negative at the Base and positive at the Emitter because for a PNP transistor, the Base terminal is always biased negative with respect to the Emitter. Also the Emitter supply voltage is positive with respect to the Collector ($V_{CE}$). So for a PNP transistor to conduct the Emitter is always more positive with respect to both the Base and the Collector.

The voltage sources are connected to a PNP transistor are as shown. This time the Emitter is connected to the supply voltage $V_{CC}$ with the load resistor, $R_L$ which limits the maximum current flowing through the device connected to the Collector terminal. The Base voltage $V_B$ which is biased negative with respect to the Emitter and is connected to the Base resistor $R_B$, which again is used to limit the maximum Base current.

To cause the Base current to flow in a PNP transistor the Base needs to be more negative than the Emitter (current must leave the base) by approx 0.7 volts for a silicon device or 0.3 volts for a germanium device with the formulas used to calculate the Base resistor, Base current or Collector current are the same as those used for an equivalent NPN transistor and is given as.

$$I_C = I_E - I_B$$

$$I_C = \beta I_B \quad I_B = \frac{I_C}{\beta}$$

Generally, the PNP transistor can replace NPN transistors in most electronic circuits, the only difference is the polarities of the voltages, and the directions of the current flow. PNP transistors can also be used as switching devices and an example of a PNP transistor switch is shown below.
A PNP Transistor Circuit

![PNP transistor circuit diagram](image)

The **Output Characteristics Curves** for a PNP transistor look very similar to those for an equivalent NPN transistor except that they are rotated by 180° to take account of the reverse polarity voltages and currents, (the currents flowing out of the Base and Collector in a PNP transistor are negative). The same dynamic load line can be drawn onto the I-V curves to find the PNP transistors operating points.

**Transistor Matching**

![Transistor matching diagram](image)
Complementary Transistors

You may think what is the point of having a PNP Transistor, when there are plenty of NPN Transistors available that can be used as an amplifier or solid-state switch?. Well, having two different types of transistors "PNP" and "NPN", can be a great advantage when designing amplifier circuits such as the Class B Amplifier which uses "Complementary" or "Matched Pair" transistors in its output stage or in reversible H-Bridge motor control circuits were we want to control the flow of current evenly in both directions.

A pair of corresponding NPN and PNP transistors with near identical characteristics to each other are called Complementary Transistors for example, a TIP3055 (NPN transistor) and the TIP2955 (PNP transistor) are good examples of complementary or matched pair silicon power transistors. They both have a DC current gain, Beta, (Ic/Ib) matched to within 10% and high Collector current of about 15A making them ideal for general motor control or robotic applications.

Also, class B amplifiers use complementary NPN and PNP in their power output stage design. The NPN transistor conducts for only the positive half of the signal while the PNP transistor conducts for negative half of the signal. This allows the amplifier to drive the required power through the load loudspeaker in both directions at the stated nominal impedance and power resulting in an output current which is likely to be in the order of several amps shared evenly between the two complementary transistors.

Identifying the PNP Transistor

We saw in the first tutorial of this transistors section, that transistors are basically made up of two Diodes connected together back-to-back. We can use this analogy to determine whether a transistor is of the PNP type or NPN type by testing its Resistance between the three different leads, Emitter, Base and Collector. By testing each pair of transistor leads in both directions with a multimeter will result in six tests in total with the expected resistance values in Ohm's given below.

1. Emitter-Base Terminals - The Emitter to Base should act like a normal diode and conduct one way only.

2. Collector-Base Terminals - The Collector-Base junction should act like a normal diode and conduct one way only.

3. Emitter-Collector Terminals - The Emitter-Collector should not conduct in either direction.

Transistor resistance values for a PNP transistor and a NPN transistor
<table>
<thead>
<tr>
<th>Between Transistor Terminals</th>
<th>PNP</th>
<th>NPN</th>
</tr>
</thead>
<tbody>
<tr>
<td>Collector</td>
<td>$R_{\text{HIGH}}$</td>
<td>$R_{\text{HIGH}}$</td>
</tr>
<tr>
<td>Collector</td>
<td>$R_{\text{LOW}}$</td>
<td>$R_{\text{HIGH}}$</td>
</tr>
<tr>
<td>Emitter</td>
<td>$R_{\text{HIGH}}$</td>
<td>$R_{\text{HIGH}}$</td>
</tr>
<tr>
<td>Emitter</td>
<td>$R_{\text{LOW}}$</td>
<td>$R_{\text{HIGH}}$</td>
</tr>
<tr>
<td>Base</td>
<td>$R_{\text{HIGH}}$</td>
<td>$R_{\text{LOW}}$</td>
</tr>
<tr>
<td>Base</td>
<td>$R_{\text{HIGH}}$</td>
<td>$R_{\text{LOW}}$</td>
</tr>
</tbody>
</table>

Then we can define a **PNP Transistor** as being normally "OFF" but a small output current and negative voltage at its Base (B) relative to its Emitter (E) will turn it "ON" allowing a much large Emitter-Collector current to flow. PNP transistors conduct when $V_e$ is much greater than $V_c$.

In the next tutorial about **Bipolar Transistors** instead of using the transistor as an amplifying device, we will look at the operation of the transistor in its saturation and cut-off regions when used as a solid-state switch. Bipolar transistor switches are used in many applications to switch a DC current "ON" or "OFF" such as LED’s which require only a few milliamps at low DC voltages, or relays which require higher currents at higher voltages.

### 1.3 The Transistor as a Switch

When used as an AC signal amplifier, the transistors Base biasing voltage is applied in such a way that it always operates within its "active" region, that is the linear part of the output characteristics curves are used. However, both the NPN & PNP type bipolar transistors can be made to operate as "ON/OFF" type solid state switches by biasing the transistors base differently to that of a signal amplifier. Solid state switches are one of the main applications for the use of transistors, and **transistor switches** can be used for controlling high power devices such as motors, solenoids or lamps, but they can also used in digital electronics and logic gate circuits.

If the circuit uses the **Bipolar Transistor as a Switch**, then the biasing of the transistor, either NPN or PNP is arranged to operate the transistor at both sides of the I-V characteristics curves we have seen previously. The areas of operation for a transistor switch are known as the **Saturation Region** and the **Cut-off Region**. This means then that we can ignore the operating Q-point biasing and voltage divider circuitry required for amplification, and use the transistor as a switch by driving it back and forth between its "fully-OFF" (cut-off) and "fully-ON" (saturation) regions as shown below.

**Basic NPN Transistor Switching Circuit**
The circuit resembles that of the Common Emitter circuit we looked at in the previous tutorials. The difference this time is that to operate the transistor as a switch the transistor needs to be turned either fully "OFF" (cut-off) or fully "ON" (saturated). An ideal transistor switch would have infinite circuit resistance between the Collector and Emitter when turned "fully-OFF" resulting in zero current flowing through it and zero resistance between the Collector and Emitter when turned "fully-ON", resulting in maximum current flow. In practice when the transistor is turned "OFF", small leakage currents flow through the transistor and when fully "ON" the device has a low resistance value causing a small saturation voltage ($V_{CE}$) across it. Even though the transistor is not a perfect switch, in both the cut-off and saturation regions the power dissipated by the transistor is at its minimum.

In order for the Base current to flow, the Base input terminal must be made more positive than the Emitter by increasing it above the 0.7 volts needed for a silicon device. By varying this Base-Emitter voltage ($V_{BE}$), the Base current is also altered and which in turn controls the amount of Collector current flowing through the transistor as previously discussed. When maximum Collector current flows the transistor is said to be Saturated. The value of the Base resistor determines how much input voltage is required and corresponding Base current to switch the transistor fully "ON".

Example No1

Using the transistor values from the previous tutorials of: $\beta = 200$, $I_c = 4mA$ and $I_b = 20\mu A$, find the value of the Base resistor ($R_b$) required to switch the load "ON" when the input terminal voltage exceeds 2.5$V$.

$$R_B = \frac{V_{In} - V_{BE}}{I_B} = \frac{2.5V - 0.7V}{20 \times 10^{-6}} = 90k\Omega$$

The next lowest preferred value is: 82k$\Omega$, this guarantees the transistor switch is always saturated.

Example No2
Again using the same values, find the minimum Base current required to turn the transistor "fully-ON" (saturated) for a load that requires 200mA of current when the input voltage is increased to 5.0V. Also calculate the new value of Rb.

Transistor Base current:

\[ I_B = \frac{I_C}{\beta} = \frac{200\text{mA}}{200} = 1\text{mA} \]

Transistor Base resistance:

\[ R_B = \frac{V_{\text{In}} - V_{\text{BE}}}{I_B} = \frac{5.0\text{V} - 0.7\text{V}}{1 \times 10^{-3}} = 4.3k\Omega \]

Transistor switches are used for a wide variety of applications such as interfacing large current or high voltage devices like motors, relays or lamps to low voltage digital logic IC's or gates like AND gates or OR gates. Here, the output from a digital logic gate is only +5v but the device to be controlled may require a 12 or even 24 volts supply. Or the load such as a DC Motor may need to have its speed controlled using a series of pulses (Pulse Width Modulation). Transistor switches will allow us to do this faster and more easily than with conventional mechanical switches.

**Digital Logic Transistor Switch**

![Diagram of Digital Logic Transistor Switch]

The base resistor, Rb is required to limit the output current from the logic gate.

**PNP Transistor Switch**

We can also use PNP transistors as switches, the difference this time is that the load is connected to ground (0v) and the PNP transistor switches power to it. To turn the PNP transistor as a switch "ON" the Base terminal is connected to ground or zero volts (LOW) as shown.
PNP Transistor Switching Circuit

The equations for calculating the Base resistance, Collector current and voltages are exactly the same as for the previous NPN transistor switch. The difference this time is that we are switching power with a PNP transistor (sourcing current) instead of switching ground with an NPN transistor (sinking current).

Unijunction transistor

Although a unijunction transistor is not a thyristor, this device can trigger larger thyristors with a pulse at base B1. A unijunction transistor is composed of a bar of N-type silicon having a P-type connection in the middle. See Figure below(a). The connections at the ends of the bar are known as bases B1 and B2; the P-type mid-point is the emitter. With the emitter disconnected, the total resistance RBBO, a datasheet item, is the sum of RB1 and RB2 as shown in Figure below(b). RBBO ranges from 4-12kΩ for different device types. The intrinsic standoff ratio η is the ratio of RB1 to RBBO. It varies from 0.4 to 0.8 for different devices. The schematic symbol is Figure below(c)

Unijunction transistor: (a) Construction, (b) Model, (c) Symbol
The Unijunction emitter current vs voltage characteristic curve (Figure below(a) ) shows that as $V_E$ increases, current $I_E$ increases up to $I_P$ at the peak point. Beyond the peak point, current increases as voltage decreases in the negative resistance region. The voltage reaches a minimum at the valley point. The resistance of RB1, the saturation resistance is lowest at the valley point.

IP and IV, are datasheet parameters; For a 2n2647, IP and IV are 2µA and 4mA, respectively. [AMS] VP is the voltage drop across RB1 plus a 0.7V diode drop; see Figure below(b). VV is estimated to be approximately 10% of VBB.

Unijunction transistor: (a) emitter characteristic curve, (b) model for $V_P$.

The relaxation oscillator in Figure below is an application of the unijunction oscillator. RE charges CE until the peak point. The unijunction emitter terminal has no effect on the capacitor until this point is reached. Once the capacitor voltage, $V_E$, reaches the peak voltage point $V_P$, the lowered emitter-base1 E-B1 resistance quickly discharges the capacitor. Once the capacitor discharges below the valley point $V_V$, the E-RB1 resistance reverts back to high resistance, and the capacitor is free to charge again.

Unijunction transistor relaxation oscillator and waveforms. Oscillator drives SCR.
During capacitor discharge through the E-B1 saturation resistance, a pulse may be seen on the external B1 and B2 load resistors, Figure above. The load resistor at B1 needs to be low to not affect the discharge time. The external resistor at B2 is optional. It may be replaced by a short circuit. The approximate frequency is given by $1/T = RC$. A more accurate expression for frequency is given in Figure above.

The charging resistor $R_E$ must fall within certain limits. It must be small enough to allow $I_P$ to flow based on the $V_{BB}$ supply less $V_P$. It must be large enough to supply $I_V$ based on the $V_{BB}$ supply less $V_V$.

[MHW] The equations and an example for a 2n2647:

$$V_P = 0.7 + \eta V_{MB}$$
$$V_V = 0.10(V_{BB})$$

$$\frac{V_{BB} \cdot V_V}{I_V} < R_E < \frac{V_{BB} \cdot V_P}{I_P}$$

$$\frac{10 - 1}{8mA} < R_E < \frac{10 - 8.2}{2\mu A}$$

$$1.125k < R_E < 900k$$

Programmable Unijunction Transistor (PUT): Although the unijunction transistor is listed as obsolete (read expensive if obtainable), the programmable unijunction transistor is alive and well. It is inexpensive and in production. Though it serves a function similar to the unijunction transistor, the PUT is a three terminal thyristor. The PUT shares the four-layer structure typical of thyristors shown in Figure below. Note that the gate, an N-type layer near the anode, is known as an “anode gate”. Moreover, the gate lead on the schematic symbol is attached to the anode end of the symbol.

The characteristic curve for the programmable unijunction transistor in Figure above is similar to that of the unijunction transistor. This is a plot of anode current $I_A$ versus anode voltage $V_A$. The gate lead voltage sets, programs, the peak anode voltage $V_P$. As anode current increases, voltage increases up to the peak point. Thereafter, increasing current results in decreasing voltage, down to the valley point.
The PUT equivalent of the unijunction transistor is shown in Figure below. External PUT resistors $R_1$ and $R_2$ replace unijunction transistor internal resistors $R_{B1}$ and $R_{B2}$, respectively. These resistors allow the calculation of the intrinsic standoff ratio $\eta$.

$$R_{BB0} = R_1 + R_2$$
$$\eta = \frac{R_1}{R_1 + R_2}$$
$$V_S = \eta V_{BB}$$
$$V_P = V_T + V_S$$
$$R_G = \frac{R_1 R_2}{R_1 + R_2}$$

PUT equivalent of unijunction transistor

Figure below shows the PUT version of the unijunction relaxation oscillator Figure previous. Resistor $R$ charges the capacitor until the peak point, Figure previous, then heavy conduction moves the operating point down the negative resistance slope to the valley point. A current spike flows through the cathode during capacitor discharge, developing a voltage spike across the cathode resistors. After capacitor discharge, the operating point resets back to the slope up to the peak point.

**Problem:** What is the range of suitable values for $R$ in Figure above, a relaxation oscillator? The charging resistor must be small enough to supply enough current to raise the anode to $V_P$ the peak point (Figure previous) while charging the capacitor. Once $V_P$ is reached, anode voltage decreases as current increases (negative resistance), which moves the operating point to the valley. It is the job of the capacitor to supply the valley current $I_V$. Once it is discharged, the operating point resets back to the upward slope to the peak point. The resistor must be large enough so that it will never supply the high valley current $I_P$. If the charging resistor ever could supply that much current, the resistor would supply the valley current after
the capacitor was discharged and the operating point would never reset back to the high resistance condition to the left of the peak point.

We select the same \( V_{BB} = 10V \) used for the unijunction transistor example. We select values of \( R_1 \) and \( R_2 \) so that \( \eta \) is about \( 2/3 \). We calculate \( \eta \) and \( V_S \). The parallel equivalent of \( R_1, R_2 \) is \( R_G \), which is only used to make selections from Table below. Along with \( V_S = 10 \), the closest value to our 6.3, we find \( V_T = 0.6V \), in Table below and calculate \( V_P \).

\[
\eta = \frac{R_1}{R_1 + R_2} \quad \eta = \frac{27}{27 + 16} = 0.6279
\]

\[
V_S = \eta V_{BB} \quad V_S = 0.6279(10) = 6.279V
\]

\[
R_G = \frac{R_1R_2}{R_1 + R_2} \quad R_G = \frac{27k \cdot 16k}{27k + 16k} = 10k
\]

For \( R_G = 10k \) and \( V_S = 10V \), \( V_T = 0.6V \)

\[
V_P = V_T + V_S \quad V_P = 0.6 + 6.3 = 6.9V
\]

We also find \( I_P \) and \( I_V \), the peak and valley currents, respectively in Table below. We still need \( V_V \), the valley voltage. We used 10% of \( V_{BB} = 1V \), in the previous unijunction example. Consulting the datasheet, we find the forward voltage \( V_F = 0.8V \) at \( I_F = 50mA \). The valley current \( I_V = 70\mu A \) is much less than \( I_F = 50mA \). Therefore, \( V_V \) must be less than \( V_F = 0.8V \). How much less? To be safe we set \( V_V = 0V \). This will raise the lower limit on the resistor range a little.

\[
V_V = 0.10(V_{BB}) \quad \text{not used} \quad V_V = 0V
\]

\[
\frac{V_{BB} - V_V}{I_V} < R_E < \frac{V_{BB} - V_P}{I_P} \quad \frac{10 - 0}{70\mu A} < R_E < \frac{10 - 6.9}{4\mu A} \quad 143k < R_E < 755k
\]

Choosing \( R > 143k \) guarantees that the operating point can reset from the valley point after capacitor discharge. \( R < 755k \) allows charging up to \( V_P \) at the peak point.

Figure below show the PUT relaxation oscillator with the final resistor values. A practical application of a PUT triggering an SCR is also shown. This circuit needs a \( V_{BB} \) unfiltered supply (not shown) divided down from the bridge rectifier to reset the relaxation oscillator after each power zero crossing. The variable resistor should have a minimum resistor in series with it to prevent a low pot setting from hanging at the valley point.
PUT relaxation oscillator with component values. PUT drives SCR lamp dimmer.

1.4 Silicon-Controlled Rectifiers, or SCRs

Shockley diodes are curious devices, but rather limited in application. Their usefulness may be expanded, however, by equipping them with another means of latching. In doing so, each becomes true amplifying devices (if only in an on/off mode), and we refer to these as silicon-controlled rectifiers, or SCRs.

The progression from Shockley diode to SCR is achieved with one small addition, actually nothing more than a third wire connection to the existing PNPN structure. (Figure below)

The Silicon-Controlled Rectifier (SCR)

If an SCR’s gate is left floating (disconnected), it behaves exactly as a Shockley diode. It may be latched by break-over voltage or by exceeding the critical rate of voltage rise between anode and cathode, just as with the Shockley diode. Dropout is accomplished by reducing current until one or both internal transistors fall into cutoff mode, also like the Shockley diode. However, because the gate terminal connects directly to the base of the lower transistor, it may be used as an alternative means to latch the SCR. By applying a small voltage between gate and cathode, the lower transistor will be forced on by the resulting base current, which will cause the upper transistor to conduct, which then supplies the lower transistor’s base with current so that it no longer needs to be activated by a gate voltage. The necessary gate current to initiate latch-up, of course, will be much lower than the current through the SCR from cathode to anode, so the SCR does achieve a measure of amplification.
This method of securing SCR conduction is called triggering, and it is by far the most common way that
SCRs are latched in actual practice. In fact, SCRs are usually chosen so that their breakover voltage is far
beyond the greatest voltage expected to be experienced from the power source, so that it can be turned on
only by an intentional voltage pulse applied to the gate.

It should be mentioned that SCRs may sometimes be turned off by directly shorting their gate and cathode
terminals together, or by "reverse-triggering" the gate with a negative voltage (in reference to the
cathode), so that the lower transistor is forced into cutoff. I say this is "sometimes" possible because it
involves shunting all of the upper transistor's collector current past the lower transistor's base. This
current may be substantial, making triggered shut-off of an SCR difficult at best. A variation of the SCR,
called a Gate-Turn-Off thyristor, or GTO, makes this task easier. But even with a GTO, the gate current
required to turn it off may be as much as 20% of the anode (load) current! The schematic symbol for a
GTO is shown in the following illustration: (Figure below)

1.5 The Gate Turn-Off Thyristors (GTO)

SCRs and GTOs share the same equivalent schematics (two transistors connected in a positive-feedback
fashion), the only differences being details of construction designed to grant the NPN transistor a greater
β than the PNP. This allows a smaller gate current (forward or reverse) to exert a greater degree of control
over conduction from cathode to anode, with the PNP transistor's latched state being more dependent
upon the NPN's than vice versa. The Gate-Turn-Off thyristor is also known by the name of Gate-
Controlled Switch, or GCS.

A rudimentary test of SCR function, or at least terminal identification, may be performed with an
ohmmeter. Because the internal connection between gate and cathode is a single PN junction, a meter
should indicate continuity between these terminals with the red test lead on the gate and the black test
lead on the cathode like this: (Figure below)
Rudimentary test of SCR

All other continuity measurements performed on an SCR will show "open" ("OL" on some digital multimeter displays). It must be understood that this test is very crude and does not constitute a comprehensive assessment of the SCR. It is possible for an SCR to give good ohmmeter indications and still be defective. Ultimately, the only way to test an SCR is to subject it to a load current.

If you are using a multimeter with a "diode check" function, the gate-to-cathode junction voltage indication you get may or may not correspond to what's expected of a silicon PN junction (approximately 0.7 volts). In some cases, you will read a much lower junction voltage: mere hundredths of a volt. This is due to an internal resistor connected between the gate and cathode incorporated within some SCRs. This resistor is added to make the SCR less susceptible to false triggering by spurious voltage spikes, from circuit "noise" or from static electric discharge. In other words, having a resistor connected across the gate-cathode junction requires that a strong triggering signal (substantial current) be applied to latch the SCR. This feature is often found in larger SCRs, not on small SCRs. Bear in mind that an SCR with an internal resistor connected between gate and cathode will indicate continuity in both directions between those two terminals: (Figure below)

Larger SCRs have gate to cathode resistor.

"Normal" SCRs, lacking this internal resistor, are sometimes referred to as sensitive gate SCRs due to their ability to be triggered by the slightest positive gate signal.

The test circuit for an SCR is both practical as a diagnostic tool for checking suspected SCRs and also an excellent aid to understanding basic SCR operation. A DC voltage source is used for powering the circuit, and two pushbutton switches are used to latch and unlatch the SCR, respectively: (Figure below)
Actuating the normally-open "on" pushbutton switch connects the gate to the anode, allowing current from the negative terminal of the battery, through the cathode-gate PN junction, through the switch, through the load resistor, and back to the battery. This gate current should force the SCR to latch on, allowing current to go directly from cathode to anode without further triggering through the gate. When the "on" pushbutton is released, the load should remain energized.

Pushing the normally-closed "off" pushbutton switch breaks the circuit, forcing current through the SCR to halt, thus forcing it to turn off (low-current dropout).

If the SCR fails to latch, the problem may be with the load and not the SCR. A certain minimum amount of load current is required to hold the SCR latched in the "on" state. This minimum current level is called the holding current. A load with too great a resistance value may not draw enough current to keep an SCR latched when gate current ceases, thus giving the false impression of a bad (unlatchable) SCR in the test circuit. Holding current values for different SCRs should be available from the manufacturers. Typical holding current values range from 1 milliamp to 50 milliamps or more for larger units.

For the test to be fully comprehensive, more than the triggering action needs to be tested. The forward breakover voltage limit of the SCR could be tested by increasing the DC voltage supply (with no pushbuttons actuated) until the SCR latches all on its own. Beware that a breakover test may require very high voltage: many power SCRs have breakover voltage ratings of 600 volts or more! Also, if a pulse voltage generator is available, the critical rate of voltage rise for the SCR could be tested in the same way: subject it to pulsing supply voltages of different V/time rates with no pushbutton switches actuated and see when it latches.

In this simple form, the SCR test circuit could suffice as a start/stop control circuit for a DC motor, lamp, or other practical load: (Figure below)
Another practical use for the SCR in a DC circuit is as a crowbar device for overvoltage protection. A "crowbar" circuit consists of an SCR placed in parallel with the output of a DC power supply, for placing a direct short-circuit on the output of that supply to prevent excessive voltage from reaching the load. Damage to the SCR and power supply is prevented by the judicious placement of a fuse or substantial series resistance ahead of the SCR to limit short-circuit current: (Figure below)

Some device or circuit sensing the output voltage will be connected to the gate of the SCR, so that when an overvoltage condition occurs, voltage will be applied between the gate and cathode, triggering the SCR and forcing the fuse to blow. The effect will be approximately the same as dropping a solid steel crowbar directly across the output terminals of the power supply, hence the name of the circuit.

Most applications of the SCR are for AC power control, despite the fact that SCRs are inherently DC (unidirectional) devices. If bidirectional circuit current is required, multiple SCRs may be used, with one or more facing each direction to handle current through both half-cycles of the AC wave. The primary reason SCRs are used at all for AC power control applications is the unique response of a thyristor to an alternating current. As we saw, the thyratron tube (the electron tube version of the SCR) and the DIAC, a hysteretic device triggered on during a portion of an AC half-cycle will latch and remain on throughout the remainder of the half-cycle until the AC current decreases to zero, as it must to begin the next half-cycle. Just prior to the zero-crossover point of the current waveform, the thyristor will turn off due to insufficient current (this behavior is also known as natural commutation) and must be fired again during the next cycle. The result is a circuit current equivalent to a "chopped up" sine wave. For review, here is
the graph of a DIAC's response to an AC voltage whose peak exceeds the breakover voltage of the DIAC: (Figure below)

With the DIAC, that breakover voltage limit was a fixed quantity. With the SCR, we have control over exactly when the device becomes latched by triggering the gate at any point in time along the waveform. By connecting a suitable control circuit to the gate of an SCR, we can "chop" the sine wave at any point to allow for time-proportioned power control to a load.

Take the circuit in Figure below as an example. Here, an SCR is positioned in a circuit to control power to a load from an AC source.

Being a unidirectional (one-way) device, at most we can only deliver half-wave power to the load, in the half-cycle of AC where the supply voltage polarity is positive on the top and negative on the bottom. However, for demonstrating the basic concept of time-proportional control, this simple circuit is better than one controlling full-wave power (which would require two SCRs).

With no triggering to the gate, and the AC source voltage well below the SCR's breakover voltage rating, the SCR will never turn on. Connecting the SCR gate to the anode through a standard rectifying diode (to prevent reverse current through the gate in the event of the SCR containing a built-in gate-cathode resistor), will allow the SCR to be triggered almost immediately at the beginning of every positive half-cycle: (Figure below).
Gate connected directly to anode through a diode; nearly complete half-wave current through load.

We can delay the triggering of the SCR, however, by inserting some resistance into the gate circuit, thus increasing the amount of voltage drop required before enough gate current triggers the SCR. In other words, if we make it harder for electrons to flow through the gate by adding a resistance, the AC voltage will have to reach a higher point in its cycle before there will be enough gate current to turn the SCR on. The result is in Figure below.

Resistance inserted in gate circuit; less than half-wave current through load.

With the half-sine wave chopped up to a greater degree by delayed triggering of the SCR, the load receives less average power (power is delivered for less time throughout a cycle). By making the series gate resistor variable, we can make adjustments to the time-proportioned power: (Figure below)
Increasing the resistance raises the threshold level, causing less power to be delivered to the load. Decreasing the resistance lowers the threshold level, causing more power to be delivered to the load.

Unfortunately, this control scheme has a significant limitation. In using the AC source waveform for our SCR triggering signal, we limit control to the first half of the waveform's half-cycle. In other words, it is not possible for us to wait until after the wave's peak to trigger the SCR. This means we can turn down the power only to the point where the SCR turns on at the very peak of the wave: (Figure below)

Raising the trigger threshold any more will cause the circuit to not trigger at all, since not even the peak of the AC power voltage will be enough to trigger the SCR. The result will be no power to the load.

An ingenious solution to this control dilemma is found in the addition of a phase-shifting capacitor to the circuit: (Figure below)
Addition of a phase-shifting capacitor to the circuit

The smaller waveform shown on the graph is voltage across the capacitor. For the sake of illustrating the phase shift, I am assuming a condition of maximum control resistance where the SCR is not triggering at all with no load current, save for what little current goes through the control resistor and capacitor. This capacitor voltage will be phase-shifted anywhere from 0° to 90° lagging behind the power source AC waveform. When this phase-shifted voltage reaches a high enough level, the SCR will trigger.

With enough voltage across the capacitor to periodically trigger the SCR, the resulting load current waveform will look something like Figure below)

![Phase-shifted signal triggers SCR into conduction.](image)

Because the capacitor waveform is still rising after the main AC power waveform has reached its peak, it becomes possible to trigger the SCR at a threshold level beyond that peak, thus chopping the load current wave further than it was possible with the simpler circuit. In reality, the capacitor voltage waveform is a bit more complex that what is shown here, its sinusoidal shape distorted every time the SCR latches on. However, what I'm trying to illustrate here is the delayed triggering action gained with the phase-shifting RC network; thus, a simplified, undistorted waveform serves the purpose well.

SCRs may also be triggered, or "fired," by more complex circuits. While the circuit previously shown is sufficient for a simple application like a lamp control, large industrial motor controls often rely on more sophisticated triggering methods. Sometimes, pulse transformers are used to couple a triggering circuit to the gate and cathode of an SCR to provide electrical isolation between the triggering and power circuits: (Figure below)

![Transformer coupling of trigger signal provides isolation.](image)
When multiple SCRs are used to control power, their cathodes are often not electrically common, making it difficult to connect a single triggering circuit to all SCRs equally. An example of this is the controlled bridge rectifier shown in Figure below.

![Controlled Bridge Rectifier Diagram](image)

In any bridge rectifier circuit, the rectifying diodes (in this example, the rectifying SCRs) must conduct in opposite pairs. SCR₁ and SCR₃ must be fired simultaneously, and SCR₂ and SCR₄ must be fired together as a pair. As you will notice, though, these pairs of SCRs do not share the same cathode connections, meaning that it would not work to simply parallel their respective gate connections and connect a single voltage source to trigger both: (Figure below)

![Parallel Gate Connection Diagram](image)

This strategy will not work for triggering SCR₂ and SCR₄ as a pair.

Although the triggering voltage source shown will trigger SCR₄, it will not trigger SCR₂ properly because the two thyristors do not share a common cathode connection to reference that triggering voltage. Pulse transformers connecting the two thyristor gates to a common triggering voltage source will work, however: (Figure below)
Transformer coupling of the gates allows triggering of SCR$_2$ and SCR$_4$.

Bear in mind that this circuit only shows the gate connections for two out of the four SCRs. Pulse transformers and triggering sources for SCR$_1$ and SCR$_3$, as well as the details of the pulse sources themselves, have been omitted for the sake of simplicity.

Controlled bridge rectifiers are not limited to single-phase designs. In most industrial control systems, AC power is available in three-phase form for maximum efficiency, and solid-state control circuits are built to take advantage of that. A three-phase controlled rectifier circuit built with SCRs, without pulse transformers or triggering circuitry shown, would look like Figure below.

1.6 Recommended Questions

1. Explain the clipping above and below the reference voltage in a basic parallel clipper.
2. Explain the clipping above and below the reference voltage in the basic series clipper.
3. Explain the clippers with voltage divider circuit.
4. Explain the negative & positive clamper circuit.
5. Draw and explain the characteristics of Schottky diode.
6. What is feature of a Varactor diode?
7. What is a voltage multiplier circuit? Explain the operation of a full wave voltage doubler circuit. (Jan-2007)
8. Define diffusion capacitance. Derive an expression for the same. (July-2007)
9. Draw the piece wise linear V-I characteristics of a P-N junction diode. Give the circuit model for the ON state and OFF state. (July-2007)
10. Discuss voltage doubler circuit. (Jan-2008)
11. Define regulation and derive equation for a full wave circuit. (Jan-2008)
Unit – 2: Bipolar Junction Transistors

2.1 Bipolar junction transistor (BJT)

The bipolar junction transistor (BJT) was the first solid-state amplifier element and started the solid-state electronics revolution. Bardeen, Brattain and Shockley, while at Bell Laboratories, invented it in 1948 as part of a post-war effort to replace vacuum tubes with solid-state devices. Solid-state rectifiers were already in use at the time and were preferred over vacuum diodes because of their smaller size, lower weight and higher reliability. A solid-state replacement for a vacuum triode was expected to yield similar advantages. The work at Bell Laboratories was highly successful and culminated in Bardeen, Brattain and Shockley receiving the Nobel Prize in 1956.

Their work led them first to the point-contact transistor and then to the bipolar junction transistor. They used germanium as the semiconductor of choice because it was possible to obtain high purity material. The extraordinarily large diffusion length of minority carriers in germanium provided functional structures despite the large dimensions of the early devices.

Since then, the technology has progressed rapidly. The development of a planar process yielded the first circuits on a chip and for a decade, bipolar transistor operational amplifiers, like the 741, and digital TTL circuits were for a long time the workhorses of any circuit designer.

The spectacular rise of the MOSFET market share during the last decade has completely removed the bipolar transistor from center stage. Almost all logic circuits, microprocessor and memory chips contain exclusively MOSFETs.

Nevertheless, bipolar transistors remain important devices for ultra-high-speed discrete logic circuits such as emitter coupled logic (ECL), power-switching applications and in microwave power amplifiers. Heterojunction bipolar transistors (HBTs) have emerged as the device of choice for cell phone amplifiers and other demanding applications.

In this chapter we first present the structure of the bipolar transistor and show how a three-layer structure with alternating n-type and p-type regions can provide current and voltage amplification. We then present the ideal transistor model and derive an expression for the current gain in the forward active mode of operation. Next, we discuss the non-ideal effects, the modulation of the base width and recombination in the depletion region of the base-emitter junction. A discussion of transit time effects, BJT circuit models, HBTs, BJT technology and bipolar power devices completes this chapter.

Structure and principle of operation

A bipolar junction transistor consists of two back-to-back p-n junctions, who share a thin common region with width, $w_B$. Contacts are made to all three regions, the two outer regions called the emitter and collector and the middle region called the base. The structure of an npn
bipolar transistor is shown in Figure (a). The device is called “bipolar” since its operation involves both types of mobile carriers, electrons and holes.

Since the device consists of two back-to-back diodes, there are depletion regions between the quasi-neutral regions. The width of the quasi neutral regions in the emitter, base and collector are indicated with the symbols \( w_E', w_B' \) and \( w_C' \) and are calculated from

\[
\begin{align*}
  w_E' &= w_E - x_{n, BE} \\
  w_B' &= w_B - x_{p, BE} - x_{p, BC}
\end{align*}
\]
\[ w_C' = w_C - x_{n,BC} \]

where the depletion region widths are given by:

\[
x_{n,BE} = \sqrt{\frac{2e_s (\phi_{BE} - V_{BE})}{q} N_B \left( \frac{1}{N_B + N_E} \right)}
\]

\[
x_{p,BE} = \sqrt{\frac{2e_s (\phi_{BE} - V_{BE})}{q} N_E \left( \frac{1}{N_B + N_E} \right)}
\]

\[
x_{p,BC} = \sqrt{\frac{2e_s (\phi_{BC} - V_{BC})}{q} N_C \left( \frac{1}{N_B + N_C} \right)}
\]

\[
x_{n,BC} = \sqrt{\frac{2e_s (\phi_{BC} - V_{BC})}{q} N_B \left( \frac{1}{N_B + N_C} \right)}
\]

With

\[
\phi_{BE} = V_I \ln \frac{N_B N_B}{n_i^2}
\]

\[
\phi_{BC} = V_I \ln \frac{N_C N_B}{n_i^2}
\]

The sign convention of the currents and voltage is indicated on Figure 5.2.1(a). The base and collector current are positive if a positive current goes into the base or collector contact. The emitter current is positive for a current coming out of the emitter contact. This also implies that the emitter current, \( I_E \), equals the sum of the base current, \( I_B \), and the collector current, \( I_C \):

\[ I_E = I_C + I_B \]

The base-emitter voltage and the base-collector voltage are positive if a positive voltage is applied to the base contact relative to the emitter and collector respectively.

The operation of the device is illustrated with Figure 5.2.1(b). We consider here only the forward active bias mode of operation, obtained by forward biasing the base-emitter junction and reverse biasing the base-collector junction. To simplify the discussion further, we also set \( V_{CE} = 0 \). The corresponding energy band diagram is shown in Figure 5.2.2. Electrons diffuse from the...
emitter into the base and holes diffuse from the base into the emitter. This carrier diffusion is identical to that in a p-n junction. However, what is different is that the electrons can diffuse as minority carriers through the quasi-neutral base. Once the electrons arrive at the base-collector depletion region, they are swept through the depletion layer due to the electric field. These electrons contribute to the collector current. In addition, there are two more currents, the base recombination current, indicated on Figure 5.2.2 by the vertical arrow, and the base-emitter depletion layer recombination current, $I_{r,d}$ (not shown).

Figure 5.2.2. : Energy band diagram of a bipolar transistor biased in the forward active mode.

The total emitter current is the sum of the electron diffusion current, $I_{E,n}$, the hole diffusion current, $I_{E,p}$ and the base-emitter depletion layer recombination current, $I_{r,d}$.

$$I_B = I_{E,n} + I_{E,p} + I_{r,d}$$

The total collector current is the electron diffusion current, $I_{E,n}$, minus the base recombination current, $I_{r,B}$.

$$I_C = I_{E,n} - I_{r,B}$$

The base current is the sum of the hole diffusion current, $I_{E,p}$, the base recombination current, $I_{r,B}$ and the base-emitter depletion layer recombination current, $I_{r,d}$.

$$I_B = I_{E,p} + I_{r,B} + I_{r,d}$$

The transport factor, is defined as the ratio of the collector and emitter current:
Using Kirchoff’s current law and the sign convention shown in Figure 5.2.1(a), we find that the base current equals the difference between the emitter and collector current. The current gain, \( \alpha \), is defined as the ratio of the collector and base current and equals:

\[
\alpha = \frac{I_C}{I_B}
\]

This explains how a bipolar junction transistor can provide current amplification. If the collector current is almost equal to the emitter current, the transport factor, \( \beta \), approaches one. The current gain, \( \beta \), can therefore become much larger than one.

To facilitate further analysis, we now rewrite the transport factor, \( \beta \), as the product of the emitter efficiency, \( \gamma_E \), the base transport factor, \( \alpha_T \), and the depletion layer recombination factor, \( \delta_r \).

\[
\alpha = \alpha_T \gamma_E \delta_r
\]

The emitter efficiency, \( \gamma_E \), is defined as the ratio of the electron current in the emitter, \( I_{E,n} \), to the sum of the electron and hole current diffusing across the base-emitter junction, \( I_{E,n} + I_{E,p} \).

\[
\gamma_E = \frac{I_{E,n}}{I_{E,n} + I_{E,p}}
\]

The base transport factor, \( \alpha_T \), equals the ratio of the current due to electrons injected in the collector, to the current due to electrons injected in the base.

\[
\alpha_T = \frac{I_{E,n} - I_{r,B}}{I_{E,n}}
\]

Recombination in the depletion-region of the base-emitter junction further reduces the current gain, as it increases the emitter current without increasing the collector current. The depletion layer recombination factor, \( \delta_r \), equals the ratio of the current due to electron and hole diffusion across the base-emitter junction to the total emitter current:

\[
\delta_r = \frac{I_E - I_{r,d}}{I_E}
\]
Example 5.1  A bipolar transistor with an emitter current of 1 mA has an emitter efficiency of 0.99, a base transport factor of 0.995 and a depletion layer recombination factor of 0.998. Calculate the base current, the collector current, the transport factor and the current gain of the transistor.

Solution: The transport factor and current gain are:

\[ \alpha = \gamma E \alpha_T \delta_T = 0.99 \times 0.995 \times 0.998 = 0.983 \]

and

\[ \beta = \frac{\alpha}{1 - \alpha} = 58.1 \]

The collector current then equals

\[ I_C = \alpha I_E = 0.983 \text{ mA} \]

And the base current is obtained from:

\[ I_B = I_E - I_C = 17 \mu \text{A} \]

2.2 The Field Effect Transistor

In the Bipolar Junction Transistor tutorials, we saw that the output Collector current of the transistor is proportional to input current flowing into the Base terminal of the device, thereby making the bipolar transistor a "CURRENT" operated device (Beta model). The Field Effect Transistor, or simply FET however, uses the voltage that is applied to their input terminal, called the Gate to control the current flowing through them resulting in the output current being proportional to the input voltage. As their operation relies on an electric field (hence the name field effect) generated by the input Gate voltage, this then makes the Field Effect Transistor a "VOLTAGE" operated device.
The **Field Effect Transistor** is a three terminal unipolar semiconductor device that has very similar characteristics to those of their Bipolar Transistor counterparts i.e., high efficiency, instant operation, robust and cheap and can be used in most electronic circuit applications to replace their equivalent bipolar junction transistors (BJT) cousins.

Field effect transistors can be made much smaller than an equivalent BJT transistor and along with their low power consumption and power dissipation makes them ideal for use in integrated circuits such as the CMOS range of digital logic chips.

We remember from the previous tutorials that there are two basic types of Bipolar Transistor construction, NPN and PNP, which basically describes the physical arrangement of the P-type and N-type semiconductor materials from which they are made. This is also true of FET's as there are also two basic classifications of Field Effect Transistor, called the N-channel FET and the P-channel FET.

The field effect transistor is a three terminal device that is constructed with no PN-junctions within the main current carrying path between the Drain and the Source terminals, which correspond in function to the Collector and the Emitter respectively of the bipolar transistor. The current path between these two terminals is called the "channel" which may be made of either a P-type or an N-type semiconductor material. The control of current flowing in this channel is achieved by varying the voltage applied to the Gate. As their name implies, Bipolar Transistors are "Bipolar" devices because they operate with both types of charge carriers, Holes and Electrons. The Field Effect Transistor on the other hand is a "Unipolar" device that depends only on the conduction of electrons (N-channel) or holes (P-channel).

The **Field Effect Transistor** has one major advantage over its standard bipolar transistor cousins, in that their input impedance, (R\text{in}) is very high, (thousands of Ohms), while the BJT is comparatively low. This very high input impedance makes them very sensitive to input voltage signals, but the price of this high sensitivity also means that they can be easily damaged by static electricity. There are two main types of field effect transistor, the **Junction Field Effect Transistor** or **JFET** and the **Insulated-gate Field Effect Transistor** or **IGFET**), which is more commonly known as the standard **Metal Oxide Semiconductor Field Effect Transistor** or **MOSFET** for short.
2.3 The Junction Field Effect Transistor

We saw previously that a bipolar junction transistor is constructed using two PN-junctions in the main current carrying path between the Emitter and the Collector terminals. The Junction Field Effect Transistor (JUGFET or JFET) has no PN-junctions but instead has a narrow piece of high-resistivity semiconductor material forming a "Channel" of either N-type or P-type silicon for the majority carriers to flow through with two ohmic electrical connections at either end commonly called the Drain and the Source respectively.

There are two basic configurations of junction field effect transistor, the N-channel JFET and the P-channel JFET. The N-channel JFET's channel is doped with donor impurities meaning that the flow of current through the channel is negative (hence the term N-channel) in the form of electrons. Likewise, the P-channel JFET's channel is doped with acceptor impurities meaning that the flow of current through the channel is positive (hence the term P-channel) in the form of holes. N-channel JFET's have a greater channel conductivity (lower resistance) than their equivalent P-channel types, since electrons have a higher mobility through a conductor compared to holes. This makes the N-channel JFET's a more efficient conductor compared to their P-channel counterparts.

We have said previously that there are two ohmic electrical connections at either end of the channel called the Drain and the Source. But within this channel there is a third electrical connection which is called the Gate terminal and this can also be a P-type or N-type material forming a PN-junction with the main channel. The relationship between the connections of a junction field effect transistor and a bipolar junction transistor are compared below.

2.4 Comparison of connections between a JFET and a BJT

<table>
<thead>
<tr>
<th>Bipolar Transistor</th>
<th>Field Effect Transistor</th>
</tr>
</thead>
<tbody>
<tr>
<td>Emitter - (E) &gt;&gt; Source - (S)</td>
<td></td>
</tr>
<tr>
<td>Base - (B) &gt;&gt; Gate - (G)</td>
<td></td>
</tr>
<tr>
<td>Collector - (C) &gt;&gt; Drain - (D)</td>
<td></td>
</tr>
</tbody>
</table>

The symbols and basic construction for both configurations of JFETs are shown below.
The semiconductor "channel" of the Junction Field Effect Transistor is a resistive path through which a voltage $V_{DS}$ causes a current $I_D$ to flow. The JFET can conduct current equally well in either direction. A voltage gradient is thus formed down the length of the channel with this voltage becoming less positive as we go from the Drain terminal to the Source terminal. The PN-junction therefore has a high reverse bias at the Drain terminal and a lower reverse bias at the Source terminal. This bias causes a "depletion layer" to be formed within the channel and whose width increases with the bias.

The magnitude of the current flowing through the channel between the Drain and the Source terminals is controlled by a voltage applied to the Gate terminal, which is a reverse-biased. In an N-channel JFET this Gate voltage is negative while for a P-channel JFET the Gate voltage is positive. The main difference between the JFET and a BJT device is that when the JFET junction is reverse-biased the Gate current is practically zero, whereas the Base current of the BJT is always some value greater than zero.
Bias arrangement for an N-channel JFET and corresponding circuit symbols.

The cross sectional diagram above shows an N-type semiconductor channel with a P-type region called the Gate diffused into the N-type channel forming a reverse biased PN-junction and it is this junction which forms the depletion region around the Gate area when no external voltages are applied. JFETs are therefore known as depletion mode devices. This depletion region produces a potential gradient which is of varying thickness around the PN-junction and restrict the current flow through the channel by reducing its effective width and thus increasing the overall resistance of the channel itself. The most-depleted portion of the depletion region is in between the Gate and the Drain, while the least-depleted area is between the Gate and the Source. Then the JFETs channel conducts with zero bias voltage applied (i.e. the depletion region has near zero width).

With no external Gate voltage \( V_G = 0 \), and a small voltage \( V_{DS} \) applied between the Drain and the Source, maximum saturation current \( I_{DSS} \) will flow through the channel from the Drain to the Source restricted only by the small depletion region around the junctions.

If a small negative voltage \( -V_{GS} \) is now applied to the Gate the size of the depletion region begins to increase reducing the overall effective area of the channel and thus reducing the current flowing through it, a sort of "squeezing" effect takes place. So by applying a reverse bias voltage increases the width of the depletion region which in turn reduces the conduction of the channel. Since the PN-junction is reverse biased, little current will flow into the gate connection. As the Gate voltage \( -V_{GS} \) is made more negative, the width of the channel decreases until no more current flows between the Drain and the Source and the FET is said to be "pinched-off" (similar
to the cut-off region for a BJT). The voltage at which the channel closes is called the "pinch-off voltage", \( V_P \).

**JFET Channel Pinched-off**

![JFET Channel Pinched-off Diagram]

In this pinch-off region the Gate voltage, \( V_{GS} \) controls the channel current and \( V_{DS} \) has little or no effect.

**JFET Model**

The result is that the FET acts more like a voltage controlled resistor which has zero resistance when \( V_{GS} = 0 \) and maximum "ON" resistance \( R_{DS} \) when the Gate voltage is very negative. Under normal operating conditions, the JFET gate is always negatively biased relative to the source.

It is essential that the Gate voltage is never positive since if it is all the channel current will flow to the Gate and not to the Source, the result is damage to the JFET. Then to close the channel:

- No Gate voltage \( V_{GS} \) and \( V_{DS} \) is increased from zero.
- No \( V_{DS} \) and Gate control is decreased negatively from zero.
- \( V_{DS} \) and \( V_{GS} \) varying.

The P-channel Junction Field Effect Transistor operates the same as the N-channel above, with the following exceptions: 1). Channel current is positive due to holes, 2). The polarity of the biasing voltage needs to be reversed.

The output characteristics of an N-channel JFET with the gate short-circuited to the source is given as
Output characteristic V-I curves of a typical junction FET

The voltage $V_{GS}$ applied to the Gate controls the current flowing between the Drain and the Source terminals. $V_{GS}$ refers to the voltage applied between the Gate and the Source while $V_{DS}$ refers to the voltage applied between the Drain and the Source. Because a Junction Field Effect Transistor is a voltage controlled device, "NO current flows into the gate!" then the Source current ($I_S$) flowing out of the device equals the Drain current flowing into it and therefore ($I_D = I_S$).

The characteristics curves example shown above, shows the four different regions of operation for a JFET and these are given as:

- **Ohmic Region** - When $V_{GS} = 0$ the depletion layer of the channel is very small and the JFET acts like a voltage controlled resistor.
- **Cut-off Region** - This is also known as the pinch-off region were the Gate voltage, $V_{GS}$ is sufficient to cause the JFET to act as an open circuit as the channel resistance is at maximum.
• Saturation or Active Region - The JFET becomes a good conductor and is controlled by the Gate-Source voltage, \( V_{GS} \) while the Drain-Source voltage, \( V_{DS} \) has little or no effect.

• Breakdown Region - The voltage between the Drain and the Source, \( V_{DS} \) is high enough to causes the JFET’s resistive channel to break down and pass uncontrolled maximum current.

The characteristics curves for a P-channel junction field effect transistor are the same as those above, except that the Drain current \( I_D \) decreases with an increasing positive Gate-Source voltage, \( V_{GS} \).

The Drain current is zero when \( V_{GS} = V_P \). For normal operation, \( V_{GS} \) is biased to be somewhere between \( V_P \) and 0. Then we can calculate the Drain current, \( I_D \) for any given bias point in the saturation or active region as follows:

\[
I_D = I_{DSS} \left[ 1 - \frac{V_{GS}}{V_P} \right]^2
\]

Note that the value of the Drain current will be between zero (pinch-off) and \( I_{DSS} \) (maximum current). By knowing the Drain current \( I_D \) and the Drain-Source voltage \( V_{DS} \) the resistance of the channel (\( I_D \)) is given as:

\[
R_{DS} = \frac{\Delta V_{DS}}{\Delta I_D} = \frac{1}{g_m}
\]

Where: \( g_m \) is the "transconductance gain" since the JFET is a voltage controlled device and which represents the rate of change of the Drain current with respect to the change in Gate-Source voltage.

2.5 JFET Amplifier

Just like the bipolar junction transistor, JFET’s can be used to make single stage class A amplifier circuits with the JFET common source amplifier and characteristics being very similar to the BJT common emitter circuit. The main advantage JFET amplifiers have over BJT amplifiers is their high input impedance which is controlled by the Gate biasing resistive network formed by \( R_1 \) and \( R_2 \) as shown.
Biasing of JFET Amplifier

This common source (CS) amplifier circuit is biased in class A mode by the voltage divider network formed by $R_1$ and $R_2$. The voltage across the Source resistor $R_S$ is generally set at one fourth $V_{DD}$, i.e., $\frac{V_{DD}}{4}$. The required Gate voltage can then be calculated using this $R_S$ value. Since the Gate current is zero, $I_G = 0$, we can set the required DC quiescent voltage by the proper selection of resistors $R_1$ and $R_2$.

The control of the Drain current by a negative Gate potential makes the Junction Field Effect Transistor useful as a switch and it is essential that the Gate voltage is never positive for an N-channel JFET as the channel current will flow to the Gate and not the Drain resulting in damage to the JFET. The principals of operation for a P-channel JFET are the same as for the N-channel JFET, except that the polarity of the voltages need to be reversed.

In the next tutorial about Transistors, we will look at another type of Field Effect Transistor called a MOSFET whose Gate connection is completely isolated from the main current carrying channel.

2.6 CMOS

Complementary Metal-Oxide-Silicon circuits require an nMOS and pMOS transistor technology on the same substrate. To this end, an n-type well is provided in the p-type substrate. Alternatively one can use a p-well or both an n-type and p-type well in a low-doped substrate. The gate oxide, poly-silicon gate and source-drain contact metal are typically shared between the pMOS and nMOS technology, while the source-drain implants must be done separately. Since CMOS circuits contain pMOS devices, which are affected by the lower hole mobility, CMOS circuits are not faster than their all-nMOS counter parts. Even when scaling the size of the pMOS devices so that they provide the same current, the larger pMOS device has a higher capacitance.

The CMOS advantage is that the output of a CMOS inverter can be as high as the power supply voltage and as low as ground. This large voltage swing and the steep transition between logic levels yield large operation margins and therefore also a high circuit yield. In addition, there is no power dissipation in either logic state. Instead the power dissipation occurs only when a
transition is made between logic states. CMOS circuits are therefore not faster than nMOS circuits but are more suited for very/ultra large-scale integration (VLSI/ULSI).

CMOS circuits have one property, which is very undesirable, namely latchup. Latchup occurs when four alternating p-type and n-type regions are brought in close proximity. Together they form two bipolar transistors, one npn and one pnp transistor. The base of each transistor is connected to the collector of the other, forming a cross-coupled thyristor-like combination. As a current is applied to the base of one transistor, the current is amplified by the transistor and provided as the base current of the other one. If the product of the current gain of both transistors is larger than unity, the current through both devices increases until the series resistances of the circuit limits the current. Latchup therefore results in excessive power dissipation and faulty logic levels in the gates affected. In principle, this effect can be eliminated by separating the n-type and p-type device. A more effective and less space-consuming solution is the use of trenches, which block the minority carrier flow. A deep and narrow trench is etched between all n-type and p-type wells, passivated and refilled with an insulating layer.

**VMOS Transistors and UMOS**

The VMOS transistor, named after the V-shaped groove, is a vertical MOSFET with high current handling capability as well as high blocking voltage. It consists of a double diffused n+ p layer, which is cut by a V-shaped groove. The V-groove is easily fabricated by anisotropically etching a (100) silicon surface using a concentrated KOH solution. The V-groove is then coated with a gate oxide, followed by the gate electrode. As the V-groove cuts through the double diffused layer, it creates two vertical MOSFETs, one on each side of the groove. The combination of the V-groove with the double diffused layers results in a short gate length, which is determined by the thickness of the p-type layer. The vertical structure allows the use of a low-doped drain region, which results in a high blocking voltage.

![Insulated Gate Bipolar Transistor (IGBT): a) equivalent circuit and b) device cross-section.](image-url)
2.7 Recommended Questions

1. Draw and explain the base biased amplifier.
2. Draw a typical common emitter amplifier and explain the function of each component in it.
3. Explain the single stage CE amplifier.
4. Draw the dc and ac equivalent circuits for base bias, voltage divider and two supply emitter bias transistor circuits.
5. For the VBD amplifier shown below, draw the DC equivalent circuit and find the
6. DC quantities \( I_E \), \( V_E \), \( V_{CE} \) and \( V_C \). (Jul-2009)
7. Explain with a neat circuit diagram and a.c equivalent circuit, the working of base biased amplifier. (Jan-2010)
8. Sketch the output waveform for the clipper circuit of Fig.Q2 (b) shown below.
9. Assume silicon diode and obtain the peak magnitude of the output waveform. (Jan-2010)
10. Find the voltage gain and output voltage across the load resistor for the given circuit of FigQ2( C ) (Jan-2010)
11. Explain with a neat circuit diagram, voltage divider bias amplifier by mentioning the importance of bypass capacitor. (Jun-2010)
Unit – 3: Photodiodes

3.1 Photodiodes

Photodiodes are junction semiconductor light sensors that generate current or voltage when the PN junction in the semiconductor is illuminated by light of sufficient energy.

The spectral response of the photodiode is a function of the bandgap energy of the material used in its construction. The cut-off wavelength of the photodiode is given by

$$\lambda_c = \frac{1240}{E_g}$$

where $\lambda_c$ is the cut-off wavelength (nm) and $E_g$ is the bandgap energy (eV).

Photodiodes are mostly constructed using silicon, germanium, indium gallium arsenide (InGaAs), lead sulphide (PbS) and mercury cadmium telluride (HgCdTe). Depending upon their construction there are several types of photodiodes. They are,

3.2 PN Photodiode

PN Photodiodes comprise a PN junction as shown in figure.

When light with sufficient energy strikes the diode, photo-induced carriers are generated which include electrons in the conduction band of the P-type material and holes in the valence band of the N-type material.
When photodiode is reverse biased, the photo-induced electrons will move down the potential hill from P-side to the N-side. Similarly, the photo-induced holes will add to the current flow by moving across the junction to the P-side from the N-side.

PN Photodiodes are used for precision photometry applications like medical instrumentation, analytical instruments, semiconductor tools and industrial measurement systems.

### 3.3 PIN Photodiode

In PIN photodiodes, an extra high resistance intrinsic layer is added between the P and N layers. This has the effect of reducing the transit or diffusion time of the photo-induced electron-hole pairs which in turn results in improved response time.

PIN photodiodes feature low capacitance, thereby offering high bandwidth making them suitable for high speed photometry and optical communication applications.

### 3.4 Schottky Photodiode

In Schottky-type Photodiodes, a thin gold coating is sputtered onto the N-material to forma Schottky effect PN junction. Schottky Photodiodes have enhanced ultraviolet (UV) response.

**Avalanche Photodiode (APD)**

APDs are high-speed, high-sensitivity photodiodes utilizing an internal gain mechanism that functions by applying a relatively higher reverse bias voltage than that is applied in the case of PIN photodiodes.

APDs are so constructed to provide a very uniform junction that exhibits the avalanche effect at reverse-bias voltages between 30V - 200V. The electron-hole pairs that are generated by incident photons are accelerated by the high electric field to force the new electrons to move from the valance band to the conduction band.
Avalanche Photodiode

They offer excellent signal-to-noise ratio. Hence, they are used in variety of applications requiring high sensitivity such as long distance optical communication and measurement.

**VI Characteristics of Photodiode**

The VI characteristics of photodiode is as shown in figure below.
VI Characteristics of Photodiode

VI characteristics of photodiode is similar that of a conventional diode, but when light strikes, curve shifts downwards with increasing intensity of light.

If the photodiode terminals are shorted, a photocurrent proportional to the light intensity will flow in a direction from anode to cathode. If the circuit is open, then an open circuit voltage will be generated with the positive polarity at the anode.

It is mentioned that short circuit current is linearly proportional to the light intensity while open circuit voltage has a logarithmic relationship with the light intensity.

Photodiodes can be operated in two modes namely the Photovoltaic mode and Photoconductive mode. In Photovoltaic mode of operation, no bias voltage is applied and due to the incident light, a forward voltage is produced across the photodiode. In Photoconductive operational mode, a reverse bias voltage is applied across the photodiode; this widens the depletion region resulting in higher speed of response.

In the photovoltaic mode, the photodiode is operated with zero external bias voltage and is generally used for low-speed applications or for detecting low light levels. The output voltage of photodiode circuits can be calculated by $I_{det} \times R$ and $I_{det} \times R_f$, Where $I_{det}$ is the current through the photodiode.

Solar Cells

Solar cell is device whose operation is very similar to that of a photodiode operating in the photovoltaic mode. The operating principle of solar cell is based on the photovoltaic effect.

When the PN junction of solar cell is exposed to sun light, open circuit voltage is generated. This open circuit voltage leads to the flow of electric current through a load resistor connected across it.

The incident photon energy leads to generation of electron-hole pairs. The electron-hole pairs either recombine and vanish or start drifting in the opposite directions with electrons moving towards the N-
region and holes moving towards the P-region. This accumulation of positive and negative charge carriers constitutes the open circuit voltage.

This voltage can cause a current to flow through an external load or when the junction is shorted, the result is a short circuit current whose magnitude is proportional to input light intensity. As the energy produced by the individual solar cell is very less (500mV output with a load current capability of 150mA), series-parallel arrangement of solar cells is done to get the desired output. The series combination is used to enhance the output voltage while the parallel combination is used to enhance the current gain.

### 3.5 Phototransistors

![Phototransistor Diagram]

The above figure shows the construction of Phototransistor. Phototransistors are usually connected in the common-emitter configuration with the base open and the radiation is concentrated on the region near the collector-base junction.
When there is no radiation incident on the phototransistor, the collector current is due to the thermally generated carriers, called as dark current and is given by

\[ I_c = (\beta+1)I_{co} \]

Where \( I_{co} \) is the reverse saturation current

When light is incident on the phototransistor, photocurrent is generated and the magnitude of the collector current increases. The expression for the collector current is given by

\[ I_c = (\beta+1) (I_{co}+I_{\lambda}) \]

Where \( I_{\lambda} \) is the current generated due to incident light photons.

**Phototransistor Applications**

Phototransistor can be used in two configurations, namely, the common-emitter configuration and the common-collector configuration.

In the common-emitter configuration, the output is high and goes low when light is incident on the phototransistor, whereas in common-collector configuration, the output is high and goes to low when light is incident on the phototransistor.

3.6 Light-Emitting Diodes (LED)

LED is a semiconductor PN junction diode designed to emit light when forward-biased. It is one of the most popular optoelectronic source. LEDs consume very little power and are inexpensive.
When PN junction is forward biased, the electrons in the N-type material and the holes in the P-type material travel towards the junction. Some of these holes and electrons recombine with each other and in the process radiate energy. The energy will be released either in the form of photons of light. Gallium Phosphide (GaP), Gallium Arsenide (GaAs) and Gallium arsenide Phosphide (GaAsP) are used in the construction of LEDs.

In the absence of an externally applied voltage, the N-type material contains electrons while the P-type material contains holes that can act as current carriers. When the diode is forward-biased, the energy levels shift and there is significant increase in the concentration of electrons in the conduction band on the N-side and that of holes in valance band on the P-side. The electrons and holes combine near the junction to release energy in the form of photons. The process of light emission in LED is spontaneous, i.e., the photons emitted are not in phase and travel in different directions.

The energy of the photon resulting from this recombination is equal to the bandgap energy of the semiconductor material and is expressed by:
The wavelength (nm) and bandgap energy (eV).

**LED Characteristic Curve**

As the LED is operated in the forward-biased mode, the VI characteristics in the forward-biased region are shown. VI characteristics of LEDs are similar to that of conventional PN junction diodes except that the cut-in voltage in the case of LEDs is in the range of 1.3-3V as compared to 0.7V for silicon diodes and 0.3V for germanium diodes.

**LED Parameters**

1. **Forward Voltage (V>F):** It is the DC voltage across the LED when it is ON.
2. **Candle Power (CP):** It is a measure of the luminous intensity or the brightness of the light emitted by the LED. It is a non-linear function of LED current and the value of CP increases with increase in the current flowing through the LED.
3. **Radiant Power Output (P_o):** It is the light power of the LED.
4. **Peak Spectral Emission (λ_P):** It is the wavelength where the intensity of light emitted by the LED is maximum.
5. **Spectral Bandwidth:** It is the measure of concentration of color brightness around the LEDs nominal wavelength.

**LED Drive Circuit**
Connecting LEDs in Parallel

Connecting LEDs in Series

LEDs are operated in the forward-biased mode. As the current through the LED changes very rapidly with change in forward voltage above the threshold voltage, LEDs are current-driven devices. The resistor (R) is used limit the current flowing through the device. A silicon diode can be placed inversely parallel to the LED for reverse polarity voltage protection.

The current that will flow through the LED is given by

\[ I = \frac{(V_{CC} - V_F)}{R} \]

The value of the resistor (R) to be connected is given by

\[ R = \left[ V_{CC} - (V_{F1} + V_{F2} + V_{F3} + \ldots + V_{Fn}) \right] \left[ \frac{1}{I} \right] \]

Liquid Crystal Displays (LCD)

Liquid Crystals are materials that exhibit properties of both solids and liquids, that is, they are an intermediate phase of matter. They can be classified into three different groups: nematic, smectic and cholestric. Nematic liquid crystals are generally used in the fabrication of liquid crystal displays (LCDs) with the twisted nematic material being the most common.

Construction of an LCD

Construction of LCD Display

Operation of LCD Display
An LCD display consists of liquid-crystal fluid, conductive electrodes, a set of polarizers and a glass casing. The outermost layers are the polarizers which are housed on the outer surface of the glass casing. The polarizer attached to the front glass is referred to as the front polarizer, while the one attached to the rear glass is the rear polarizer.

On the inner surface of the glass casing, transparent electrodes are placed in the shape of desired image. The electrode attached to the front glass is referred to as the segment electrode while the one attached to the rear glass is the backplane or the common electrode. The liquid crystal is sandwiched between the two electrodes.

**Operation**

The basic principle of operation of LCD is to control the transmission of light by changing the polarization of the light passing through the liquid crystal with the help of an externally applied voltage. As LCDs do not emit their own light, backlighting is used to enhance the legibility of the display in dark conditions.

LCDs have the capability to produce both positive as well as negative images. A positive image is defined as a dark image on a light background. In a positive image display, the front and rear polarizers are perpendicular to each other. Light entering the display is guided by the orientation of the liquid crystal molecules that are twisted by 90° from the front glass plate to the rear glass plate. This twist allows the incoming light to pass through the second polarizer. When a light is applied to the display, the liquid crystal molecules straighten out and stop redirecting the light. As a result, light travels straight through and is filtered out by the second polarizer. Therefore, no light can pass through, making this region darker compared to the rest of the screen. Hence, in order to display characters or graphics, voltage is applied to the desired regions, making them dark and visible to the eye.

A negative image is a light image on a dark background. In negative image displays, the front and the rear polarizers are aligned parallel to each other.

**Driving an LCD**

LCD can be classified as direct-drive and multiplex-drive displays depending upon the technique used to drive them. Direct-drive displays, also known as static-drive displays, have an independent driver for each pixel. The voltage in this case is a square waveform having two voltage levels, namely, ground and Vcc. As the display size increases, the drive circuitry becomes very complex. Hence, multiplex drive circuits are used for larger size displays. These displays reduce the total number of interconnections between the LCD and the driver. They have more than one backplane and the driver produces an amplitude-varying, time-synchronized waveform for both segment and backplanes.

**LCD Response Time**

The LCD response time is defined by the ON and OFF response times.
ON time refers to the time required by an OFF pixel to become visible after the application of proper drive voltage. The OFF time is defined as the time required by the ON pixel to turn OFF after the application of proper drive voltage.

The response time of LCDs varies widely with temperature and increase rapidly at low operating temperatures.

**Liquid Crystal Display Types**

LCDs are non-emissive devices, that is, they do not generate light on their own. Depending upon the mode of transmission of light in an LCD, they are classified:

**Reflective LCD displays:**

Reflective LCD displays have a reflector attached to the rear polarizer which reflects incoming light evenly back into the display. These displays rely on the ambient light to operate and do not work in the dark conditions. They produce only positive images. The front and the rear polarizers are perpendicular to each other. These types of displays are commonly used in calculators and digital wrist watches.

**Transmissive LCD displays:**

In transmissive LCD displays, back light is used as the light source. Most of these displays operate in the negative mode, that is, the text will be displayed in light color and the background is dark color. Transmissive displays are good for very low light level conditions. They offer very poor contrast when used in direct sunlight and provide good picture quality indoors. They are generally used in medical devices, electronic test and measuring equipments and in laptops.

**Transreflective LCD displays**
Transreflective displays are a combination of reflective and transmissive displays. A white or silver translucent material is applied to the rear of the display, which reflects some of the ambient light back to the observer. It also allows the backlight to pass through. They are good for displays operating in varying light conditions.

**Active and Passive LCD displays**

LCD displays are classified as passive LCD displays and active LCD displays depending upon the nature of the activation circuit.

Passive displays use components that do not supply their own energy to turn ON or OFF the desired pixel.

Active displays use an active device such as a transistor or diode in each which acts like a switch that precisely controls the voltage that each pixel receives.

**Advantages and Disadvantages**

- LCD displays are not active sources of light
- They offer very low power consumption, low operating voltages and good flexibility
- Their response time is too slow for many applications
- They offer limited viewing angle and are temperature sensitive

**3.7 Cathode Ray Tube Displays**

Cathode Ray Tube (CRT) displays are used in a wide range of systems ranging from consumer electronic systems like television and computer monitors to measuring instruments like oscilloscopes to military systems like radar and so on. CRT display is a specialized vacuum tube in which the images are produced when the electron beam strikes the fluorescent screen. CRT displays can be monochrome displays as well as colored displays.
Monochrome CRT displays comprise a single electron gun, a fluorescent screen and an internal or external mechanism to accelerate and deflect the electron beam. The electron gun produces a narrow beam of electrons that are accelerated by the anodes. There are two sets of deflecting coils, namely, the horizontal coil and the vertical coil. These coils produce an extremely low frequency electromagnetic field in the horizontal and vertical directions to adjust the direction of the electron beam. CRT tubes also have a mechanism to vary the intensity of the electron beam. In order to produce moving pictures in natural colors on the screen, complex signals are applied to the deflecting coils and to the circuitry responsible for controlling the intensity of the electron beam. This results in movement of the spot from right to left and from top to bottom of the screen. The speed of the spot movement is so fast that the person viewing the screen sees a constant image on the entire screen.

Color CRT displays comprises three electron guns, one each for the three primary colors namely red, blue and green. The CRT produces three overlapping images, one in red, one in green and one in blue. This is referred to as the RGB color model.

Advantages and Disadvantages

- CRT displays offer very high resolution and these displays emit their own light; therefore, they have very high value of peak luminance.
- They offer wide viewing angles of the order of 180°.
- CRT displays are bulky and consume significant power.
- They require high voltages to operate and they cause fatigue and strain to the human eye.

3.8 Emerging Display Technologies

**Organic Light-Emitting Diodes (OLEDs)**

OLEDs are composed of a light-emitting organic material sandwiched between two conducting plates, one of N-type material and the other of P-type material. When an electric potential is applied between these plates, holes are ejected from the P-type plate and electrons are ejected from the N-type plate. Recombination of these holes and electrons, energy is released in the form of light photons. The wavelength of light emitted depends upon the bandgap energy of the semiconductor material used. OLEDs can be classified into three types, namely, small molecule OLEDs (SMOLED), polymer LEDs (PLED) and dendrimer OLEDs.

**Digital Light Processing Technology (DLP)**

DLP technology makes use of an optical semiconductor device referred to as digital micromirror device (DMD) which is basically a precise light switch that can digitally modulate light through a large number of microscopic mirrors arranged in a rectangular array. These mirrors are mounted on tiny hinges and can be tilted away or towards the light source with the help of DMD chip and thus projecting a light or a dark pixel on the screen. Use of DLP technology is currently limited to large projection systems.
Plasma Display Panels (PDP)

Plasma displays are composed of millions of cells sandwiched between two panels of glass. Two electrodes, namely, the address electrodes and display electrodes, are also placed between the two glass plates covering the entire screen. The address electrodes are printed on the rear glass plate and the transparent display electrodes are located above the cells along the front glass plate. These electrodes are perpendicular to each other forming a grid network. Each cell is filled with xenon and neon gas mixture. The electrodes intersecting at a specific cell are charged to excite the gas mixture in that cell. When the gas mixture is excited, plasma is created releasing ultraviolet light which then excites the phosphor electrons located on the sides of the cells. These electrons in turn release visible light and return to their lower energy state. Each pixel is composed of three cells containing red, green and blue phosphors.

Plasma displays offer advantages like each pixel generates its own light offering large viewing angles, generates super image quality and the image quality is not affected by the area of the display, but these displays are fragile in nature and are susceptible to burn-out from static images.

Field Emission Displays

FEDs function like CRT displays with the main difference being that these displays use millions of small electron guns to emit electrons at the screen instead of just one as in case of CRT. FED displays produce the same quality of image as produced by the CRT displays without being bulky as the CRT displays. These displays can be as thin as LCD and as large as Plasma.

Electronic Ink Displays

Electronic ink displays, also referred to as electronic paper, are active matrix displays making use of pigments that resemble the ink used in print.

3.9 Optocouplers

An optocoupler, also referred to as an optoisolator, is a device that uses a short optical transmission path to transfer signals between the elements of a circuit. Optocouplers are sealed units that house an optical transmitting device and a photosensitive device that are coupled together optically. The optical path may be air or a dielectric waveguide.

![Optical Path Diagram](image-url)
Optocoupler Parameters

The important parameters that define the performance of an optocoupler are:

1. Forward Optocoupling Efficiency: It is specified in terms of current transfer ratio (CTR). CTR is the ratio of the output current to the input current.
2. Isolation Voltage: It is the maximum permissible DC potential that can be allowed to exist between the input and the output circuits.
3. Input Current: It is the maximum permissible forward current that is allowed to flow into the transmitting LED.
4. \( V_{CE\text{max}} \): It is the transistor’s maximum collector-emitter voltage rating. It limits the supply voltage that can be applied to the output circuit.
5. Bandwidth: It determines the maximum signal frequency that can be successfully passed through the optocouplers. The bandwidth of an optocoupler depends upon its switching speed.

Optocoupler Configurations

Non-Latching Optocoupler Configuration:

- Photodiode
- Phototransistor
- Photo-Darlington Transistors
- Photoconductor
- PhotoFET

Latching Optocoupler Configuration:

- PhotoSCR
- PhotoDIAC
- PhotoTRIAC
Advantages

- Complete electrical isolation between input and output circuit
- Less response time of optocouplers enable data transmission in MHz range
- Capable of wideband signal transmission
- Unidirectional signal transfer, output does not loop back to the input circuit
- Easy interface with logic devices
- Compact and light weight
- Noise, transients, contact bounce etc. are completely eliminated

3.10 Recommended Questions

1. Derive the equation for $A_v$, $A_i$, $Z_i$, and $Z_o$ for a common emitter amplifier.
2. Derive the expression for voltage gain for two stage CE amplifier.
3. Draw and explain the working of swamped amplifier.
4. What are the advantages of swamped amplifier?
5. Draw and explain the working of two-stage feedback amplifier.
6. Derive the equation for $A_v$, $A_i$, $Z_i$, and $Z_o$ for the CC amplifier.
7. Derive the equation for $A_v$, $A_i$, $Z_i$, and $Z_o$ for CB amplifier.
8. Draw and explain the two transistor regulator.
9. Draw and explain the Zener follower voltage regulator.
10. Draw and explain the cascaded CE and CC stages of an amplifier.
11. Derive the expression for voltage gain of the CE-CC cascaded amplifier.
12. Explain how h-parameter can be obtained from the transistor characteristics.
13. The transistor amplifier shown in the figure below uses a transistor whose h-parameters are $h_{ie} = 1.1k\Omega$, $h_{re} = 50$, $h_{ce} = 2.5 \times 10^{-4}$ and $1/h_{oe} = 40k\Omega$. Calculate $A_i = I_o/I_i$, $A_v$, $A_{vs}$, $R_o$ and $R_i$. (Aug-2005)
14. Draw the ‘h’ parameter equivalent circuit for a typical common emitter amplifier and Derive the expression for $A_i$, $R_i$, and $A_v$.
15. For the amplifier shown in the figure below calculate $R_i$, $R'_i$, $A_v$, $A_{vs}$ and $A_i = -(I_2/I_1)$. The transistor parameters are $h_{ie} = 1.1k\Omega$, $h_{re} = 2.5 \times 10^{-4}$, $h_{ce} = 50$ and $h_{oe} = 25\mu A/V$. (Aug-2004, 2005)
16. Explain how h-parameters can be obtained from the static characteristics of a transistor. (Jan-2007)
UNIT – 4: Small Signal Analysis of Amplifiers

4.1 Basic FET Amplifiers

In the last chapter, we described the operation of the FET, in particular the MOSFET, and analyzed and designed the dc response of circuits containing these devices. In this chapter, we emphasize the use of FETs in linear amplifier applications. Although a major use of MOSFETs is in digital applications, they are also used in linear amplifier circuits.

There are three basic configurations of single-stage or single-transistor FET amplifiers. These are the common-source, source-follower, and common-gate configurations. We investigate the characteristics of each configuration and show how these properties are used in various applications. Since MOSFET integrated circuit amplifiers normally use MOSFETs as load devices instead of resistors because of their small size, we introduce the technique of using MOSFET enhancement or depletion devices as loads. These three configurations form the building blocks for more complex amplifiers, so gaining a good understanding of these three amplifier circuits is an important goal of this chapter.

In integrated circuit systems, amplifiers are usually connected in series or cascade, forming a multistage configuration, to increase the overall voltage gain or to provide a particular combination of voltage gain and output resistance. We consider a few of the many possible multistage configurations, to introduce the analysis methods required for such circuits, as well as their properties.

4.2 THE MOSFET AMPLIFIER

We discussed the reasons linear amplifiers are necessary in analog electronic systems. In this chapter, we continue the analysis and design of linear amplifiers that use field-effect transistors as the amplifying device. The term small signal means that we can linearize the ac equivalent circuit. We will define what is meant by small signal in the case of MOSFET circuits. The term linear amplifiers means that we can use superposition so that the dc analysis and ac analysis of the circuits can be performed separately and the total response is the sum of the two individual responses.

The mechanism with which MOSFET circuits amplify small time-varying signals was introduced in the last chapter. In this section, we will expand that discussion using the graphical technique, dc load line, and ac load line. In the process, we will develop the various small-signal parameters of linear circuits and the corresponding equivalent circuits.

There are four possible equivalent circuits that can be used.
The most common equivalent circuit that is used for the FET amplifiers is the transconductance amplifier, in which the input signal is a voltage and the output signal is a current.

**Graphical Analysis, Load Lines, and Small-Signal Parameters**

Figure 6.1 shows an NMOS common-source circuit with a time-varying voltage source in series with the dc source. We assume the time-varying input signal is sinusoidal. Figure 6.2 shows the transistor characteristics, dc load line, and Q-point, where the dc load line and Q-point are functions of $v_{GS}$, $V_{DD}$, $R_D$ and the transistor parameters.
For the output voltage to be a linear function of the input voltage, the transistor must be biased in the saturation region. Note that, although we primarily use n-channel, enhancement-mode MOSFETs in our discussions, the same results apply to the other MOSFETs.

Also shown in Figure 6.2 are the sinusoidal variations in the gate-to-source voltage, drain current, and drain-to-source voltage, as a result of the sinusoidal source voltage. The total gate-to-source voltage is the sum of \( V_{GSQ} \) and \( v_i \). As \( v_i \) increases, the instantaneous value of \( v_{GS} \) increases, and the bias point moves up the load line. A larger value of \( v_{GS} \) means a larger drain current and a smaller value of \( v_{DS} \). Once the Q-point is established, we can develop a mathematical model for the sinusoidal, or small-signal, variations in the gate-to-source voltage, drain-to-source voltage, and drain current.

The time-varying signal source in Figure 6.1 generates a time-varying component of the gate-to-source voltage. For the FET to operate as a linear amplifier, the transistor must be biased in the saturation region, and the instantaneous drain current and drain-to-source voltage must also be confined to the saturation region.
Transistor Parameters

The instantaneous gate-to-source voltage is

\[ v_{GS} = V_{GSQ} + v_i = V_{GSQ} + v_{gs} \]  \hspace{1cm} (6.1)

where \( V_{GSQ} \) is the dc component and \( v_{gs} \) is the ac component. The instantaneous drain current is

\[ i_D = K_n(v_{gs} - V_{TN})^2 \]  \hspace{1cm} (6.2)

Substituting Equation (6.1) into (6.2) produces

\[ i_D = K_n[V_{GSQ} + v_{gs} - V_{TN}]^2 = K_n(V_{GSQ} - V_{TN} + v_{gs})^2 \]  \hspace{1cm} (6.3(a))

or

\[ i_D = K_n(V_{GSQ} - V_{TN})^2 + 2K_n(V_{GSQ} - V_{TN})v_{gs} + K_nv_{gs}^2 \]  \hspace{1cm} (6.3(b))

The first term in Equation (6.3(b)) is the dc or quiescent drain current \( I_{DQ} \), the second term is the time-varying drain current component that is linearly related to the signal \( v_{gs} \), and the third term is proportional to the square of the signal voltage. For a sinusoidal input signal, the squared term produces undesirable harmonics, or nonlinear distortion, in the output voltage. To minimize these harmonics, we require

\[ v_{gs} \ll 2(V_{GSQ} - V_{TN}) \]  \hspace{1cm} (6.4)

which means that the third term in Equation (6.3(b)) will be much smaller than the second term. Equation (6.4) represents the small-signal condition that must be satisfied for linear amplifiers.

Neglecting the \( v_{gs}^2 \) term, we can write Equation (6.3(b))

\[ i_D = I_{DQ} + i_d \]  \hspace{1cm} (6.5)

Again, small-signal implies linearity so that the total current can be separated into a dc component and an ac component. The ac component of the drain current is given by

\[ i_d = 2K_n(V_{GSQ} - V_{TN})v_{gs} \]  \hspace{1cm} (6.6)

The small-signal drain current is related to the small-signal gate-to-source voltage by the transconductance \( g_m \). The relationship is
\[ g_m = \frac{i_d}{v_{gs}} = 2K_n(V_{GSQ} - V_{TN}) \]  

(6.7)

The transconductance is a transfer coefficient relating output current to input voltage and can be thought of as representing the gain of the transistor.

The transconductance can also be obtained from the derivative

\[ g_m = \frac{\partial i_d}{\partial v_{GS}} \bigg|_{v_{gs}=v_{GSQ}=\text{const.}} = 2K_n(V_{GSQ} - V_{TN}) \]  

(6.8(a))

which can be written

\[ g_m = 2\sqrt{K_n I_{DQ}} \]  

(6.8(b))

The drain current versus gate-to-source voltage for the transistor biased in the saturation region is given in Equation (6.2) and is shown in Figure 6.3. The transconductance \( g_m \) is the slope of the curve. If the time-varying signal \( v_{gs} \) is sufficiently small, the transconductance \( g_m \) is a constant. With the \( Q \)-point in the saturation region, the transistor operates as a current source that is linearly controlled by \( v_{gs} \). If the \( Q \)-point moves into the nonsaturation region, the transistor no longer operates as a linearly controlled current source.

![Figure 6.3](image)

**Figure 6.3** Drain current versus gate-to-source voltage characteristics, with superimposed sinusoidal signals

As shown in Equation (6.8(a)), the transconductance is directly proportional to the conduction parameter \( K_n \), which in turn is a function of the width-to-length ratio. Therefore, increasing the width of the transistor increases the transconductance, or gain, of the transistor.
Comment: The transconductance of a bipolar transistor is \( g_m = \frac{I_C}{V_T} \), which is 38.5 mA/V for a collector current of 1 mA. The transconductance values of MOSFETs tend to be small compared to those of BJTs. However, the advantages of MOSFETs include high input impedance, small size, and low power dissipation.

**AC Equivalent Circuit**

From Figure 6.1, we see that the output voltage is

\[
v_{DS} = v_O = V_{DD} - i_D R_D
\]  

(6.9)

Using Equation (6.5), we obtain

\[
v_O = V_{DD} - (I_{DQ} + i_d) R_D = (V_{DD} - I_{DQ} R_D) - i_d R_D
\]  

(6.10)

The output voltage is also a combination of dc and ac values. The time-varying output signal is the time-varying drain-to-source voltage, or

\[
v_o = v_{ds} = -i_d R_D
\]  

(6.11)

Also, from Equations (6.6) and (6.7), we have

\[
i_d = g_m v_{gs}
\]  

(6.12)

In summary, the following relationships exist between the time-varying signals for the circuit in Figure 6.1. The equations are given in terms of the instantaneous ac values, as well as the phasors. We have

\[
v_{gs} = v_i
\]  

(6.13(a))

or

\[
V_{gs} = V_i
\]  

(6.13(b))

and

\[
i_d = g_m v_{gs}
\]  

(6.14(a))

or

\[
I_d = g_m V_{gs}
\]  

(6.14(b))

Also,

\[
v_{ds} = -i_d R_D
\]  

(6.15(a))

or

\[
V_{ds} = -I_d R_D
\]  

(6.15(b))

The ac equivalent circuit in Figure 6.4 is developed by setting the dc sources in Figure 6.1 equal to zero. The small-signal relationships are given in Equations (6.13), (6.14), and (6.15). As shown in Figure 6.1, the drain current, which is composed of ac signals superimposed on the quiescent value, flows through the voltage source \( V_{DD} \). Since the voltage across this
source is assumed to be constant, the sinusoidal current produces no sinusoidal voltage component across this element. The equivalent ac impedance is therefore zero, or a short circuit. Consequently, in the ac equivalent circuit, the dc voltage sources are equal to zero. We say that the node connecting $R_D$ and $V_{DD}$ is at signal ground.

### 4.3 Small-Signal Equivalent Circuit

Now that we have the ac equivalent circuit for the NMOS amplifier circuit, (Figure 6.4), we must develop a small-signal equivalent circuit for the transistor.

Initially, we assume that the signal frequency is sufficiently low so that any capacitance at the gate terminal can be neglected. The input to the gate thus appears as an open circuit, or an infinite resistance. Eq. 6.14 relates the small-signal drain current to the small-signal input voltage and Eq. 6.7 shows that the transconductance is a function of the Q-point. The resulting simplified small-signal equivalent circuit for the NMOS device is shown in Figure 6.5. (The phasor components are in parentheses.)

**Figure 6.5** (a) Common-source NMOS transistor with small-signal parameters and (b) simplified small-signal equivalent circuit for NMOS transistor
This small-signal equivalent circuit can also be expanded to take into account the finite output resistance of a MOSFET biased in the saturation region. This effect, discussed in the previous chapter, is a result of the nonzero slope in the $i_D$ versus $v_{DS}$ curve. We know that

$$i_D = K_n[(v_{GS} - V_{TN})^2(1 + \lambda v_{DS})]$$

(6.16)

where $\lambda$ is the channel-length modulation parameter and is a positive quantity. The small-signal output resistance, as previously defined, is

$$r_o = \left(\frac{\partial i_D}{\partial v_{DS}}\right)^{-1} \bigg|_{v_{GS} = V_{GSQ} = \text{const.}}$$

(6.17)

or

$$r_o = [\lambda K_n(V_{GSQ} - V_{TN})^2]^{-1} \cong [\lambda I_DQ]^{-1}$$

(6.18)

This small-signal output resistance is also a function of the $Q$-point parameters.

The expanded small-signal equivalent circuit of the n-channel MOSFET is shown in Figure 6.6 in phasor notation.

![Figure 6.6](image1)

**Figure 6.6** Expanded small-signal equivalent circuit, including output resistance, for NMOS transistor.

![Figure 6.7](image2)

**Figure 6.7** Small-signal equivalent circuit of common-source circuit with NMOS transistor model.

We note that the small-signal equivalent circuit for the MOSFET circuit is very similar to that of the BJT circuits.
Example 6.2  Objective: Determine the small-signal voltage gain of a MOSFET circuit.

For the circuit in Figure 6.1, assume parameters are: \( V_{GSQ} = 2.12 \text{ V}, \) \( V_{DD} = 5 \text{ V}, \) and \( R_D = 2.5 \text{ k}\Omega. \) Assume transistor parameters are: \( V_{TN} = 1 \text{ V}, \) \( K_n = 0.80 \text{ mA/V}^2, \) and \( \lambda = 0.02 \text{ V}^{-1}. \) Assume the transistor is biased in the saturation region.

Solution: The quiescent values are

\[
I_{DQ} \cong K_n(V_{GSQ} - V_{TN})^2 = (0.8)(2.12 - 1)^2 = 1.0 \text{ mA}
\]

and

\[
V_{DSQ} = V_{DD} - I_{DQ}R_D = 5 - (1)(2.5) = 2.5 \text{ V}
\]

Therefore,

\[
V_{DSQ} = 2.5 \text{ V} > V_{DS}(\text{sat}) = V_{GS} - V_{TN} = 1.82 - 1 = 0.82 \text{ V}
\]

which means that the transistor is biased in the saturation region, as initially assumed, and as required for a linear amplifier. The transconductance is

\[
g_m = 2K_n(V_{GSQ} - V_{TN}) = 2(0.8)(2.12 - 1) = 1.79 \text{ mA/V}
\]

and the output resistance is

\[
r_o = [\lambda I_{DQ}]^{-1} = [(0.02)(1)]^{-1} = 50 \text{ k}\Omega
\]

From Figure 6.7, the output voltage is

\[
V_o = -g_mV_{gs}(r_o \parallel R_D)
\]

Since \( V_{gs} = V_i, \) the small-signal voltage gain is

\[
A_v = \frac{V_o}{V_i} = -g_m(r_o \parallel R_D) = -(1.79)(50 \parallel 2.5) = -4.26
\]

Comment: Because of the relatively low value of transconductance, MOSFET circuits tend to have a lower small-signal voltage gain than comparable bipolar circuits. Also, the small-signal voltage gain contains a minus sign, which means that the sinusoidal output voltage is 180 degrees out of phase with respect to the input sinusoidal signal.

4.4 Problem-Solving Technique: MOSFET AC Analysis

Since we are dealing with linear amplifiers, superposition applies, which means that we can perform the dc and ac analyses separately. The analysis of the MOSFET amplifier proceeds as follows:

1. Analyze the circuit with only the dc sources present. This solution is the dc or quiescent solution. The transistor must be biased in the saturation region in order to produce a linear amplifier.
2. Replace each element in the circuit with its small-signal model, which means replacing the transistor by its small-signal equivalent circuit.

3. Analyze the small-signal equivalent circuit, setting the dc source components equal to zero, to produce the response of the circuit to the time-varying input signals only.

The previous discussion was for an n-channel MOSFET amplifier. The same basic analysts and equivalent circuit also applies to the p-channel transistor. Figure 6.8(a) shows a circuit containing a p-channel MOSFET.

![Diagram of p-channel MOSFET amplifier](image)

Figure 6.8 (a) Common-source circuit with PMOS transistor and (b) corresponding ac equivalent circuit

Note that the power supply voltage is connected to the source. (The subscript DD can be used to indicate that the supply is connected to the drain terminal. Here, however, $V_{DD}$ is simply the usual notation for the power supply voltage in MOSFET circuits.) Also note the change in current directions and voltage polarities compared to the circuit containing the NMOS transistor. Figure 6.8(b) shows the ac equivalent circuit, with the dc voltage sources replaced.

The final small-signal equivalent circuit of the p-channel MOSFET amplifier is shown in Figure 6.10.

![Diagram of small-signal equivalent circuit](image)

Figure 6.10 Small-signal equivalent circuit of common-source amplifier with PMOS transistor model
We also note that the expression for the small-signal voltage gain of the p-channel MOSFET amplifier is exactly the same as that for the n-channel MOSFET amplifier. The negative sign indicates that a 180-degree phase reversal exists between the output and input signals, for both the PMOS and the NMOS circuit.

4.5 Basic Transistor Amplifier Configurations

As we have seen, the MOSFET is a three-terminal device (actually 4 counting the substrate). Three basic single-transistor amplifier configurations can be formed, depending on which of the three transistor terminals is used as signal ground. These three basic configurations are appropriately called common source, common drain (source follower), and common gate. These three circuit configurations correspond to the common-emitter, emitter-follower, and common-base configurations using BJTs.

The input and output resistance characteristics of amplifiers are important in determining loading effects. These parameters, as well as voltage gain, for the three basic MOSFET circuit configurations will be determined in the following sections.

THE COMMON-SOURCE AMPLIFIER

In this section, we consider the first of the three basic circuits; the common-source amplifier. We will analyze several basic common-source circuits, and will determine small-signal voltage gain and input and output impedances.

A Basic Common-Source Configuration

For the circuit shown in Figure 6.13, assume that the transistor is biased in the saturation region by resistors $R_1$ and $R_2$, and that the signal frequency is sufficiently large for the coupling capacitor to act essentially as a short circuit. The signal source is represented by a Thevenin equivalent circuit, in which the signal voltage source $v_i$, is in series with an equivalent source resistance $R_{Si}$. As we will see, $R_{Si}$ should be much less than the amplifier input resistance, $R_i = R_1 \parallel R_2$ in order to minimize loading effects.

![Figure 6.13](image_url) Common-source circuit with voltage divider biasing and coupling capacitor
Figure 6.14 shows the resulting small-signal equivalent circuit. The small signal variables, such as the input signal voltage $V_i$, are given in phasor form.

![Small-signal equivalent circuit](image)

**Figure 6.14** Small-signal equivalent circuit, assuming coupling capacitor acts as a short circuit

The output voltage is

$$V_o = -g_m V_{gs} (r_o \parallel R_D)$$  \hspace{1cm} (6.27)

The input gate-to-source voltage is

$$V_{gs} = \left( \frac{R_i}{R_i + R_{Si}} \right) \cdot V_i$$  \hspace{1cm} (6.28)

so the small-signal voltage gain is

$$A_v = \frac{V_o}{V_i} = -g_m (r_o \parallel R_D) \cdot \left( \frac{R_i}{R_i + R_{Si}} \right)$$  \hspace{1cm} (6.29)

We can also relate the ac drain current to the ac drain-to-source voltage, as $V_{ds} = -I_d(R_D)$.

Figure 6.15 shows the dc load line, the transition point, and the Q-point, which is in the saturation region. As previously stated, in order to provide the maximum symmetrical output voltage swing and keep the transistor biased in the saturation region, the Q-point must be near the middle of the saturation region. At the same time, the input signal must be small enough for the amplifier to remain linear.

![DC load line and transition point](image)

**Figure 6.15** DC load line and transition point separating saturation and nonsaturation regions
The input and output resistances of the amplifier can be determined from Figure 6.14. The input resistance to the amplifier is $R_i = R_1 \parallel R_2$. Since the low-frequency input resistance looking into the gate of the MOSFET is essentially infinite, the input resistance is only a function of the bias resistors. The output resistance looking back into the output terminals is found by setting the independent input source $V_i$ equal to zero, which means that $V_{GS} = 0$. The output resistance is therefore $R_o = R_D \parallel r_o$.

**Example 6.3 Objective:** Determine the small-signal voltage gain and input and output resistances of a common-source amplifier.

For the circuit shown in Figure 6.13, the parameters are: $V_{DD} = 10 \text{ V}$, $R_1 = 70.9 \text{ k}\Omega$, $R_2 = 29.1 \text{ k}\Omega$, and $R_D = 5 \text{ k}\Omega$. The transistor parameters are: $V_{TN} = 1.5 \text{ V}$, $K_n = 0.5 \text{ mA/V}^2$, and $\lambda = 0.01 \text{ V}^{-1}$. Assume $R_{SI} = 4 \text{ k}\Omega$.

**Solution:** DC Calculations: The dc or quiescent gate-to-source voltage is

$$V_{GSQ} = \left( \frac{R_2}{R_1 + R_2} \right) (V_{DD}) = \left( \frac{29.1}{70.9 + 29.1} \right) (10) = 2.91 \text{ V}$$

The quiescent drain current is

$$I_{DQ} = K_n (V_{GSQ} - V_{TN})^2 = (0.5)(2.91 - 1.5)^2 = 1 \text{ mA}$$

and the quiescent drain-to-source voltage is

$$V_{DSQ} = V_{DD} - I_{DQ}R_D = 10 - (1)(5) = 5 \text{ V}$$

Since $V_{DSQ} > V_{GSQ} - V_{TN}$, the transistor is biased in the saturation region.

**Small-signal Voltage Gain:** The small-signal transconductance $g_m$ is then

$$g_m = 2K_n(V_{GSQ} - V_{TN}) = 2(0.5)(2.91 - 1.5) = 1.41 \text{ mA/V}$$

and the small-signal output resistance $r_o$ is

$$r_o \simeq \left[ \lambda I_{DQ} \right]^{-1} = [(0.01)(1)]^{-1} = 100 \text{ k}\Omega$$

The amplifier input resistance is

$$R_i = R_1 \parallel R_2 = 70.9 \parallel 29.1 = 20.6 \text{ k}\Omega$$

From Figure 6.14 and Equation (6.29), the small-signal voltage gain is

$$A_v = -g_m(r_o \parallel R_D) \left( \frac{R_i}{R_2 + R_{SI}} \right) = -1.41(100 \parallel 5) \left( \frac{20.6}{20.6 + 4} \right)$$

or

$$A_v = -5.62$$

**Input and Output Resistances:** As already calculated, the amplifier input resistance is

$$R_i = R_1 \parallel R_2 = 70.9 \parallel 29.1 = 20.6 \text{ k}\Omega$$

and the amplifier output resistance is

$$R_o = R_D \parallel r_o = 5 \parallel 100 = 4.76 \text{ k}\Omega$$

**Comment:** The resulting $Q$-point is in the center of the load line but not in the center of the saturation region. Therefore, this circuit does not achieve the maximum symmetrical output voltage swing in this case.
**Discussion:** The small-signal input gate-to-source voltage is

\[ V_{gs} = \left( \frac{R_i}{R_i + R_{Sl}} \right) \cdot V_i = \left( \frac{20.6}{20.6 + 4} \right) \cdot V_i = (0.837) \cdot V_i \]

Since \( R_{Sl} \) is not zero, the amplifier input signal \( V_{gs} \) is approximately 84 percent of the signal voltage. This is again called a loading effect. Even though the input resistance to the gate of the transistor is essentially infinite, the bias resistors greatly influence the amplifier input resistance and loading effect.

**Common-Source Amplifier with Source Resistor**

A source resistor \( R_s \) tends to stabilize the Q-point against variations in transistor parameters (Figure 6.18).

![Common-source circuit with source resistor and positive and negative supply voltages](image)

**Figure 6.18** Common-source circuit with source resistor and positive and negative supply voltages

If, for example, the value of the conduction parameter varies from one transistor to another, the Q-point will not vary as much if a source resistor is included in the circuit. However, as shown in the following example, a source resistor also reduces the signal gain. This same effect was observed in BJT circuits when an emitter resistor was included.

The circuit in Figure 6.18 is an example of a situation in which the body effect (not discussed) should be taken into account. The substrate (not shown) would normally be connected to the -5 V supply, so that the body and substrate terminals are not at the same potential. However, in the following example, we will neglect this effect.
Example 6.5 Objective: Determine the small-signal voltage gain of a common-source circuit containing a source resistor.

Consider the circuit in Figure 6.18. The transistor parameters are $V_{TN} = 0.8$ V, $K_n = 1$ mA/V$^2$, and $\lambda = 0$.

Solution: From the dc analysis of the circuit, we find that $V_{GSQ} = 1.50$ V, $I_{DQ} = 0.50$ mA, and $V_{DSQ} = 6.25$ V. The small-signal transconductance is

$$g_m = 2K_n(V_{GS} - V_{TN}) = 2(1)(1.50 - 0.8) = 1.4$$ mA/V

and the small-signal resistance is

$$r_o \equiv \left[\lambda I_{DQ}\right]^{-1} = \infty$$

Figure 6.19 shows the resulting small-signal equivalent circuit.

![Diagram](image)

**Figure 6.19** Small-signal equivalent circuit of NMOS common-source amplifier with source resistor

The output voltage is

$$V_o = -g_m R_D$$

Writing a KVL equation from the input around the gate–source loop, we find

$$V_i = V_{gs} + (g_m V_{gs})R_S = V_{gs}(1 + g_m R_S)$$

or

$$V_{gs} = \frac{V_i}{1 + g_m R_S}$$
The small-signal voltage gain is
\[ A_v = \frac{V_o}{V_i} = \frac{-g_m R_D}{1 + g_m R_S} \]

We may note that if \( g_m \) were large, then the small-signal voltage gain would be approximately
\[ A_v \approx \frac{-R_D}{R_S} \]

Substituting the appropriate parameters into the actual voltage gain expression, we find
\[ A_v = \frac{-(1.4)(7)}{1 + (1.4)(0.5)} = -5.76 \]

**Comment:** A source resistor reduces the small-signal voltage gain. However, as discussed in the last chapter, the \( Q \)-point is more stabilized against variations in the transistor parameters. We may note that the approximate voltage gain gives \( A_v \approx -R_D/R_S = -14 \). Since the transconductance of MOSFETs is generally low, the approximate gain expression is a poor one at best.

**Discussion:** We mentioned that including a source resistor tends to stabilize the circuit characteristics against any changes in transistor parameters. If, for example, the conduction parameter \( K_n \) varies by \( \pm 20 \) percent, we find the following results.

<table>
<thead>
<tr>
<th>( K_n ) (mA/V²)</th>
<th>( g_m ) (mA/V)</th>
<th>( A_v )</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.8</td>
<td>1.17</td>
<td>-5.17</td>
</tr>
<tr>
<td>1.0</td>
<td>1.40</td>
<td>-5.76</td>
</tr>
<tr>
<td>1.2</td>
<td>1.62</td>
<td>-6.27</td>
</tr>
</tbody>
</table>

The change in \( K_n \) produces a fairly large change in \( g_m \). The resulting change in the voltage gain is approximately \( \pm 9.5 \) percent. This change is larger than might be expected because the initial value of \( g_m \) is smaller than that of the bipolar circuit.

**Common-Source Circuit with Source Bypass Capacitor**

A source bypass capacitor added to the common-source circuit with a source resistor will minimize the loss in the small-signal voltage gain, while maintaining \( Q \)-point stability. The \( Q \)-point stability can be further increased by replacing the source resistor with a constant-current source. The resulting circuit is shown in Figure 6.22, assuming an ideal signal source. If the signal frequency is sufficiently large so that the bypass capacitor acts essentially as an ac short-circuit, the source will be held at signal ground.
Example 6.6 Objective: Determine the small-signal voltage gain of a circuit biased with a constant-current source and incorporating a source bypass capacitor.

For the circuit shown in Figure 6.22, the transistor parameters are: \( V_{TN} = 0.8 \text{ V}, \) \( K_n = 1 \text{ mA/V}^2, \) and \( \lambda = 0. \)

Solution: Since the dc gate current is zero, the dc voltage at the source terminal is \( V_S = -V_{GSQ}, \) and the gate-to-source voltage is determined from

\[
I_D = I_Q = K_n (V_{GSQ} - V_{TN})^2
\]

or

\[
0.5 = (1)(V_{GSQ} - 0.8)^2
\]

which yields

\[
V_{GSQ} = -V_S = 1.51 \text{ V}
\]

The quiescent drain-to-source voltage is

\[
V_{DSQ} = V_{DD} - I_Q R_D - V_S = 5 - (0.5)(7) - (-1.51) = 3.01 \text{ V}
\]

The transistor is therefore biased in the saturation region.

The small-signal equivalent circuit is shown in Figure 6.23. The output voltage is
4.6 The Source-Follower Amplifier

The second type of MOSFET amplifier to be considered is the common-drain circuit. An example of this circuit configuration is shown in Figure 6.28.

As seen in this figure, the output signal is taken off the source with respect to ground and the drain is connected directly to $V_{DD}$. Since $V_{DD}$ becomes signal ground in the ac equivalent circuit, we get the name common drain, but the more common name is a source follower. The reason for this name will become apparent as we proceed through the analysis.

Small-Signal Voltage Gain

\[ V_o = -g_m V_{gs} R_D \]

Since $V_{gs} = V_i$, the small-signal voltage gain is

\[ A_v = \frac{V_o}{V_i} = -g_m R_D = -(1.4)(7) = -9.8 \]

**Comment:** Comparing the small-signal voltage gain of 9.8 in this example to the 5.76 calculated in Example 6.5, we see that the magnitude of the gain increases when a source bypass capacitor is included.
\[ V_o = (g_m V_{gs})(R_s || r_o) \]  

(6.30)

Writing a KVL equation from input to output results in the following:

\[ V_{in} = V_{gs} + V_o = V_{gs} + g_m V_{gs}(R_s || r_o) \]  

(6.31(a))

Therefore, the gate-to-source voltage is

\[ V_{gs} = \frac{V_{in}}{1 + g_m(R_s || r_o)} = \left[ \frac{1}{g_m} + \frac{1}{g_m + (R_s || r_o)} \right] \cdot V_{in} \]  

(6.31(b))

Equation (6.31(b)) is written in the form of a voltage-divider equation, in which the gate-to-source of the NMOS device looks like a resistance with a value of \(1/g_m\). More accurately, the effective resistance looking into the source terminal (ignoring \(r_o\)) is \(1/g_m\). The voltage \(V_{in}\) is related to the source input voltage \(V_i\) by

\[ V_{in} = \left( \frac{R_i}{R_i + R_{Si}} \right) \cdot V_i \]  

(6.32)

where \(R_i = R_1 || R_2\) is the input resistance to the amplifier.

Substituting Equations (6.31(b)) and (6.32) into (6.30), we have the small-signal voltage gain:

\[ A_v = \frac{V_o}{V_i} = \frac{g_m(R_s || r_o)}{1 + g_m(R_s || r_o)} \cdot \left( \frac{R_i}{R_i + R_{Si}} \right) \]  

(6.33(a))

or

\[ A_v = \frac{1}{g_m + R_s || r_o} \cdot \left( \frac{R_i}{R_i + R_{Si}} \right) \]  

(6.33(b))

which again is written in the form of a voltage-divider equation. An inspection of Equation 6.33(b) shows that the magnitude of the voltage gain is always less than unity. This result is consistent with the results of the BJT emitter-follower circuit.
Example 6.7  **Objective:** Calculate the small-signal voltage gain of the source-follower circuit in Figure 6.28.

Assume the circuit parameters are $V_{DD} = 12$ V, $R_1 = 162$ kΩ, $R_2 = 463$ kΩ, and $R_S = 0.75$ kΩ, and the transistor parameters are $V_{TN} = 1.5$ V, $K_n = 4$ mA/V$^2$, and $\lambda = 0.01$ V$^{-1}$. Also assume $R_{SI} = 4$ kΩ.

**Solution:** The dc analysis results are $I_D = 7.97$ mA and $V_{GSQ} = 2.91$ V. The small-signal transconductance is therefore

$$g_m = 2K_n(V_{GSQ} - V_{TN}) = 2(4)(2.91 - 1.5) = 11.3 \text{ mA/V}$$

and the small-signal transistor resistance is

$$r_o \approx \frac{1}{\lambda I_D} = \frac{1}{(0.01)(7.97)} = 12.5 \text{ kΩ}$$

The amplifier input resistance is

$$R_i = R_1 || R_2 = 162 || 463 = 120 \text{ kΩ}$$

The small-signal voltage gain then becomes

$$A_v = \frac{g_m(R_S || r_o)}{1 + g_m(R_S || r_o)} \cdot \frac{R_i}{R_i + R_{SI}} = \frac{(11.3)(0.75 || 12.5)}{1 + (11.3)(0.75 || 12.5)} \cdot \frac{120}{120 + 4} = +0.860$$

**Comment:** The magnitude of the small-signal voltage gain is less than 1. An examination of Equation (6.33(b)) shows that this is always true. Also, the voltage gain is positive, which means that the output signal voltage is in phase with the input signal voltage. Since the output signal is essentially equal to the input signal, the circuit is called a source follower.

**Discussion:** The expression for the voltage gain of the source follower is essentially identical to that of the bipolar emitter follower. Since the transconductance of the BJT is, in general, larger than that of the MOSFET, the voltage gain of the emitter follower will be closer to unity than that of the MOSFET source follower.

Although the voltage gain is slightly less than 1, the source follower is an extremely useful circuit because the output resistance is less than that of a common-source circuit. A small output resistance is desirable when the circuit is to act as an ideal voltage source and drive a load circuit without suffering any loading effects.
4.7 Input and Output impedance

The input resistance $R_i$, as defined in Figure 6.29(b), is the Thevenin equivalent resistance of the bias resistors. Even though the input resistance to the gate of the MOSFET is essentially infinite, the input bias resistances do create a loading effect. This same effect was seen in the common-source circuits.

To calculate the output resistance, we set all independent small-signal sources equal to zero, apply a test voltage to the output terminals, and measure a test current. Figure 6.31 shows the circuit we will use to determine the output resistance of the source follower shown in Figure 6.28.

![Figure 6.31 Equivalent circuit of NMOS source follower, for determining output resistance](www.vtucs.com)

We set $V_i = 0$ and apply a test voltage $V_x$. Since there are no capacitances in the circuit, the output impedance is simply an output resistance, which is defined as

$$R_o = \frac{V_x}{I_x}$$

Writing a KCL equation at the output source terminal produces

$$I_x + g_m V_{gs} = \frac{V_x}{R_s} + \frac{V_x}{r_o}$$

Since there is no current in the input portion of the circuit, we see that $V_{gs} = -V_x$. Therefore, Equation (6.35) becomes

$$I_x = V_x \left( g_m + \frac{1}{R_s} + \frac{1}{r_o} \right)$$

or

$$\frac{I_x}{V_x} = g_m + \frac{1}{R_s} + \frac{1}{r_o}$$

The output resistance is then

$$R_o = \frac{1}{g_m} || R_s || r_o$$

From Figure 6.31, we see that the voltage $V_{gs}$ is directly across the current source $g_m V_{gs}$. This means that the effective resistance of the device is $1/g_m$. The output resistance given by Equation (6.37) can therefore be written directly. This result also means that the resistance looking into the source terminal (ignoring $r_o$) is $1/g_m$, as previously noted.
Example 6.9 Objective: Calculate the output resistance of a source-follower circuit.

Consider the circuit shown in Figure 6.28 with circuit and transistor parameters given in Example 6.7.

Solution: The results of Example 6.7 are: $R_S = 0.75 \, \text{k} \Omega$, $r_o = 12.5 \, \text{k} \Omega$, and $g_m = 11.3 \, \text{m} \text{A/V}$. Using Figure 6.31 and Equation (6.37), we find

$$R_o = \frac{1}{g_m} || R_S || r_o = \frac{1}{11.3} || 0.75 || 12.5$$

or

$$R_o = 0.0787 \, \text{k} \Omega = 78.7 \, \Omega$$

Comment: The output resistance of a source-follower circuit is dominated by the transconductance parameter. Also, because the output resistance is very low, the source follower tends to act like an ideal voltage source, which means that the output can drive another circuit without significant loading effects.

4.8 The Common-Gate Configuration

The third amplifier configuration is the common-gate circuit. To determine the small-signal voltage and current gains, and the input and output impedances, we will use the same small-signal equivalent circuit for the transistor that was used previously. The dc analysis of the common-gate circuit is the same as that of previous MOSFET circuits.

Small-Signal Voltage and Current Gains

In the common-gate configuration, the input signal is applied to the source terminal and the gate is at signal ground. The common-gate configuration shown in Figure 6.344 is biased with a constant-current source $I_Q$.

![Common-gate circuit](image)

**Figure 6.34** Common-gate circuit
The gate resistor $R_G$ prevents the buildup of static charge on the gate terminal, and the capacitor $C_G$ ensures that the gate is at signal ground. The coupling capacitor $C_{C1}$ couples the signal to the source, and coupling capacitor $C_{C2}$ couples the output voltage to load resistance $R_L$.

The small-signal equivalent circuit is shown in Figure 6.35. The small-signal transistor resistance $r_o$ is assumed to be infinite.

![Small-signal equivalent circuit of common-gate amplifier](image)

**Figure 6.35** Small-signal equivalent circuit of common-gate amplifier

The output voltage is

$$V_o = -(g_mV_{gs})(R_D \parallel R_L)$$

Writing the KVL equation around the input, we find

$$V_i = I_i R_S - V_{gs}$$

where $I_i = -g_m V_{gs}$. The gate-to-source voltage can then be written as

$$V_{gs} = \frac{-V_i}{1 + g_m R_S}$$

The small-signal voltage gain is found to be

$$A_v = \frac{V_o}{V_i} = \frac{g_m (R_D \parallel R_L)}{1 + g_m R_S}$$

Also, since the voltage gain is positive, the output and input signals are in phase.

In many cases, the signal input to a common-gate circuit is a current. Figure 6.36 shows the small-signal equivalent common-gate circuit with a Norton equivalent circuit as the signal source. We can calculate a current gain. The output current $I_o$ can be written

$$I_o = \left(\frac{R_D}{R_D + R_L}\right)(-g_m V_{gs})$$

At the input we have

$$I_i + g_m V_{gs} + \frac{V_{gs}}{R_S} = 0$$
Input and Output Impedance

In contrast to the common-source and source-follower amplifiers, the common-gate circuit has a low input resistance because of the transistor. However, if the input signal is a current, a low input resistance is an advantage. The input resistance is defined as

\[ g_m = 2K_n(V_{GSQ} - V_{TN}) = 2(1)(2 - 1) = 2 \text{ mA/V} \]

From Equation (6.45), we can write the output current as

\[ I_o = I_i \left( \frac{R_D}{R_D + R_L} \right) \cdot \left( g_m R_{SI} \right) \cdot \left( \frac{1}{1 + g_m R_{SI}} \right) \]

The output voltage is \( V_o = I_o R_L \), so we find

\[ V_o = I_i \left[ \frac{R_L R_D}{R_D + R_L} \right] \cdot \left( \frac{g_m R_{SI}}{1 + g_m R_{SI}} \right) \]

\[ = \left[ \frac{(10)(4)}{4 + 10} \right] \cdot \left[ \frac{(2)(50)}{1 + (2)(50)} \right] \cdot (0.1) \sin \omega t \]

or

\[ V_o = 0.283 \sin \omega t \text{ V} \]

**Comment:** As with the BJT common-base circuit, the MOSFET common-gate amplifier is useful if the input signal is a current.
4.9 The Three Basic Amplifier Configurations: Summary and Comparison

Table 6.1 is a summary of the small-signal characteristics of the three amplifier configurations.

<table>
<thead>
<tr>
<th>Configuration</th>
<th>Voltage gain</th>
<th>Current gain</th>
<th>Input resistance</th>
<th>Output resistance</th>
</tr>
</thead>
<tbody>
<tr>
<td>Common source</td>
<td>$A_v &gt; 1$</td>
<td>—</td>
<td>$R_{TH}$</td>
<td>Moderate to high</td>
</tr>
<tr>
<td>Source follower</td>
<td>$A_v \approx 1$</td>
<td>—</td>
<td>$R_{TH}$</td>
<td>Low</td>
</tr>
<tr>
<td>Common gate</td>
<td>$A_v &gt; 1$</td>
<td>$A_I \approx 1$</td>
<td>Low</td>
<td>Moderate to high</td>
</tr>
</tbody>
</table>

The input resistance looking directly into the gate of the common-source and source-follower circuits is essentially infinite at low to moderate signal frequencies. However, the input resistance of these discrete amplifiers is the Thevenin equivalent resistance $R_{TH}$ of the bias resistors. In contrast, the input resistance to the common-gate circuit is generally in the range of a few hundred ohms.

The output resistance of the source follower is generally in the range of a few hundred ohms. The output resistance of the common-source and common-gate configurations is dominated by the resistance $R_D$. The specific characteristics of these single-stage amplifiers are used in the design of multistage amplifiers.

In the last chapter, we considered three all-MOSFET inverters and plotted the voltage transfer characteristics. All three inverters use an n-channel enhancement-mode driver transistor. The three types of load devices are an n-channel enhancement-mode device, an n-channel depletion-mode device, and a p-channel enhancement-mode device. The MOS transistor used as a load device is referred to as an active load. We mentioned that these three circuits can be used as amplifiers.

In this section, we revisit these three circuits and consider their amplifier characteristics. We will emphasize the small-signal equivalent circuits. This section serves as an introduction to more advanced MOS integrated circuit amplifier designs considered in Part II of the text.

**NMOS Amplifiers with Enhancement Load**

The characteristics of an n-channel enhancement toad device were presented in the last chapter. Figure 6.38(a) shows an NMOS enhancement load transistor.
Figure 6.38  (a) NMOS enhancement-mode transistor with gate and drain connected in a load device configuration and (b) current–voltage characteristics of NMOS enhancement load transistor.

and Figure 6.38(b) shows the current–voltage characteristics. The threshold voltage is $V_{TNL}$.

Figure 6.39(a) shows an NMOS amplifier with an enhancement load.
The driver transistor is $M_D$ and the load transistor is $M_L$. The characteristics of transistor $M_D$ and the load curve are shown in Figure 6.39(b). The load curve is essentially the mirror image of the $i$-$v$ characteristic of the load device. Since the $i$-$v$ characteristics of the load device are nonlinear, the load curve is also nonlinear. The load curve intersects the voltage axis at $V_{DD} - V_{TNL}$, which is the point where the current in the enhancement load device goes to zero. The transition point is also shown on the curve.

The voltage transfer characteristic is also useful in visualizing the operation of the amplifier. This curve is shown in Figure 6.39(c). When the enhancement-mode driver first begins to conduct, it is biased in the saturation region. For use as an amplifier, the circuit Q-point should be in this region, as shown in both Figures 6.39(b) and (c).

We can now apply the small-signal equivalent circuits to find the voltage gain. In the discussion of the source follower, we found that the equivalent resistance looking into the source terminal (with $R_S = \infty$) was $R_O = (1 / g_m) \parallel r_o$. The small-signal equivalent circuit of the inverter is given in Figure 6.40, where the subscripts D and L refer to the driver and load transistors, respectively. We are again neglecting the body effect of the load transistor.
The small-signal voltage gain is

\[ A_v = \frac{V_o}{V_i} = -g_{mD} \left( \frac{1}{g_{mL} || r_{oL}} \right) \]  

(6.49)

Since, generally, \( 1/g_{mL} \ll r_{oL} \) and \( 1/g_{mD} \ll r_{oD} \), the voltage gain, to a good approximation is given by

\[ A_v = \frac{-g_{mD}}{g_{mL}} = -\sqrt{\frac{K_{nD}}{K_{nL}}} = -\sqrt{\frac{(W/L)_D}{(W/L)_L}} \]  

(6.50)

The voltage gain, then, is related to the size of the two transistors.

### 4.10 Recommended Questions

1. Which amplifiers are classified as power amplifiers? Explain the general features of a power amplifier.
2. Give the expression for dc power input, ac power output and efficiency of a series fed, directly, coupled class A amplifier.
3. When the power dissipation is maximum, in class A amplifiers? What is the power dissipation rating of a transistor?
4. Explain with neat circuit diagram, the working of a transformer coupled class A power amplifier.
5. Prove that the maximum efficiency of a transformer coupled class A amplifier is 50%.
6. What is harmonic distortion? How the output signal gets distorted due to the harmonic distortion.
7. Draw a neat circuit diagram of push pull class B amplifier. Explain its working.
8. Draw the circuit diagram of class B push pull amplifier and discuss
   a. Its merits.
   b. Cross-over distortion
9. Prove that the maximum efficiency of a class B amplifier is 78.5%.
10. Write a short note on class D amplifier.
11. Give the classification of multistage amplifier. Explain the various distortions in amplifiers. (July-2007)
Unit – 5: Large Signal Amplifiers

Large signal or power amplifiers provide power amplification and are used in applications to provide sufficient power to the load or the power device. The output power delivered by these amplifiers is of the order of few watts to few tens of watts. They handle moderate-to-high levels of current and voltage signals as against small levels of current and voltage signals in the case of small signal amplifiers.

5.1 Classification of Large Signal Amplifiers

On the basis of their circuit configurations and principle of operation, amplifiers are classified into different classes.

Class A Amplifiers

A class A power amplifier is defined as a power amplifier in which output current flows for the full-cycle (360°) of the input signal. In other words, the transistor remains forward biased throughout the input cycle. The active device in a class A amplifier operates during the whole during the whole of the input cycle and the output signal is an amplified replica of the input signal with no clipping. The amplifying element is so biased that it operates over the linear region of its output characteristics during full period of the input cycle and is always conducting to some extent.

Class A amplifiers offer very poor efficiency and a maximum of 50% efficiency is possible in these amplifiers.

Class B Amplifier

Class B amplifiers operate only during the half of the input cycle. Class B amplifiers offer much improved efficiency over class A amplifiers with a possible maximum of 78.5%. They also create a large amount of distortion.
Class AB Amplifier

In a class AB amplifier, the amplifying device conducts for a little more than half of the input waveform. They sacrifice some efficiency over class B amplifiers but they offer better linearity than class B amplifier. They offer much more efficiency than class A amplifiers.

Class C Amplifiers

Class C amplifiers conducts for less than half cycle of the input signal resulting in a very high efficiencies up to 90%. But they are associated with very high level of distortion at the output. Class C amplifiers operate in two modes, namely, the tuned mode and the unturned mode.

Class D Amplifiers

Class D amplifiers use the active device in switching mode to regulate the output power. These amplifiers offer high efficiencies and do not require heat sinks and transformers. These amplifiers use pulse width modulation (PWM), pulse density modulation or sigma delta modulation to convert the input signal into a
string of pulses. The pulse width of the PWM output waveform at any time instant is directly proportional to the amplitude of the input signal.

5.2 Large signal Amplifier Characteristics

The main characteristics that define the performance of a power amplifier are efficiency, distortion level and output power.

Efficiency

Efficiency of an amplifier is defined as the ability of the amplifier to convert the DC input power of the supply into an AC output power that can be delivered to the load. The expression for efficiency is given by

\[
\eta = \frac{P_o}{P_i} \times 100\%
\]

Where, \( P_o \) is the AC output power delivered to the load and \( P_i \) is the DC input power.

Harmonic Distortion

Distortion in large signal amplifiers is mainly caused due to harmonic distortion. Harmonic distortion refers to the distortion in the amplitude of the output signal of an amplifier caused due to the non-linearity in the characteristics of the active device used for amplification. In other words, the active device does not equally amplify all portions of the input signal over its positive and negative excursions. The distortion is more in the case of a large input signal level.

The total current of an amplifier is given by

\[
I_{\text{tot}} = I_o + I_q
\]
Where \( I_o \) is the alternating portion of the output current and 

\[ I_Q \] is the output DC current under quiescent condition.

Where 
\[ I_o = A_0 + A_1 \cos \omega t + A_2 \cos 2\omega t \]

\[ \therefore I_{ot} = I_Q + A_0 + A_1 \cos \omega t + A_2 \cos 2\omega t \]

Where \( A_0 \) is the extra DC component due to rectification of the signal;

\( A_1 \) is the amplitude of the desired signal at the fundamental input signal frequency \( \omega \);

\( A_2 \) is the amplitude of the desired signal at the fundamental input signal frequency \( 2\omega \);

Second harmonic distortion is a measure of the relative amount of second harmonic component to the fundamental frequency component and is expressed as 

\[ D_2 = \frac{A_2}{A_1} \times 100\% \]

III\(^0\) the third harmonic distortion component (D3) is given by 

\[ D_3 = \frac{A_3}{A_1} \times 100\% \]

The fourth harmonic distortion component (D3) is given by 

\[ D_4 = \frac{A_4}{A_1} \times 100\% \]

The total harmonic distortion (D) is given by the square root of the mean square of the individual harmonic components.

\[ D = \sqrt{D_2^2 + D_3^2 + D_4^2 + ...} \]

The power delivered at the fundamental frequency is 

\[ P_1 = \frac{A_1^2 R_L}{2} \]

The total power output is given by 

\[ P = (A_1^2 + A_2^2 + A_3^2 + ... \times \frac{R_L}{2} \]

Therefore, 

\[ P = (1 + D_2^2 + D_3^2 + ...) \times P_1 \]

\[ \therefore P = (1 + D^2) \times P_1 \]
5.3 Feedback Amplifiers

Classification of Amplifiers

On the basis of the input and output parameters of interest, amplifiers are classified as voltage amplifiers, current amplifiers, transresistance and transconductance amplifiers.

In the case of a voltage amplifier, a small change in the input voltage produces a large change in the output voltage. Voltage gain, which is the ratio of the change in the output voltage to change in the input voltage, is the gain parameter.

In the case of a current amplifier, a small change in the input current produces a large change in the output current. Current gain, which is the ratio of change in the output current to the change in the input current, is the gain parameter.

In the case of a transresistance amplifier, a small change in the input current produces a large change in the output voltage. Ratio of the change in the output voltage to change in the input current is the gain parameter. The gain parameter has the units of resistance.

In the case of a transconductance amplifier, a small change in the input voltage produces a large change in the output current. Ratio of the change in the output current to the change in the input voltage is the gain parameter. The gain parameter has the units of conductance.

Amplifier with Negative Feedback

In a negative feedback amplifier, a sample of the output signal is fed back to the input and the feedback signal is combined with the externally applied input signal in a subtractor circuit as shown in the figure below.

![Feedback Network Diagram]

The actual signal $X_i$, applied to the basic amplifier is the difference of the externally applied input signal $X_s$ and the feedback signal $X_f$. The generalized representation of input and output signals is intended to indicate that the signal could either be a current or a voltage signal.
The gain parameters could be a voltage gain, current gain, transresistance or transconductance. The gain has two values namely gain of the basic amplifier without feedback and the gain of the amplifier with feedback.

From the above circuit, the actual signal applied to the amplifier input $X_i$ is the difference of externally applied input signal $X_s$ and the feedback signal $X_f$. It is given by

$$X_i = X_s - X_f$$  ---  (1)

Also, $X_i = \beta X_o$, $X_o = AX_i$ and $X_o = A_i X_s$

Substituting for $X_i$ and $X_f$ in eq (1) we get

$$\frac{X_o}{A} = X_s - \beta X_o$$

Simplifying the equation we get

$$A_f = \frac{X_o}{X_s} = \frac{A}{1 + \beta A}$$

The feedback is expressed in decibels. It is given by

$$dB \text{ of feedback} = 20 \log \left( \frac{A_f}{A} \right) = 20 \log \left( \frac{1}{1 + \beta A} \right)$$

The three important assumptions to be satisfied by a feedback network is,

1. The input signal is transmitted to the output through the amplifier only and not through the feedback network. That is, forward transmission through feedback network is zero. This further implies that if the gain of the amplifier were reduced to zero, the output must drop to zero.
2. The feedback signal is transmitted from output to input through feedback network only. That is, reverse transmission through amplifier is zero.
3. The feedback factor is independent of load and source resistance.
Advantages of Negative Feedback

1. Gain parameter is independent of variation in values of components used in building the basic amplifier.
2. Bandwidth is increased.
3. Distortion is reduced and linearity is increased.
4. Noise is reduced.
5. Input resistance can be reduced or decreased depending on the feedback topology.
6. Output resistance can be reduced or decreased depending on the feedback topology.

5.4 Stability of Gain Factor

The introduction of negative feedback makes the amplifier insensitive to variations in the values of the components and parameters of the active devices used in building the amplifier. The expression that relates percentage variation in the gain parameter of the amplifier with feedback to the percentage variation in the same without feedback can be derived as follows.

$$A_f = \frac{A}{1 + \beta A}$$

Differentiating the above equation with respect to $A$ we get,

$$\left| \frac{dA_f}{A_f} \right| = \left| \frac{1}{1 + \beta A} \cdot \frac{dA}{A} \right|$$

$|(1+\beta A)|$ is called the desensitivity parameter $D$. thus, percentage variation in gain with feedback is equal to the percentage variation in gain without feedback divided by the desensitivity parameter $D$.

Effect on bandwidth

Bandwidth increases with introduction of negative feedback. Increase in bandwidth results from the fact that amplifier exhibit a constant gain-bandwidth product. Reduction in gain is, therefore, accompanied by increase in bandwidth. Bandwidth increases by the same desensitivity factor $D = 1+\beta A$ by which the gain reduces. Bandwidth of amplifier with feedback is given by

$$(BW)_f = BW \times (1+\beta A)$$

Effect on Non-Linear Distortion

It can be proved that non-linear distortion decreases by the desensitivity factor $D = 1+\beta A$. Let us assume that the distortion levels without and with negative feedback are $D_2$ and $D_{2f}$, respectively. $D_2$ is the distortion contributed by the active device. In the presence of feedback, $D_{2f}$ appears as $-\beta AD_{2f}$ at the input of the amplifier. The following gives the expression for $D_{2f}$. 
\[ D_2^2\beta AD_2f = \frac{D_2}{1 + \beta A} \]

**Effect on Noise**

Introduction of negative feedback acts on the noise generated in the amplifier in the same manner as it does on the non-linear distortion. Let us assume \( N_f \) and \( N \) are the noise levels with and without negative feedback respectively. Reduction in noise is governed by

\[ N_f = \frac{N}{1 + \beta A} \]

**Effect on Input Resistance**

Input resistance in the case of an amplifier with negative feedback is affected depending upon how the feedback signal is connected to the source of external input signal. The input resistance increases if the feedback signal is connected in series with the source of input and decreases if the feedback signal is connected across it in shunt.

Thus in the case of voltage-series and current-series feedback, \( R_i \) is the input resistance without feedback and the input resistance without feedback \( R_{if} \) is given by

\[ R_{if} = R_i \times (1 + \beta A) \]

Input resistance in the case of voltage-shunt and current-shunt feedback is given by

\[ R_{if} = \frac{R_i}{1 + \beta A} \]

**Effect on Output Resistance**

Output resistance in the case of an amplifier with negative feedback is affected depending upon how the feedback signal is connected to the source of external input signal. The output resistance increases in the case of output is current and decreases in the case of output is voltage.

Thus in the case of current-series and current-shunt feedback, \( R_o \) is the input resistance without feedback and the input resistance without feedback \( R_{of} \) is given by

\[ R_{of} = R_o \times (1 + \beta A) \]

Input resistance in the case of voltage-series and voltage-shunt feedback is given by

\[ R_{of} = \frac{R_o}{1 + \beta A} \]
5.5 Feedback Topologies

On the basis of the nature of sampled signal and the mode in which it is fed back to the input, there are four feedback topologies. They are

1. Voltage-series feedback topology (series-shunt)
2. Voltage-shunt feedback topology (shunt-shunt)
3. Current-series feedback topology (series-series)
4. Current-shunt feedback topology (shunt-series)

**Voltage-Series Feedback (Series-Shunt)**

Schematic arrangement of voltage-series (series-shunt) feedback

Equivalent circuit of voltage-series feedback
5.6 Recommended Questions

1. List the important features of MOSFET.
2. With the help of neat diagrams explain the operation and drain characteristics of n-channel depletion type MOSFET. Explain clearly the mechanism of “Pinch-off Condition”. Sketch the drain characteristics and define $r_d$.
3. Explain the constructional features of a depletion mode MOSFET and explain its basic operation.
4. With the help of neat diagram explain the operation of an n-channel enhancement type MOSFET.
5. Sketch the drain characteristics of MOSFET for different values of $V_{GS}$ and mark different region of operation.
6. Draw and explain the drain characteristics of n-channel enhancement type MOSFET.
7. Sketch the graphic symbols for: n-channel JFET, p-channel JFET, n-channel enhancement type MOSFET, p-channel enhancement type MOSFET, n-channel depletion type MOSFET and p-channel depletion type MOSFET.
8. Explain the structure of the depletion mode MOSFET. And the D-MOSFET curves. (July-2008)
UNIT – 6: Sinusoidal Oscillators

6.1 Principles of oscillators

Here we consider the principles of oscillators that produce approximately sinusoidal waveforms. (Other oscillators, such as multivibrators, operate somewhat differently.) Because the waveforms are sinusoidal, we use phasor analysis.

A sinusoidal oscillator ordinarily consists of an amplifier and a feedback network. Let’s consider the following idealized configuration to begin understanding the operation of such oscillators.

We begin consideration of sinusoidal oscillators by conducting a somewhat artificial thought experiment with this configuration. Suppose that initially, as shown in the figure, the switch, S, connects the input of the amplifier is connected to the driver. Suppose, furthermore, that the complex constant, \( F_{vv} \), in the feedback network is adjusted (designed) to make the output of the feedback network exactly equal \( V_{in} \), the input voltage provided by the driver circuit. Then suppose that, instantaneously (actually, in a time negligible in comparison to the period of the sinusoids), the switch S disconnects the driver circuit from the input of the amplifier and immediately connects the identical voltage, \( V_{in} \), supplied by the feedback network to the input of the amplifier. A circuit, of course, cannot distinguish between two identical voltages and, therefore, the amplifier continues to behave as before. Specifically, it continues to produce a sinusoidal output. Now, however, it produces the sinusoidal output without connection to a driver. The amplifier is now self-driven, or self-excited, and functions as an oscillator.

We now analyze the behavior of the amplifier when it is connected to produce self-excited oscillations to develop a consistency condition that must be satisfied if such operation is to be possible. First, we note that the output voltage of the amplifier can be written in terms of the input voltage, \( V_{in} \):

\[
V_{out} = \frac{Z_L}{Z_L + Z_o} A_v V_{in} \equiv AV_{in}
\]
where we have written the gain, $A$, of the amplifier, under load, as

$$A = \frac{Z_L}{Z_L + Z_o} A_v$$

But the output of the feedback network is $F_v V_{out}$ where $F_v$ has been chosen so that

$$F_v V_{out} = V_{in}$$

If we substitute this result into our earlier result for $V_{out}$, we find

$$V_{out} = AV_{in} = AF_v V_{out}$$

or

$$(1 - AF_v) V_{out} = 0$$

Of course, $V_{out} \neq 0$ for a useful oscillator so we must have

$$AF_v = 1$$

Although it is usually summarized as requiring the complex loop gain to be unity as a condition of oscillation, let’s examine this condition, known as the Barkhausen condition for oscillation, to gain a better understanding of what it means. To begin with, $A$ and $F_v$ are complex numbers that can be written in polar form:

$$A = |A| e^{j\phi}$$

$$F_v = |F_v| e^{j\delta}$$

Thus, the Barkhausen condition can be written as

$$AF_v = |A| e^{j\phi} |F_v| e^{j\delta} = |A||F_v| e^{j(\phi + \delta)} = 1$$

or

$$|A||F_v| e^{j(\phi + \delta)} = 1$$

This equation, being complex, gives two real equations, one from the magnitude and one from the angle:

Magnitude: $|A||F_v| = 1$
Angle: $\phi + \theta = n 2\pi, \ n = 0, \pm 1, \pm 2, \ldots$

The magnitude portion of the Barkhausen condition requires a signal that enters the amplifier and undergoes amplification by some factor to be attenuated by the same factor by the feedback network before the signal reappears at the input to the amplifier. The magnitude condition therefore ensures that the amplitude of oscillation remains constant over time. If it were true that $|A| F_{in} < 1$, then the amplitude of the oscillations would gradually decrease each time the signal passed around the loop through the amplifier and the feedback network. Similarly, if it were true that $|A| F_{in} > 1$, then the amplitude of the oscillations would gradually increase each time the signal passed around the loop through the amplifier and the feedback network. Only if $|A| F_{in} = 1$ does the amplitude of the oscillations remain steady.

The angle portion of the Barkhausen condition requires that the feedback network complement any phase shift experienced by a signal when it enters the amplifier and undergoes amplification so that the total phase shift around the signal loop through the amplifier and the feedback network totals to 0, or to what amounts to the same thing, an integral multiple of $2\pi$. Without this condition, signals would interfere destructively as they travel around the signal loop and oscillation would not persist because of the lack of reinforcement. Because the phase shift around the loop usually depends on frequency, the angle part of the Barkhausen condition usually determines the frequency at which oscillation is possible. In principle, the angle condition can be satisfied by more than one frequency. In laser feedback oscillators (partially reflecting mirrors provide the feedback), indeed, the angle part of the Barkhausen condition is often satisfied by several closely spaced, but distinct, frequencies. In electronic feedback oscillators, however, the circuit usually can satisfy the angle part of the Barkhausen condition only for a single frequency.

Although the Barkhausen condition is useful for understanding basic conditions for oscillation, the model we used to derive it gives an incomplete picture of how practical oscillators operate. For one thing, it suggests that we need a signal source to start up an oscillator. That is, it seems that we need an oscillator to make an oscillator. Such a circumstance would present, of course, a very inconvenient version of the chicken-and-the-egg dilemma. Second, the model suggests that the amplitude of the oscillations can occur at any amplitude, the amplitude apparently being determined by the amplitude at which the amplifier was operating when it was switched to self-excitation. Practical oscillators, in contrast, start by themselves when we flip on a switch, and a particular oscillator always gives approximately the same output amplitude unless we take specific action to adjust it in some way.

Let’s first consider the process through which practical oscillators start themselves. The key to understanding the self-starting process is to realize that in any practical circuit, a variety of processes produce noise voltages and currents throughout the circuit. Some of the noise, called Johnson noise, is the result of the tiny electric fields produced by the random thermal motion of electrons in the components. Other noise results during current flow because of the discrete charge on electrons, the charge carriers. This noise is analogous to the acoustic noise that results from dumping a shovel-full of marbles onto a concrete sidewalk, in comparison to that from dumping a shovel-full of sand on the same sidewalk. The lumpiness of the mass of the marbles produces more noise than the less lumpy grains of sand. The lumpiness of the charge on the electrons leads to electrical noise, called shot noise, as they carry electrical current. Transient voltages and currents produced during start-up by power supplies and other circuits also can produce noise in the circuit. In laser feedback oscillators, the noise to initiate oscillations is provided by spontaneous emission of photons. Amplification in lasers occurs through the process of stimulated emission of photons.
Whatever the source, noise signals can be counted upon to provide a small frequency component at any frequency for which the Barkhausen criterion is satisfied. Oscillation begins, therefore, as this frequency component begins to loop through the amplifier and the feedback network. The difficulty, of course, is that the amplitude of the oscillations is extremely small because the noise amplitude at any particular frequency is likely to be measured in microvolts. In practice, therefore, we design the oscillator so that loop gain, \( |A||F_{vv}| \), is slightly greater than one:

\[
|A||F_{vv}| > 1
\]

With the loop gain slightly greater than one, the small noise component at the oscillator frequency is amplified slightly each time it circulates around the loop through the amplifier and the feedback network, and hence gradually builds to useful amplitude. A problem would occur if the amplitude continued to build toward infinite amplitude as the signal continued to circulate around the loop. Our intuition tells us, of course, that the amplitude, in fact, is unlikely to exceed some fraction of the power supply voltage (without a step-up transformer or some other special trick), but more careful consideration of how the amplitude is limited in practical oscillators provides us with some useful additional insight.

Initially, let’s consider the amplifier by itself, without the feedback network. Suppose that we drive the amplifier with a sinusoidal generator whose frequency is the same as that of the oscillator in which the amplifier is to be used. With the generator, suppose we apply sinusoids of increasing amplitude to the amplifier input and observe its output with an oscilloscope. For sufficiently large inputs, the output becomes increasingly distorted as the amplitude of the driving sinusoid becomes larger. For large enough inputs, we expect the positive and negative peaks of the output sinusoids to become clipped so that the output might even resemble a square wave more than a sinusoid. Suppose we then repeat the experiment but observe the oscillator output with a tuned voltmeter set to measure the sinusoidal component of the output signal at the fundamental oscillator frequency, the only frequency useful for maintaining self-excited oscillations when the amplifier is combined with the feedback circuit. Then, we would measure an input/output characteristic curve for the amplifier at the fundamental oscillator frequency something like that shown in the following sketch.

![Input/Output Characteristic Curve](https://example.com/characteristic_curve.png)

From the curve above, note that, at low levels of input amplitude, \( |V_{in}| \), the output amplitude, \( |V_{out,f}| \), of the sinusoidal component at the fundamental frequency increases in direct proportion to the input amplitude. At sufficiently high output levels, however, note that a given increment in input amplitude
produces a diminishing increase in the output amplitude (at the fundamental frequency). Physically, as the output becomes increasingly distorted at larger amplitudes, harmonic components with frequencies at multiples of the fundamental frequency necessarily increase in amplitude. Because the total amplitude is limited to some fraction of the power supply voltage, the sinusoidal component at the fundamental frequency begins to grow more slowly as the input amplitude increases and causes the amplitude of the distortion components to increase, as well. Thus, the amplitude of the component of the output at the fundamental frequency eventually must decrease as the input amplitude increases to accommodate the growing harmonic terms that accompany the rapidly worsening distortion. Effectively, the magnitude of the voltage gain, $|A|$, for the fundamental frequency decreases at large amplitudes.

Now let's reconsider the amplifier in its oscillator environment, that is, with the feedback network designed so that $|A||F_{vv}| > 1$. As the oscillations build up from noise and increase to larger and larger amplitudes, they eventually reach amplitudes at which the magnitude of the voltage gain, $|A|$, begins to decrease. As a consequence, the loop gain, $|A||F_{vv}|$, begins to decrease. The amplitude of the oscillations grows until the decreasing $|A|$ reduces the loop gain, $|A||F_{vv}|$, to unity:

$$|A||F_{vv}| = 1$$

At that point, the oscillations cease growing and their amplitude becomes stable, at least as long as the gain characteristic of the amplifier shown in the curve above do not change.

In summary, the small signal loop gain in practical amplifiers is chosen so that $|A||F_{vv}| > 1$ and oscillations grow from small noise components at the oscillator frequency. The output climbs along the input/output characteristic curve of the amplifier at the fundamental frequency until the voltage gain drops enough to make $|A||F_{vv}| = 1$, at which point the oscillator amplitude stops growing and maintains a steady level. The amplifier input/output characteristic curve therefore explains why practical oscillators operate at approximately the same amplitude each time we turn them on. Note designers sometimes add nonlinearities, voltage-limiting circuits with diodes, for example, to gain more direct control of the oscillator amplitude.

The analysis of oscillator operation based on the input/output characteristic of the amplifier at the fundamental frequency can help illuminate one more aspect of the operation of practical sinusoidal oscillators: distortion in the output waveform. From the discussion above, it is clear that the higher the oscillations climb along the input/output characteristic curve, the more distortion in the output worsens. In addition, it is clear that the more the loop gain, $|A||F_{vv}|$, exceeds unity at small amplitudes, the higher the oscillations climb along the curve, and the more distorted the output will become, before the amplitude of the oscillations stabilizes. Thus, it is clear that, in the design process, $|A||F_{vv}|$ should not be chosen to exceed unity very much, even at small signals. On the other hand, if $|A||F_{vv}|$ is chosen too close to unity in an effort to reduce distortion, then even small changes in the amplification characteristics at some later time can preclude oscillation if they cause the loop gain, $|A||F_{vv}|$, to drop below unity. Such changes can easily be caused by, for example, changes in temperature or aging of components. The designer must therefore choose a compromise value of $|A||F_{vv}|$ to realize low distortion, but reliable operation, as well.
If the oscillator is to operate at a single frequency, it may be possible to have our cake and eat it too by choosing the value of $|A|F_{vv}$ well above unity to achieve reliable operation and then purifying the oscillator output with a tuned filter, such as an LC resonant circuit. This solution is not very convenient if the oscillator must operate over a wide range of frequencies, however, because a band pass filter with a wide tuning range can be difficult to realize in practice.

As a final perspective on the Barkhausen condition, we note that when

$$AF_{vv} = 1$$

or

$$\frac{Z_L}{Z_L + Z_o} F_{vv} = 1$$

then our earlier result for the voltage gain, $G_v$, with feedback,

$$G_v = \frac{A}{1 - AF_{vv}} \left\{ \frac{Z_I + Z_{th,dr}}{1 - \frac{Z_L}{Z_I + Z_L} AF_{vv}} \right\}$$

is infinite because the denominator in the numerator of the curly brackets is zero. In a naïve sense, then, we can say that the gain with feedback becomes infinite when the Barkhausen criterion is satisfied. The naïve perspective, then, is that oscillation corresponds to infinite gain with feedback.

This perspective is not particularly useful, except that it emphasizes that the feedback in sinusoidal oscillators is positive and, thereby, increases the gain of the amplifier instead of decreasing it, as negative feedback does. It is interesting to note, however, that positive feedback need not produce oscillations. If the feedback is positive, but $|A|F_{vv} < 1$, then oscillations die out and are not sustained. In this regime, the gain of the amplifier can be increased considerably by positive feedback. Historically, Edwin H. Armstrong, the person who first understood the importance of DeForest's vacuum triode as a dependent or controlled source, used positive feedback to obtain more gain from a single, costly, vacuum triode in a high frequency amplifier before Black applied negative feedback to audio amplifiers. As vacuum triodes became more readily available at reasonable cost, however, the use of positive feedback to obtain increased gain fell out of favor because the increased gain it produced was accompanied by enhanced noise in the output, in much the same way that the decreased gain produced by negative feedback was accompanied by reduced noise in the amplifier output. In practice, you got better results at a reasonable cost by using amplifiers with negative feedback, even though they required more vacuum triodes than would be necessary with positive feedback.
We now analyze a variety of sinusoidal oscillator circuits in detail to determine the frequency of possible oscillation and the condition on circuit components necessary to achieve slightly more than unity loop gain and, thereby, useful oscillations.

6.2 Phase shift Oscillator

We begin our consideration of practical oscillators with the phase shift oscillator, one that conforms fairly closely to our idealized model of sinusoidal oscillators.

The phase shift oscillator satisfies Barkhausen condition with an angle of $2\pi$. The inverting amplifier provides a phase shift of $\pi$. The three identical $RC$ sections (recall that the inverting input to the operational amplifier is a virtual ground so that $V_{\text{in}} \approx 0$) each provide an additional phase shift of $\pi / 3$ at the frequency of oscillation so that the phase shift around the loop totals to $2\pi$.

We begin the analysis by using the usual result for an inverting opamp configuration to express the output voltage, $V_{out}$, in terms of the input voltage, $V_{in}$, to the inverting amplifier:

$$V_{out} = -\frac{R_F}{R} V_{in} = -AV_{in}$$

where

$$A = \frac{R_F}{R}$$

Next, we write node equations to find the output of the feedback network in terms of the input to the feedback network. Oddly enough, the figure shows that the input to the feedback network is $V_{out}$ and that the output of the feedback network is $V_{in}$. To achieve a modest increase in notational simplicity, we use Laplace transform notation, although we will neglect transients and eventually substitute $s = j\omega$ and specialize to phasor analysis because we are interested only in the steady-state sinusoidal behavior of the circuit.
(1) \[ \left[ V_i(s) - V_{out}(s) \right] C_s + \frac{V_i(s)}{R} + \left[ V_i(s) - V_2(s) \right] C_s = 0 \]

(2) \[ \left[ V_2(s) - V_1(s) \right] C_s + \frac{V_2(s)}{R} + \left[ V_2(s) - V_{in}(s) \right] C_s = 0 \]

(3) \[ \left[ V_{in}(s) - V_2(s) \right] C_s + \frac{V_{in}(s)}{R} = 0 \]

Collecting terms, we find:

(1)' \[ 2 Cs + \frac{1}{R} V_i(s) + [ -Cs ] V_2(s) + [0] V_{in}(s) = s C \ V_{out}(s) \]

(2)' \[ [ -Cs ] V_i(s) + \left[ 2 Cs + \frac{1}{R} \right] V_2(s) + [ -Cs ] V_{in}(s) = 0 \]

(3)' \[ [0] V_i(s) + [ -Cs ] V_2(s) + \left[ Cs + \frac{1}{R} \right] V_{in}(s) = 0 \]

In matrix form,

\[
\begin{bmatrix}
2Cs + \frac{1}{R} & -Cs & 0 \\
-Cs & 2Cs + \frac{1}{R} & -Cs \\
0 & -Cs & Cs + \frac{1}{R}
\end{bmatrix}
\begin{bmatrix}
V_i(s) \\
V_2(s) \\
V_{in}(s)
\end{bmatrix}
= 
\begin{bmatrix}
sC \ V_{out}(s) \\
0 \\
0
\end{bmatrix}
\]

We calculate Cramer's delta as a step towards calculating the output of the feedback network, \( V_{in}(s) \), in terms of the input to the feedback network, \( V_{out}(s) \).

\[
\Delta = \begin{vmatrix}
2Cs + \frac{1}{R} & -Cs & 0 \\
-Cs & 2Cs + \frac{1}{R} & -Cs \\
0 & -Cs & Cs + \frac{1}{R}
\end{vmatrix}
\]
\[ \Delta = \left( 2Cs + \frac{1}{R} \right) \begin{bmatrix} 2Cs + \frac{1}{R} & -Cs & 0 \\ -Cs & Cs + \frac{1}{R} & -Cs \\ 0 & Cs & Cs + \frac{1}{R} \end{bmatrix} \]

\[ \Delta = \left( 2Cs + \frac{1}{R} \right) \left[ \left( 2Cs + \frac{1}{R} \right) \left( Cs + \frac{1}{R} \right) - (-Cs)^2 \right] + Cs \left[ -Cs \left( Cs + \frac{1}{R} \right) \right] \]

\[ \Delta = \left( 2Cs + \frac{1}{R} \right) \left[ 2(Cs)^2 + \frac{3}{R} Cs + \frac{1}{R^2} - (Cs)^2 \right] - (Cs)^3 - \frac{(Cs)^2}{R} \]

\[ \Delta = 2(Cs)^3 + \frac{6}{R} (Cs)^2 + \frac{2}{R^2} (Cs) + \frac{1}{R} (Cs)^2 + \frac{3}{R^2} (Cs) + \frac{1}{R^3} - (Cs)^3 - \frac{1}{R} (Cs)^2 \]

\[ \Delta = (Cs)^3 + \frac{6}{R} (Cs)^2 + \frac{5}{R^2} (Cs) + \frac{1}{R^3} \]

We now use this result in Cramer's rule to solve our set of equations for \( V_{in}(s) \) in terms of \( V_{out}(s) \).

\[ V_{in}(s) = \frac{1}{\Delta} \begin{bmatrix} 2Cs + \frac{1}{R} & -Cs & sCV_{out}(s) \\ -Cs & 2Cs + \frac{1}{R} & 0 \\ 0 & -Cs & 0 \end{bmatrix} \]

\[ V_{in}(s) = sCV_{out}(s) \frac{1}{\Delta} \begin{bmatrix} -Cs & 2Cs + \frac{1}{R} \\ 0 & -Cs \end{bmatrix} \]

\[ V_{in}(s) = \frac{(Cs)^3}{(Cs)^3 + \frac{6}{R} (Cs)^2 + \frac{5}{R^2} (Cs) + \frac{1}{R^3}} V_{out}(s) \]

\[ V_{in}(s) = \frac{s^3}{s^3 + \frac{6}{(RC)}s^2 + \frac{5}{(RC)^2}s + \frac{1}{(RC)^3}} V_{out}(s) \]

For an alternative solution with MATLAB, enter the following commands:

```matlab
syms s R C V1 V2 Vin V Vout
```
A=[2*C*s+1/R -C*s 0; -C*s 2*C*s+1/R -C*s; 0 -C*s C*s+1/R];

b=[s*C*Vout; 0; 0];

V=Ab;

Vin=V(3)

The result is:

\[ \text{Vin} = \frac{C^3 s^3 R^3 Vout}{C^3 s^3 R^3 + 6 C^2 s^2 R^2 + 5 C s R + 1} \]

That is,

\[ V_{in}(s) = \frac{C^3 s R^3}{C^3 s^3 R^3 + 6 C^2 s^2 R^2 + 5 C s R + 1} V_{out}(s) \]

\[ V_{in}(s) = \frac{s^3}{s^3 + \frac{6}{RC} s^2 + \frac{5}{(RC)^2} s + \frac{1}{(RC)^3}} V_{out}(s) \]

This expression, recall, gives the output of the feedback network, \( V_{in}(s) \), in terms of the input to the feedback network, \( V_{out}(s) \). Recall, also, that in phasor notation, the output of the amplifier, \( V_{out} \), in terms of the input to the amplifier, \( V_{in} \), is given by

\[ V_{out} = -AV_{in} \]

In Laplace transform notation, this equation becomes

\[ V_{out}(s) = -AV_{in}(s) \]

If we eliminate \( V_{in}(s) \) between the input/output equations for the amplifier and for the feedback network, we find:

\[ \frac{1}{A} V_{out}(s) = \frac{s^3}{s^3 + \frac{6}{RC} s^2 + \frac{5}{(RC)^2} s + \frac{1}{(RC)^3}} V_{out}(s) \]

Because \( V_{out}(s) \neq 0 \) if the oscillator is to provide useful output, we must require
\[
1 = \frac{-As^3}{s^3 + \frac{6}{(RC)}s^2 + \frac{5}{(RC)^2}s + \frac{1}{(RC)^3}}
\]

This required consistency condition is tantamount to the Barkhausen condition for oscillation.

\[
s^3 + \frac{6}{(RC)}s^2 + \frac{5}{(RC)^2}s + \frac{1}{(RC)^3} = -As^3
\]

\[
(A+1)s^3 + \frac{6}{(RC)}s^2 + \frac{5}{(RC)^2}s + \frac{1}{(RC)^3} = 0
\]

Because we are interested in the sinusoidal steady state, we specialize to phasor analysis by substituting \( s = j\omega \):

\[
-(A+1)j\omega^3 + \frac{6}{(RC)}(-\omega)^2 + \frac{5}{(RC)^2}j\omega + \frac{1}{(RC)^3} = 0
\]

\[
\left\{ -\frac{6}{(RC)}\omega^2 + \frac{1}{(RC)^3} + j\omega \left[ -(A+1)\omega^2 + \frac{5}{(RC)} \right] \right\} = 0
\]

This complex equation is in rectangular form. We obtain two separate equations by setting the real and imaginary parts to zero, separately. First, let’s set the real part of the equation to zero:

\[
-\frac{6}{(RC)}\omega^2 + \frac{1}{(RC)^3} = 0
\]

\[
-6\omega^2 + \frac{1}{(RC)^2} = 0
\]

\[
\omega^2 = \frac{1}{6(\text{RC})^2}
\]

\[
\omega = \frac{1}{\sqrt{6RC}}
\]

Thus, the frequency of the possible oscillations is determined by the values of the components in the feedback network. Whether or not oscillations will actually occur depends upon whether or not the second equation we obtain from the equation above is satisfied. To obtain this equation, we set the imaginary part of the equation to zero:
\[-(A+1)\omega^2 + \frac{5}{(RC)^2} = 0\]

\[(A+1)\omega^2 = \frac{5}{(RC)^2}\]

But, of course, we have already discovered the only possible value of \(\omega\):

\[\omega = \frac{1}{\sqrt{6RC}}\]

If we use this value in our equation, we find:

\[(A+1)\frac{1}{6(RC)^2} = \frac{5}{(RC)^2}\]

\[(A+1) = 30\]

Thus, we require

\[A = 29\]

to realize a loop gain through the amplifier and the feedback network of unity. In a practical oscillator, of course, \(A\) should be larger so that the oscillations will build up from noise. Perhaps it should be chosen to lie in the low 30’s to ensure reliable operation without too much distortion of the essentially sinusoidal output voltage. Recall that the gain-bandwidth product for a 741 operational amplifier is roughly 1 MHz. If we were to use it to realize an inverting amplifier with a gain of 30-something, then the bandwidth of the resulting amplifier would be no more than about 30 kHz. Thus, a 741 operational amplifier can be used to realize a phase shift oscillator in the audio range, but not much higher. A 2N3904 BJT, in contrast, has a gain-bandwidth product of several hundred MHz and could be used to realize an amplifier with a gain of 30-something with a much larger bandwidth. Some details of the feedback circuit would change, but the arrangement would be similar. Even if we use devices with larger gain-bandwidth product, the output Thevenin resistance of the amplifier limits the maximum frequency at which the phase shift oscillator is useful in practice. The impedance of the phase shift feedback network at the oscillation frequency should be much larger than the Thevenin output impedance so that the amplifier output will not be unduly loaded and so that our simple theory will apply. As a consequence, the minimum resistance of the resistors with value, \(R\), in the phase shift feedback network typically cannot be less than about 1000\(\Omega\). The minimum value of the capacitors, \(C\), must be much larger than stray or parasitic capacitances in the circuit and hence typically should be no smaller than about 1000 \(pF\). From the result

\[\omega = \frac{1}{\sqrt{6RC}}\]
we see, therefore, that phase shift oscillators are seldom useful at frequencies above 500kHz, regardless of the GBW of the active device in the amplifier. Note that it is inconvenient to change the frequency of phase shift oscillators because the values of a minimum of three resistors or three capacitors must be changed simultaneously.

### 6.3 Band Pass Oscillators

If the output of a band pass amplifier is fed back to its input, it may oscillate at a frequency within its pass band. If the pass band is narrow, the frequency will occur near the center frequency of the amplifier pass band. To the extent that the band pass filter is effective in attenuating signals with frequencies outside its pass band, it suppresses the harmonic content (distortion) in the output waveform that results as the oscillations grow to the point that they are limited by nonlinearities in the amplifier. As a consequence of this harmonic suppression, band pass oscillators can provide more nearly sinusoidal outputs than other types of practical oscillators.

To consider the requirements for oscillation in detail, recall that for a second order band pass filter that

\[ V_{out}(s) = T(s) V_{in}(s) \]

where

\[ T(s) = \frac{V_{out}(s)}{V_{in}(s)} = \frac{n_1 s}{s^2 + \left(\frac{\omega_0}{Q}\right)s + \omega_0^2} \]

Suppose that we connect the output directly to the input so that

\[ v_{out}(t) \equiv v_{in}(t) \]

Our feedback network in this case is thus a simple piece of wire. In the Laplace transform domain,

\[ V_{out}(s) \equiv V_{in}(s) \]

so that

\[ T(s) = 1 \]

This condition, tantamount to the Barkhausen condition, requires

\[ T(s) = \frac{n_1 s}{s^2 + \left(\frac{\omega_0}{Q}\right)s + \omega_0^2} = 1 \]
For the steady state sinusoidal case that describes the operation of the circuit after transients have died out, we substitute \( s = j \omega \) and obtain

\[
T(j\omega) = \frac{n_1(j\omega)}{(j\omega)^2 + \left(\frac{\omega_n}{Q}\right)(j\omega) + \omega_n^2} = \frac{j\omega n_1}{-\omega^2 + j\omega \left(\frac{\omega_n}{Q}\right) + \omega_n^2} = 1
\]

or

\[
j\omega n_1 = -\omega^2 + j\omega \left(\frac{\omega_n}{Q}\right) + \omega_n^2
\]

\[
\omega^2 + j\omega \left(n_1 - \frac{\omega_n}{Q}\right) - \omega_n^2 = 0
\]

This complex equation, of course, gives two real equations. The real part gives

\[
\omega^2 - \omega_n^2 = 0
\]

so that the frequency of oscillations, if any, must be

\[\omega = \omega_n\]

The imaginary part of the equation gives

\[
j\omega \left(n_1 - \frac{\omega_n}{Q}\right) = 0
\]

or

\[
n_1 = \frac{\omega_n}{Q}
\]

This equation is a requirement on the gain of the band pass amplifier.

As an example, recall the state variable filter circuit
which provides a band passed output of

\[
V_{BP}(s) = \frac{\omega_o s}{s^2 + \left(\frac{\omega_o}{Q}\right)s + \omega_o^2} V_{in}(s)
\]

where

\[
\omega_o = \frac{1}{RC}
\]

\[
Q = \frac{1}{3} \left(1 + \frac{R_F}{R}\right)
\]

If we wait until after the transients die out, we can substitute \(s = j\omega\) and set \(V_{BP} \equiv V_{in}\) to obtain

\[
1 = \frac{\omega_o (j\omega)}{(j\omega)^2 + \left(\frac{\omega_o}{Q}\right)(j\omega) + \omega_o^2}
\]

or

\[
1 = \frac{j\omega \omega_o}{\omega^2 + j\omega \left(\frac{\omega_o}{Q}\right) + \omega_o^2}
\]

\[-\omega^2 - j\omega \omega_o \left(1 - \frac{1}{Q}\right) + \omega_o^2 = 0
\]

The real part of the equation requires
\[ \omega = \omega_o = \frac{1}{RC} \]

and the imaginary part requires

\[ 1 = Q = \frac{1}{3} \left( 1 + \frac{R_F}{R} \right) \]

or

\[ 3 = \left( 1 + \frac{R_F}{R} \right) \]

or

\[ \frac{R_F}{R} = 2 \]

\[ R_F = 2R \]

Note that this requirement means that 1/3 of the band passed output is fed back to the input. To make the loop gain slightly greater than unity so that the oscillations will build up from noise, we need to feed back a larger fraction of the band passed output. To achieve this result, note that we should choose, in practice, \( R_F < 2R \).

We noted earlier that band pass oscillators offer the advantage of built-in harmonic suppression from its band pass filter to purify their output waveform. We note that the state variable circuit offers additional harmonic suppression if we take the output from the low pass output rather than from the band bass output. The state variable circuit has the disadvantage of greater power requirements than circuits with only one operational amplifier. With the operational amplifier state variable implementation, band pass filters are limited basically to the audio frequency range. Changing the frequency of oscillations in the state variable band pass oscillator requires changing the value of 2 capacitors simultaneously, only slightly more convenient that for the phase shift oscillator's 3 capacitors (or resistors).

### 6.4 Wien Bridge Oscillator

The Wien Bridge oscillator is a band pass oscillator that requires only one operational amplifier:
With a minor abuse of conventional notation, we represent impedances in this figure by resistances. We assume that the amplifier is non-inverting and has infinite input impedance and a real gain of \( A \) at the frequency of oscillation. For specificity, we show an operational amplifier configuration for which we saw, earlier, the gain is

\[
A = 1 + \frac{R_f}{R}
\]

but so our results will be more generally applicable, we express our results in terms of the open circuit gain, \( A \), which might be provided by a BJT or FET amplifier. As before, we assume that we are interested in the sinusoidal steady state response so that we can neglect transients.

We begin our analysis by writing the following node equation:

\[
\frac{V_{in}(s) - V_{out}(s)}{Z_1(s) + \frac{V_{in}(s)}{Z_2(s)}} = 0
\]

But

\[
V_{out}(s) = AV_{in}(s)
\]
Thus, we can write the node equation as

\[
\frac{AV_{in}(s) - AV_{out}(s)}{Z_1(s)} + \frac{AV_{in}(s)}{Z_2(s)} = 0
\]

\[
\frac{V_{out}(s) - AV_{out}(s)}{Z_1(s)} + \frac{AV_{out}(s)}{Z_2(s)} = 0
\]

\[
\left\{\frac{1 - A}{Z_1(s)} + \frac{1}{Z_2(s)}\right\}V_{out}(s) = 0
\]

Recall that

\[
Z_1(s) = R + \frac{1}{sC}
\]

\[
Z_2(s) = R \left\{\frac{1}{sC} = \frac{R}{sC}\right\} + \frac{1}{sC}
\]

Thus,

\[
\left\{\frac{1 - A}{R + \frac{1}{sC}} + sC + \frac{1}{R}\right\}V_{out}(s) = 0
\]

To investigate the sinusoidal steady state, we convert this equation to phasor notation by substituting \(s = j\omega\):

\[
\left\{\frac{1 - A}{R + \frac{1}{j\omega C}} + j\omega C + \frac{1}{R}\right\}V_{out} = 0
\]

Multiply through by \(R + \frac{1}{j\omega C}\):

\[
\left\{1 - A + \left(j\omega C + \frac{1}{R}\right)\left(R + \frac{1}{j\omega C}\right)\right\}V_{out} = 0
\]
If the oscillator is to produce useful output, then $V_{out} \neq 0$ and the curly brackets must be zero. The real and imaginary parts of the curly brackets must be zero independently. If we set the imaginary part to zero, we find:

$$j\left(\omega RC - \frac{1}{\omega RC}\right) = 0$$

or

$$\omega^2 = \frac{1}{(RC)^2}$$

Thus, the frequency of possible oscillation is given

$$\omega = \frac{1}{RC}$$

Note that if the amplification, $A$, is provided by a BJT or an FET, the frequency of oscillation can be extended beyond the audio range, just as with the phase shift oscillator. Because of the absence of the factor $\sqrt{6}$ in the denominator, however, the Wien Bridge oscillator can achieve more than twice the frequency of a phase shift oscillator, in practice.

By setting the real part of the equation to zero, we obtain:

$$A = 3$$

In practice, of course, we would choose $A$ to be slightly larger so that the oscillations can build up from noise. One of the advantages of the Wien Bridge oscillator is that it requires only a modest gain from the amplifier.

### 6.5 Colpitts and Hartley Oscillators

The Colpitts and Hartley oscillators are band pass filter oscillators in which the band pass filters are $LC$ resonant circuits. In the Colpitts oscillator, a capacitive voltage divider, which also serves as the capacitance part of the $LC$ resonant circuit, feeds back a portion of the output back into the input. In the Hartley oscillator, an inductive voltage divider, which also serves as the inductance part of the $LC$ resonant circuit, feeds back a portion of the output back into the input. With appropriate amplifiers, these
configurations can oscillate at frequencies up to a few hundred megahertz, much higher than the configurations we have considered so far.

Because these two oscillator configurations are identical topologically, we can perform much of the analysis of them simultaneously by considering the following circuit:

With a minor abuse of conventional notation, we represent impedances in this figure by resistances. We assume that the amplifier is non-inverting and has infinite input impedance, a real gain of $A$ at the frequency of oscillation and a Thevenin output resistance of $R_o$. For specificity, we show an operational amplifier configuration for which we saw, earlier, the gain is

$$A = 1 + \frac{R_f}{R}$$

but so our results will be more generally applicable, we express our results in terms of the open circuit gain, $A$, which might be provided by a BJT or FET amplifier that permits operation at high frequencies. As before, we assume that we are interested in the sinusoidal steady state response so that we can neglect transients.

Using Laplace transform notation, we have for a Colpitts oscillator that

$$Z_1(s) = \frac{1}{sC_1}$$

$$Z_2(s) = \frac{1}{sC_2}$$

$$Z_3(s) = sL$$
while for a Hartley oscillator,

\[ Z_1(s) = sL_1 \]
\[ Z_2(s) = sL_4 \]
\[ Z_3(s) = \frac{1}{sC} \]

We can easily write node equations that hold for both oscillators:

\[
(1) \quad \frac{V_{in}(s) - V_{out}(s)}{Z_1(s)} + \frac{V_{in}(s)}{Z_2(s)} = 0
\]

\[
(2) \quad \frac{V_{out}(s) - V_{in}(s)}{Z_1(s)} + \frac{V_{out}(s)}{Z_3(s)} + \frac{V_{out}(s) - AV_{in}(s)}{R_o} = 0
\]

Collecting terms, we find:

\[
(1)' \quad V_{in}(s) \left[ \frac{1}{Z_1(s)} + \frac{1}{Z_2(s)} \right] + V_{out}(s) \left[ -\frac{1}{Z_1(s)} \right] = 0
\]

\[
(2)' \quad V_{in}(s) \left[ -\frac{1}{Z_1(s)} - \frac{A}{R_o} \right] + V_{out}(s) \left[ \frac{1}{Z_1(s)} + \frac{1}{Z_3(s)} + \frac{1}{R_o} \right] = 0
\]

From equation (1)', we find:

\[
(1)'' \quad V_{in}(s) = \frac{V_{out}(s)}{\frac{1}{Z_1(s)} + \frac{1}{Z_2(s)}}
\]

We substitute equation (1)'' into (2)'

\[
\left\{ \frac{1}{Z_1(s)} + \frac{1}{Z_2(s)} \right\} \left[ -\frac{1}{Z_1(s)} - \frac{A}{R_o} \right] + \left[ \frac{1}{Z_1(s)} + \frac{1}{Z_3(s)} + \frac{1}{R_o} \right] V_{out}(s) = 0
\]
We divide through by $\frac{1}{Z_1(s)} + \frac{1}{Z_2(s)}$:

$$
\left\{ \frac{1}{Z_2(s)} - \frac{A}{R_o} + \frac{1}{Z_3(s)} + \frac{1}{Z_2(s)} \right\} V_{out}(s) = 0
$$

$$
\left\{ \frac{1}{Z_2(s)} + \frac{1}{Z_3(s)} + \frac{1}{Z_2(s)} \right\} V_{out}(s) = 0
$$

We now specialize this equation for the Colpitts oscillator, for which, recall,

$$
Z_1(s) = \frac{1}{sC_1}
$$

$$
Z_2(s) = \frac{1}{sC_2}
$$

$$
Z_3(s) = sL
$$

From these results, we obtain the following equation for the Colpitts oscillator:
To investigate the sinusoidal steady state, we convert this equation to phasor notation by substituting $s = j\omega$:

$$\left\{ j\omega C_2 + \left( 1 + \frac{C_2}{C_1} \right) \frac{1}{j\omega L} + \frac{1}{R_o} \left( -A + 1 + \frac{C_2}{C_1} \right) \right\} V_{out}(s) = 0$$

If the oscillator is to produce useful output, then $V_{out} \neq 0$ and the curly brackets must be zero. The real and imaginary parts of the curly brackets must be zero independently. If we set the imaginary part to zero, we find:

$$j\omega C_2 + \left( 1 + \frac{C_2}{C_1} \right) \frac{1}{j\omega L} = 0$$

$$\omega C_2 - \left( 1 + \frac{C_2}{C_1} \right) \frac{1}{\omega L} = 0$$

$$\omega^2 = \left( \frac{1}{C_1} + \frac{1}{C_2} \right) \frac{1}{L}$$

Thus, the frequency of possible oscillation is given

$$\omega = \sqrt{\left( \frac{1}{C_1} + \frac{1}{C_2} \right) \frac{1}{L}}$$

Because $C_1$ and $C_2$ are connected in series, this frequency is simply the resonant frequency of the $LC$ circuit, that is, the center of the pass band.

By setting the real part of the equation to zero, we obtain:

$$-A + 1 + \frac{C_2}{C_1} = 0$$

or

$$A = 1 + \frac{C_2}{C_1}$$
In practice, of course, we would choose $A$ to be slightly larger so that the oscillations can build up from noise. Because the designer can set the ratio of $C_1$ and $C_2$ to a convenient value, the gain of the amplifier, $A$, need not be especially large, a potential advantage over the other circuits we have investigated so far. In practice, it is difficult to vary the frequency by changing the values of the capacitors because their ratio should remain constant. In practical variable frequency Colpitts oscillators, therefore, the frequency is sometimes varied by partially inserting and withdrawing a low-loss ferrite core located within an inductive winding. Notice that neither the frequency nor the gain requirement for the Colpitts oscillator depend upon $R_o$, the output Thevenin resistance of the amplifier.

We can see more clearly the physical meaning of the equation that specifies the minimum amplitude if we write it as follows:

$$A = 1 + \frac{C_2}{C_1} = \frac{C_1 + C_2}{C_1}$$

or

$$1 = A \frac{C_1}{C_1 + C_2} = A \frac{1}{C_1} + \frac{1}{C_2} = A \frac{1}{j\omega C_1} + \frac{1}{j\omega C_2}$$

The quantity $\frac{1}{j\omega C_2}$ is the fraction of the output that is fed back into the input of the amplifier in a Colpitts oscillator. Thus, the real part of the equation simply requires that the loop gain through the amplifier and feedback network be unity. The imaginary part of the equation requires the phase shift around the loop to be a multiply of $2\pi$.

For a Hartley oscillator, recall that

$$Z_1(s) = sL_1$$
$$Z_2(s) = sL_4$$
$$Z_3(s) = \frac{1}{sC}$$

Thus, the equation

$$\left[ \frac{1}{Z_2(s)} + \left( 1 + \frac{Z_1(s)}{Z_2(s)} \right) \frac{1}{Z_3(s)} + \frac{1}{R_o} \left( -A + 1 + \frac{Z_1(s)}{Z_2(s)} \right) \right] V_{out}(s) = 0$$
becomes
\[
\left[ \frac{1}{sL_2} + \left( 1 + \frac{L_1}{L_2} \right)Cs + \frac{1}{R_0} \left( -A + 1 + \frac{L_1}{L_2} \right) \right] V_{out}(s) = 0
\]

To investigate the sinusoidal steady state, we convert this equation to phasor notation by substituting \( s = j\omega \):
\[
\left[ \frac{1}{j\omega L_2} + \left( 1 + \frac{L_1}{L_2} \right)j\omega C + \frac{1}{R_0} \left( -A + 1 + \frac{L_1}{L_2} \right) \right] V_{out} = 0
\]

If the oscillator is to produce useful output, then \( V_{out} \neq 0 \) and the curly brackets must be zero. The real and imaginary parts of the curly brackets must be zero independently. If we set the imaginary part to zero, we find:
\[
\frac{1}{j\omega L_2} + \left( 1 + \frac{L_1}{L_2} \right)j\omega C = 0
\]
\[
- \frac{1}{\omega L_2} + \left( 1 + \frac{L_1}{L_2} \right)\omega C = 0
\]
\[
-1 + \omega^2 CL_2 \left( 1 + \frac{L_1}{L_2} \right) = 0
\]
\[
\omega^2 C(L_2 + L_1) = 1
\]

Thus, the frequency of possible oscillation is given
\[
\omega = \frac{1}{\sqrt{(L_1 + L_2)C}}
\]

Because \( L_1 \) and \( L_2 \) are connected in series, this frequency is simply the resonant frequency of the \( LC \) circuit, that is, the center of the pass band.

By setting the real part of the equation to zero, we obtain:
\[
-A + 1 + \frac{L_1}{L_2} = 0
\]

or
In practice, of course, we would choose $A$ to be slightly larger so that the oscillations can build up from noise. Because the designer can set the ratio of $L_1$ and $L_2$ to a convenient value, the gain of the amplifier, $A$, need not be especially large, a potential advantage shared with the Colpitts oscillator. In practice, it is difficult to vary the frequency by changing the values of the inductors because their ratio should remain constant. In practical variable frequency Hartley oscillators, therefore, the frequency usually is varied by adjusting the capacitance, $C$.

We can see more clearly the physical meaning of the equation that specifies the minimum amplitude if we write it as follows:

$$A = 1 + \frac{L_1}{L_2}$$

or

$$1 = A \frac{L_2}{L_1 + L_2} = A \frac{j\omega L_2}{j\omega L_1 + j\omega L_2}$$

The quantity $\frac{j\omega L_2}{j\omega L_1 + j\omega L_2}$ is the fraction of the output that is fed back into the input of the amplifier in a Hartley oscillator. Thus, the real part of the equation simply requires that the loop gain through the amplifier and feedback network be unity. The imaginary part of the equation requires the phase shift around the loop to be a multiple of $2\pi$.

### 6.6 Piezoelectric Crystal Oscillators

When some materials are placed between conducting plates and subjected to mechanical compression, they produce an internal electric field that causes a voltage to appear between the conducting plates. A voltage also appears between the plates if the materials are subjected to mechanical tension, although the polarity of the voltage produced is opposite to that produced by compression. If the sample is subjected to neither compression nor tension, no voltage appears between the plates. Conversely, application of a voltage between the plates produces compression or tension in the material, depending on the polarity of the voltage applied. This electromechanical behavior is called the piezoelectric effect.

Excitation of high frequency mechanical vibrations in a small slab of a piezoelectric material, such as crystalline quartz, produces a damped oscillatory voltage across conducting electrodes placed on opposite faces of the material similar to that produced by an excited $LC$ resonant circuit. Indeed, the electrical behavior of a small piece of piezoelectric material placed between conducting electrodes can be modeled by the following circuit:
where the upper and lower terminals connect to the conducting electrodes attached to opposite faces of the piezoelectric material. Given this equivalent circuit, it is not hard to show that a piezoelectric crystal can form the heart of a band pass filter, and hence, the basis of a band pass oscillator.

Although piezoelectric crystal oscillators can oscillate at frequencies as low as $10 \text{kHz}$, they typically oscillate at frequencies between 1 and $10 \text{MHz}$. Their frequency range can be extended to frequencies up to a few hundred megahertz by means of special tricks. Like Colpitts and Hartley oscillators, therefore, piezoelectric crystal oscillators can operate at much higher frequencies than the various $RC$ oscillators that we considered earlier.

In addition to high frequency operation, piezoelectric crystal oscillators offer two main features, one recently important and one long important. The unrelenting trend toward miniaturization in contemporary electronics has made the capacitors and inductors required for high frequency band pass oscillators begin to seem huge and cumbersome. For typical frequencies of oscillation, a piezoelectric crystal in a practical oscillator will occupy less than $\frac{3}{100} \text{mm}^3$, hundreds of times less than the volume necessary for the coil and capacitor in Colpitts or Hartley oscillators. Since the early days of electronics, piezoelectric crystal oscillators have been known for offering incomparable frequency stability, a feature perhaps more important today than ever before.

A rather peculiar feature of piezoelectric crystal oscillators is that the crystal can vibrate mechanically not only at its fundamental frequency, but at harmonics of that frequency, as well. This phenomenon is analogous to the fact that a taut string can vibrate at multiples of the lowest possible frequency of oscillation. Oscillation at these overtones of the fundamental frequency permits oscillators with piezoelectric crystals of reasonable size to operate at frequencies up to a few hundred MHz. In this respect, overtones provide a desirable feature. Overtone vibrations at harmonic frequencies, however, also mean that a piezoelectric crystal used as a band pass filter has pass bands at harmonics of the fundamental frequency as well as at the fundamental frequency itself. Consequently, harmonic suppression in a piezoelectric crystal band pass filter oscillator is not as effective as in Colpitts or Hartley oscillators. Therefore, it is usually necessary to pass the output of a piezoelectric crystal oscillator through an $LC$ band pass filter to achieve harmonic suppression comparable to the in Colpitts or Hartley oscillators. In practice, the pattern of resonant frequencies in piezoelectric crystals is actually even a little more complicated than we just described. The three-dimensional nature of the piezoelectric crystal permits it to
vibrate at more than one “fundamental” frequency, as well as the harmonics of each one of these. Thus, piezoelectric crystals can exhibit resonant frequencies that are not obviously harmonically related.

Historically, the major disadvantage of piezoelectric oscillators was inflexibility: they operate at a single fixed frequency. Today, however, phase-locked loops and digital technology have liberated piezoelectric crystal oscillators from the severe limitation of single frequency operation and made them widely useful.

A typical quality factor, \( Q \), \((2 \pi \text{ divided by the fraction of the oscillatory energy dissipated during each cycle of oscillation})\) for a piezoelectric crystal is a few hundred thousand, about 10,000 times larger than we can usually achieve with practical \( LC \) circuits. This high \( Q \) gives a piezoelectric band pass filter narrow bandwidth. Because

\[
\Delta f = \frac{f_o}{Q}
\]

the bandwidth, \( \Delta f \), of a piezoelectric crystal filter with a center frequency, \( f_o \), of \( 1 \text{MHz} \), for example, could be less than \( 10 \text{Hz} \), an impossible achievement for \( LC \) band pass filters, for which \( Q \) values of 10 or so are typical. In the equivalent circuit above, the inductor, \( L \), may have values of a few 100\( H \), the series capacitor, \( C_s \), may have values of a few tenths of a femtofarad (\( 10^{-15} F \)), and the series resistor, \( R_s \), may have values of a few tens of \( k\Omega \). The parallel capacitance, \( C_p \), results mainly from the dielectric properties of the piezoelectric material between the conducting electrodes, often plated directly on the piezoelectric material. Its value is typically a few picofarads.

With these values in mind, let’s look at the impedance, \( Z(s) \), of the piezoelectric crystal.

\[
Z(s) = \frac{1}{sC_p + \frac{1}{sL + \frac{1}{sC_s} + R_s}} = 1 + sC_p \left[ sL + \frac{1}{sC_s} + R_s \right]
\]

\[
Z(s) = \frac{s^2LC_s + 1 + sC_sR_s}{sC_s + s^2C_pC_sL + sC_p + s^3C_pC_sR_s}
\]

\[
Z(s) = \frac{LC_s}{sC_p + \frac{C_s}{C_p} + s^2C_sL + 1 + sC_sR_s}
\]
\[ Z(s) = \frac{1}{sC_p} \left( \frac{1}{LC_p} + s^2 + \frac{1}{LC_s} + \frac{s}{L} \right) \]

\[ Z(s) = \frac{1}{sC_p} \left( \frac{s^2 + R_s}{L} + s + \frac{1}{LC_s} \right) \]

We are interested in the sinusoidal steady state response, so we substitute \( s = j\omega \) and obtain

\[ Z(j\omega) = \frac{1}{j\omega C_p} \left( -\omega^2 + \frac{R_s}{L} j\omega + \omega_s^2 \right) \]

where

\[ \omega_s = \frac{1}{\sqrt{LC_s}} \]

and

\[ \omega_p = \sqrt{\frac{1}{L} \left( \frac{1}{C_p} + \frac{1}{C_s} \right)} \geq \omega_s \]

Because \( C_p \gg C_s \), note that \( \omega_p \) is only the slightest bit larger than \( \omega_s \). Thus, we find

\[ Z(j\omega) = -j \frac{1}{\omega C_p} \left( \omega^2 - \omega_s^2 - j\omega \frac{R_s}{L} \right) \equiv R(j\omega) + jX(j\omega) \]

where \( R(j\omega) \) is the resistance and \( X(j\omega) \) is the reactance of the piezoelectric crystal at angular frequency \( \omega \) and, recall, \( \omega_p^2 > \omega_s^2 \).
The dissipative resistance, \( R(j\omega) \) is positive and shows a peak at a frequency between \( \omega_s \) and \( \omega_p \), but is otherwise uninteresting for our present purposes. Here is a sketch of the reactance, \( X(j\omega) \), of the piezoelectric crystal vs. frequency:

Note that the reactance, \( X(j\omega) \), is positive only for frequencies, \( \omega \), in the range \( \omega_s < \omega < \omega_p \). That is, for frequencies between \( \omega_s \) and \( \omega_p \), the piezoelectric crystal behaves as an inductor. (Outside this range, it behaves as a capacitor.) Because \( \omega_s \) and \( \omega_p \) are nearly coincident, the crystal behaves as an inductor over only an extremely narrow range of frequencies. Because piezoelectric crystal oscillators are designed to rely on the effective inductance for their operation, the possible frequency of oscillation is limited to the extremely narrow frequency range over which the piezoelectric crystal behaves, indeed, as an inductor. Thus, the frequency of oscillation is necessarily extremely stable.

A simple example of a piezoelectric crystal oscillator is the Pierce oscillator. Consider the following FET realization:

The popularity of this circuit is doubtless due to its apparent simplicity – it is only necessary to add a piezoelectric crystal to a fairly standard FET amplifier to form the Pierce oscillator. If the Pierce oscillator works, it is, indeed, an extremely simple oscillator. During the course of our analysis, however, we will discover that best operation is achieved if some additional components are added.

In the circuit above, the capacitor \( C_s \) is assumed to be chosen so that \( \frac{1}{\omega C_s} \ll R_s \) at the frequency of oscillation. In this case, \( C_s \) can be considered to be a short circuit for signals. As a consequence, the negative feedback for signals that \( R_j \), otherwise would provide is eliminated and the voltage gain for the FET is higher than it would be without the presence of \( C_s \). Note, however, that \( R_j \) still provides negative
feedback necessary to achieve good bias stability. Similarly, we assume that $C_d$ is chosen so that its reactance at the frequency of oscillation is small in comparison to the Thevenin output resistance of the amplifier, $R_d$: \[
\frac{1}{\omega C_d} \ll R_d.\] Thus, $C_d$ also behaves as a short circuit for signals. With these assumptions in mind, we can draw the following signal equivalent circuit for the Pierce oscillator circuit:

![Pierce Oscillator Circuit Diagram]

where $g_m$ is the mutual transconductance of the FET and, for the moment, we have neglected to include the small parasitic capacitances associated with the FET so that the diagram is simpler. If we replace the piezoelectric crystal by the equivalent circuit that we considered before and add the parasitic capacitances associated with the FET, we obtain:

![Equivalent Circuit with Parasitics]

where $R_s = R_s || R_{gs}$. Note that $C_{gd}$, the gate to drain parasitic capacitance of the FET, parallels $C_p$, the capacitance between the electrodes of the piezoelectric crystal. For $\omega_s < \omega < \omega_p$, the inductance of the piezoelectric crystal, together with the gate-to-source capacitance, $C_{gs}$, and the drain-to-source capacitance, $C_{ds}$, form a variation of the Colpitts oscillator configuration. The variation is that the junction of the capacitive voltage divider formed by $C_{gs}$ and $C_{ds}$ is connected to ground whereas the output of the feedback network is taken from the end of the inductor that is opposite to the end connected...
to the drain of the FET. At the \( LC \) resonant frequency, the effect of this variation is to reverse the sign of the voltage fed back to the input of the FET amplifier in comparison with the Colpitts configuration that we considered earlier. This sign reversal is equivalent to a phase shift of \( \pi \) that, in addition to the phase shift of \( \pi \) produced by the signal inversion in the FET, gives a loop phase shift of \( 2\pi \), necessary to satisfy the phase part of the Barkhausen condition at resonance. Because \( C_{gs} \) and \( C_{ds} \) in practical devices are extremely small (picofards or smaller), the operation of the capacitive voltage divider can be unduly affected by stray capacitances external to the transistor package. As a consequence, it is usually prudent, as we will become clear during the course of the analysis, to place larger external capacitors in parallel with them to improve reliability of the oscillator performance.

To derive the Barkhausen conditions for the Pierce oscillator, we write node equations at the gate and drain terminals of the FET:

\[
(1) \quad V_{gs}(s) \left[ \frac{1}{R_g} + sC_{gs} \right] + \frac{V_{gs}(s) - V_{out}(s)}{Z(s)} = 0
\]

\[
(2) \quad V_{out}(s) \left[ \frac{1}{R_d} + sC_{ds} \right] + \frac{V_{out}(s) - V_{gs}(s)}{Z(s)} + s_m V_{gs}(s) = 0
\]

where, recall,

\[
Z(s) = \frac{1}{sC_p} \left( \frac{R_s}{s^2 + \frac{R_s}{L}s + \omega_p^2} \right)
\]

is the impedance of the piezoelectric crystal equivalent circuit. We rewrite equations (1) and equation (2) as a pair of pair of simultaneous linear algebraic equations with \( V_{gs}(s) \) and \( V_{out}(s) \) as unknowns:

\[
(1)' \quad V_{gs}(s) \left[ \frac{1}{R_g} + sC_{gs} + \frac{1}{Z(s)} \right] - V_{out}(s) \left[ \frac{1}{Z(s)} \right] = 0
\]

\[
(2)' \quad V_{gs}(s) \left[ g_m - \frac{1}{Z(s)} \right] + V_{out}(s) \left[ \frac{1}{R_d} + sC_{ds} + \frac{1}{Z(s)} \right] = 0
\]

From equation (1)' , we solve for \( V_{gs}(s) \) in terms of \( V_{out}(s) \) :
\[ V_{gs}(s) = \frac{1}{Z(s) + \frac{1}{R_g} + sC_{gs} + \frac{1}{Z(s)}} V_{out}(s) \]

We can use this result to eliminate \( V_{gs}(s) \) from equation (2)’:

\[
\left\{ \frac{1}{Z(s)} \left( g_m - \frac{1}{Z(s)} \right) \right\} + \frac{1}{R_d} + sC_{ds} + \frac{1}{Z(s)} V_{out}(s) = 0
\]

\[
\left\{ \frac{1}{Z(s)} \left( g_m - \frac{1}{Z(s)} \right) \right\} + \left( \frac{1}{R_d} + sC_{ds} \right) + \frac{1}{R_g} + sC_{gs} + \frac{1}{Z(s)} \right\} V_{out}(s) = 0
\]

\[
\left\{ g_m Z(s) - 1 \right\} + \left( \frac{Z(s)}{R_d} + sC_{ds} Z(s) + 1 \right) \left( \frac{Z(s)}{R_g} + sC_{gs} Z(s) + 1 \right) V_{out}(s) = 0
\]

\[
\left\{ g_m Z(s) - 1 \right\} + \left( \frac{Z(s)}{R_d} + sC_{ds} Z(s) \right) + \frac{Z(s)}{R_g} + sC_{gs} Z(s) + 1
\]

\[
\left\{ g_m + Z(s) \left( \frac{1}{R_d} + sC_{ds} \right) + \frac{1}{R_g} + sC_{gs} \right\} V_{out}(s) = 0
\]

\[
\left\{ g_m + Z(s) \left[ \frac{1}{R_d} + \frac{1}{R_g} + s^2 C_{ds} C_{gs} + s \left( \frac{1}{R_d} C_{gs} + \frac{1}{R_g} C_{ds} \right) \right] + \frac{1}{R_d} + \frac{1}{R_g} + s \left( C_{ds} + C_{gs} \right) \right\} V_{out}(s) = 0
\]
\[
\left\{
\begin{aligned}
g_m + Z(s) & \left[ \frac{1}{R_d} \frac{1}{R_g} + s^2 C_{ds} C_{gs} + s \left( \frac{1}{R_d} C_{gs} + \frac{1}{R_g} C_{ds} \right) \right] \\
& + \frac{1}{R_d} + \frac{1}{R_g} + s \left( C_{ds} + C_{gs} \right)
\end{aligned}
\right\} V_{\text{out}}(s) = 0
\]

Since we are interested in the sinusoidal steady state response, we switch to phasor notation by substituting \( s = j\omega \):

\[
\left\{
\begin{aligned}
g_m + Z(j\omega) & \left[ \frac{1}{R_d} \frac{1}{R_g} + (j\omega)^2 C_{ds} C_{gs} + (j\omega) \left( \frac{1}{R_d} C_{gs} + \frac{1}{R_g} C_{ds} \right) \right] \\
& + \frac{1}{R_d} + \frac{1}{R_g} + (j\omega) \left( C_{ds} + C_{gs} \right)
\end{aligned}
\right\} V_{\text{out}} = 0
\]

At this point, we take time out to discover a simple approximate form for,

\[
Z(j\omega) \equiv R(j\omega) + jX(j\omega) = -j \frac{1}{\omega(C_p + C_{gd})} \left( \omega^2 - \omega^2_p - j\omega \frac{R_s}{L} \right)
\]

where we have added \( C_{gd} \) to \( C_p \), as mentioned above. A typical value of \( \frac{R_s}{L} \sim \frac{10^4}{100} \sim 100 \) so that for frequencies greater than \( 10^5 \text{ rad/sec} \), the \( j\omega \frac{R_s}{L} \) terms in \( Z(j\omega) \) are, except for frequencies quite near \( \omega_p \), negligible in comparison to \( \omega^2 \) and to \( \omega^2_p \) and \( \omega^2 \). The main effect of the \( j\omega \frac{R_s}{L} \) term in the denominator is to limit the magnitude of the impedance to finite values at frequencies quite near \( \omega_p \). Thus, we use the approximate form

\[
Z(j\omega) \equiv R(j\omega) + jX(j\omega) = -j \frac{1}{\omega(C_p + C_{gd})} \left( \omega^2 - \omega^2_p \right)
\]

In this approximation, therefore, \( R(j\omega) = 0 \) and

\[
Z(j\omega) = jX(j\omega)
\]

where
Thus, we can write our consistency condition as:

$$
\begin{align*}
\left\{ \left. g_m + jX(j\omega) \left[ \frac{1}{R_d} \frac{1}{R_g} - \omega^2 C_{ds} C_{gs} + j\omega \left( \frac{1}{R_d} C_{gs} + \frac{1}{R_g} C_{ds} \right) \right] \right. \right. \\
+ \frac{1}{R_d} + \frac{1}{R_g} + j\omega \left( C_{ds} + C_{gs} \right) \\
\left. \left. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \right. \r...
\[
\frac{\omega^2 - \omega_s}{\omega_p^2 - \omega^2} \left[ \frac{1}{R_d} \frac{1}{R_g} - \omega^2 C_{ds} C_{gs} \right] + \omega^2 \left( C_{ds} + C_{gs} \right) \left( C_p + C_{gd} \right) = 0
\]

\[
\frac{\omega^2 - \omega_s^2}{\omega_p^2 - \omega^2} \left[ \frac{1}{R_d C_{ds}} \frac{1}{R_g C_{gs}} - \omega^2 \right] + \omega^2 \left( \frac{1}{C_{ds}} + \frac{1}{C_{gs}} \right) \left( C_p + C_{gd} \right) = 0
\]

For proper operation of the circuit, we need to choose values so that

\[
\frac{1}{R_d C_{ds}} \frac{1}{R_g C_{gs}} \ll \omega^2
\]

to minimize dependence of the oscillator frequency, as determined by the equation above, on the resistors \( R_d \) and \( R_g \). (We prefer the frequency to depend only on the piezoelectric crystal parameters.) During design, we can satisfy the inequality by choosing large values for \( R_d \) and \( R_g \) and/or by adding external supplemental capacitances in parallel with \( C_{ds} \) and \( C_{gs} \). Supplementing the values of \( C_{ds} \) and \( C_{gs} \) with sufficiently large fixed external capacitors also has the advantage of making the equation above for the frequency of the oscillator independent of all transistor parameters. Of course, we must be careful not to make design choices that will require unrealistic values of the transconductance, \( g_m \), to realize unity loop gain. We'll return to this issue later.

If the inequality above is satisfied, then the equation for the oscillator frequency becomes

\[
\frac{\omega^2 - \omega_s^2}{\omega_p^2 - \omega^2} = -\omega_s^2 + \omega^2 \left( \frac{1}{C_{ds}} + \frac{1}{C_{gs}} \right) \left( C_p + C_{gd} \right) = 0
\]

or

\[
\frac{\omega^2 - \omega_s^2}{\omega_p^2 - \omega^2} = \left( \frac{1}{C_{ds}} + \frac{1}{C_{gs}} \right) \left( C_p + C_{gd} \right) \equiv \alpha
\]

where

\[
\alpha = \left( \frac{1}{C_{ds}} + \frac{1}{C_{gs}} \right) \left( C_p + C_{gd} \right) > 0
\]

is a positive dimensionless constant. Proceeding, we find

\[
\omega^2 - \omega_s^2 = \alpha \left( \omega_p^2 - \omega^2 \right)
\]

\[
(1 + \alpha)\omega^2 = \omega_s^2 + \alpha \omega_p^2
\]
\[ \omega^2 = \frac{\omega_s^2 + \alpha \omega_p^2}{1 + \alpha} \]

Note that, regardless of the value of \( \alpha \), the frequency of oscillation is constrained to lie between \( \omega_s \) and \( \omega_p \):

\[ \omega_s^2 < \omega^2 < \omega_p^2 \]

Thus, the frequency of oscillation is forced, by the piezoelectric crystal, to lie within a very narrow range. Remember that it is only in this frequency range that the piezoelectric crystal behaves as an inductor, a component essential for the basic Colpitts to function as an oscillator.

We now return to the equation

\[
\left\{ \begin{array}{l}
g_m - \omega X(j\omega) \left( \frac{1}{R_d} C_{gs} + \frac{1}{R_g} C_{ds} \right) + \frac{1}{R_d} + \frac{1}{R_g} \\
+ jX(j\omega) \left( \frac{1}{R_d} \frac{1}{R_g} - \omega^2 C_{ds} C_{gs} \right) + j\omega \left( C_{ds} + C_{gs} \right) \end{array} \right\} V_{out} = 0
\]

and consider the consequences of the real part of its curly brackets being zero:

\[
g_m - \omega X(j\omega) \left( \frac{1}{R_d} C_{gs} + \frac{1}{R_g} C_{ds} \right) + \frac{1}{R_d} + \frac{1}{R_g} = 0
\]

If we use our earlier result that

\[
X(j\omega) \approx -\frac{1}{\omega(C_p + C_{gd})} \frac{\omega^2 - \omega_s^2}{\omega^2 - \omega_p^2}
\]

then we see

\[
g_m + \frac{1}{(C_p + C_{gd})} \frac{\omega^2 - \omega_s^2}{\omega^2 - \omega_p^2} \left( \frac{1}{R_d} C_{gs} + \frac{1}{R_g} C_{ds} \right) + \frac{1}{R_d} + \frac{1}{R_g} = 0
\]

Recall that

\[
g_m = \frac{C_{gs} C_{ds}}{(C_p + C_{gd})} \frac{\omega_s^2 - \omega_p^2}{\omega^2 - \omega_p^2} \left( \frac{1}{R_d} C_{ds} + \frac{1}{R_g} C_{gs} \right) - \frac{1}{R_d} - \frac{1}{R_g}
\]
\[
\frac{\omega^2 - \omega_p^2}{\omega_0^2 - \omega^2} = \left( \frac{1}{C_{ds}} + \frac{1}{C_{gs}} \right) \left( C_p + C_{gd} \right) \equiv \alpha
\]

Thus,

\[
g_m = \frac{C_{gs} C_{ds}}{(C_p + C_{gd})} \left( \frac{1}{C_{ds}} + \frac{1}{C_{gs}} \right) \left( C_p + C_{gd} \right) \left( \frac{1}{R_d C_{ds}} + \frac{1}{R_g C_{gs}} \right) - \frac{1}{R_d} - \frac{1}{R_g}
\]

\[
g_m = \frac{C_{ds} + C_{gs}}{R_d C_{ds}} \left( 1 + \frac{R_d C_{ds}}{R_g C_{gs}} \right) - \frac{1}{R_d} - \frac{1}{R_g}
\]

\[
g_m = \frac{1}{R_d} \left( 1 + \frac{C_{gs}}{C_{ds}} \right) \left( 1 + \frac{R_d C_{ds}}{R_g C_{gs}} \right) - \frac{1}{R_d} - \frac{1}{R_g}
\]

Because the drain in an FET is much further away from the source than is the gate, we find, in practice, \(C_{ds} \ll C_{gs}\). In practice, it is also true that the parallel combination of the gate resistors, \(R_g \gg R_d\), the drain resistor. With little error, then, we can write

\[
g_m = \frac{1}{R_d} \left( 1 + \frac{C_{gs}}{C_{ds}} \right) - \frac{1}{R_d}
\]

or

\[
g_m = \frac{1}{R_d} \frac{C_{gs}}{C_{ds}}
\]

This equation gives the critical minimum value of \(g_m\) necessary to achieve unity loop gain. To make the loop gain slightly larger to improve reliability in the operation of the oscillator, we should choose \(g_m\) to be larger than the minimum value given by the equation:

\[
g_m = \frac{1}{R_d} \frac{C_{gs}}{C_{ds}}
\]

Notice that the minimum transconductance, \(g_m\), depends only on the ratio of the capacitances. Note also that the required \(g_m\) can be reduced by increasing \(R_d\) and/or \(C_{ds}\). Increasing \(R_d\) increases the voltage gain of the FET amplifier, as we have seen. For a given operating point for the FET, however, increasing
$R_d$ means increasing the power supply voltage, which may not be easy to do. Increasing $C_{ds}$ increases the fraction of the output fed back into the input. Increasing $C_{ds}$ also helps us to satisfy the inequality

$$\frac{1}{R_g C_{ds}} \frac{1}{R_g C_{gs}} << \omega^2$$

Thus, it makes sense to add an external capacitor in parallel with $C_{ds}$ to decrease the $g_m$ required for oscillation and to reduce the sensitivity of the frequency of oscillation to values of the resistors $R_d$ and $R_g$. If it does not make satisfying the above inequality too difficult, we can add a supplemental external capacitor in parallel with $C_{gs}$ to make the frequency of oscillation essentially independent to the values of any parameters except those of the piezoelectric crystal, which are remarkable stable. If extreme frequency stability is required, the piezoelectric crystal can be placed in a controlled temperature oven to reduce even further the already slight variation of the piezoelectric crystal parameters caused by changes in temperature. A further advantage of increasing $C_{ds}$ is that it, together with $R_d$, provides low pass filtering of the output to reduce harmonic content and prevent oscillations at overtones.

Recall that the transconductance, $g_m$, for an FET is given by

$$g_m = \frac{2 I_{ds}}{V_p} \left(1 - \frac{V_{gsq}}{V_p}\right)$$

where $V_p$ and $I_{ds}$ are parameters for a particular FET and $V_{gsq}$ is the bias value of the gate-source voltage. The value of $g_m$ for a particular FET is greatest when $V_{gsq} = 0$. To achieve this condition, the Pierce oscillator circuit that we showed initially is often modified by setting $R_s = 0$ and letting $R_{b2} \rightarrow \infty$. With $R_s = 0$ note that $C_s$ is unnecessary.

If we redraw our initial circuit to reflect the addition of supplemental external capacitors for $C_{ds}$ and $C_{gs}$, as well as the removal of $R_s$, $R_{b2}$ and $C_s$, we have:
where $C_{dxx}$ and $C_{gxx}$ are the supplemental external capacitors for $C_{ds}$ and $C_{gs}$, respectively. This configuration, with external capacitors, is sometimes called a Colpitts crystal oscillator, rather than a Pierce crystal oscillator. In practice, a CMOS inverter is often substituted for the FET. Note that despite the changes, the signal equivalent circuit that we analyzed still applies, and hence all of our analysis still holds, provided only that we replace $C_{dxx}$ and $C_{gxx}$ with $C_{dxx} + C_{dxs}$ and $C_{gxx} + C_{gxs}$, respectively.
6.7 Recommended Questions

1. What is positive and negative feedback?
2. What are voltage amplifier and current amplifier? Give their equivalent circuit.
3. Show that for current series feedback amplifier input and output resistances are increased by a factor \((1+A\beta)\) with feedback.
4. Show that for voltage shunt feedback amplifier transresistance gain, \(R_1\) and \(R_o\) are decreased by a factor \((1+A\beta)\) with feedback.
5. Draw the block schematic of amplifier with negative feedback.
6. Draw the circuit diagram of a current series feedback circuit and derive expression for voltage gain and output resistance, and input resistance.
7. Draw the circuit diagram of voltage series feedback and derive expressions for input resistance and output resistance.
8. Draw the circuit diagram of a voltage shunt feedback using BJT and derive expression for voltage gain with feedback.
9. Draw the circuit for current shunt amplifier and justify the type of feedback. Derive the expression for \(A_v\), \(\beta\), \(R_i\) and \(R_o\) for the circuit.
10. Explain the concepts of feedback as applied to electronic amplifier circuits. What are the advantages and disadvantages of positive and negative feedback?
11. Derive an expression for input and output resistance of a voltage shunt feedback amplifier. (July-2007)
UNIT – 7: Linear Power Supplies, Switched mode Power Supplies

7.1 Regulated power supply

The regulated power supply converts the standard 220 volts, 50 or 60 Hz AC available at wall outlets into a constants DC voltage. It is one of the most common electronics circuits that we can find. The DC voltage produce by a power supply is used to power all the types of electronic circuits, such that television receiver, stereo system, CD players and laboratory equipment.

The regulated dual voltage DC power supply is to be used for the FM receiver. Two regulators, one positive and the other negative, provide the positive voltage required for the receiver circuits and the dual polarity voltage for the op-amp circuits.

The regulated power supply is to provide the necessary dc voltage and current, with low levels of ac ripple and with stability and regulation.

There are various methods of achieving a stable dc voltage from ac mains. The two methods are more commonly used. These are used;

(i) a linear voltage regulator and
(ii) A switching mode regulator.

Several types of both linear and switching regulators are available in integrated circuit (IC) form. By using the linear voltage regulator method, we must get the regulated dual dc power supply.

Fig: (1) Block Diagram of the Regulated Dual Voltage DC Power Supply

7.2 POWER SUPPLY FILTER

A power supply filter ideally eliminates the fluctuations in the output voltage of a half –wave rectifier and produces a constant-level dc voltage. The 60Hz pulsating dc output of a half-wave rectifier or the 120Hz pulsating output of a full-wave rectifier must be filtered to reduce the large voltage variations. Fig: (4.1) illustrates the filtering concepts showing a nearly smooth dc output voltage from the filter. The small amount of fluctuation in the filter output voltage is called ripple.
Capacitor Filter

A half wave rectifier with a capacitor filter is shown in **Fig: 4.2.** During the positive first quarter-circle of the input, the diode is forward bias and presents a low resistance path, allowing the capacitor to charge to within \(0.7V\) of the input peak. When the input begins to decrease below its peak, the capacitor retains its charge and the diode becomes reversed biased since the cathode is more positive than the anode. During the remaining part of the cycle, the capacitor can discharge only through the load resistor at a rate determined by the \(R_LC\) time constant.

(a) Initial charging of capacitor (diode is forward-biased) happens only once when power is turn on.

(b) Discharging through \(R_L\) after peak of positive alternation (diode is reverse biased)
Ripple Voltage

The capacitor quickly at the beginning of a cycle and slowing discharges after the positive peak. The variation in the output voltage due to charging and discharging is called the ripple voltage.

There are many type of circuit to regulate a certain dc voltage. Discrete circuits can be constructed using feed back transistors to get a voltage regulator. There also exists many IC types of voltage regulators. The well-known types of voltage regulator ICs are;

1. The 78XX series - for positive regulators
(2) The **79XX** series - for negative regulators

(3) The **LM 317** - for adjustable positive regulators

(4) The **LM 337** - for adjustable negative regulators

### 7.3 Fixed Positive Linear Voltage Regulators

The **78XX** series of IC regulators is representative of three terminal devices that provide a fixed positive output voltage. The three terminals are input, output and ground as indicated in the standard fixed voltage configuration in **Fig: (5.1.a)**.The last two digits in the part number designate the output voltage. For example, the ‘7805’ is a +5V regulator. Other available output voltages are given in **Table: 5.1**.

Capacitors although not always necessary are sometime used on the input and output as indicated in **Fig: (5.1.b)**. The output capacitor acts basically as a line filter to improve transient response. The input capacitor is use to prevent unwanted oscillations when the regulator is some distance from the power supply filter such that the line has a significant inductance.

The **78XX** can produce output current in excess of 1A when used with an adequate heat sink. The **78LXX** series can provide up to 100mA, the **78MXX** series can provide up to 500mA, and the **78TXX** series can provide in excess of 3A.

The input voltage must be at least 2V above the output voltage in order to maintain regulation. The circuits have internal thermal overload protection and short-circuit current-limiting features. Thermal overload occurs when the internal power dissipation becomes excessive and the temperature of the device exceeds a certain value.

![Fig: (5.1.a) Pin Layout](image1)

![Fig: (5.1.b) Standard configuration](image2)

<table>
<thead>
<tr>
<th>TYPE NUMBER</th>
<th>OUTPUT VOLTAGE</th>
</tr>
</thead>
<tbody>
<tr>
<td>7805</td>
<td>+5V</td>
</tr>
</tbody>
</table>
### Fixed Negative Linear Voltage Regulators

The 79XX series is typical of three-terminals IC regulators that provide a fixed negative output voltage. This series is the negative counterpart of the 78XX series and shares most of the same features and characteristics. **Fig: (5.2.a & b) and Table (5.2)** indicate the pin layout, the standard configuration and part numbers with corresponding output voltage that are available.

#### Table (5.2) 79XX series

<table>
<thead>
<tr>
<th>Type Number</th>
<th>Output Voltage</th>
</tr>
</thead>
<tbody>
<tr>
<td>7905</td>
<td>-5V</td>
</tr>
<tr>
<td>7905.2</td>
<td>-5.2V</td>
</tr>
<tr>
<td>7906</td>
<td>-6V</td>
</tr>
<tr>
<td>7908</td>
<td>-8V</td>
</tr>
<tr>
<td>7912</td>
<td>-12V</td>
</tr>
<tr>
<td>7915</td>
<td>-15V</td>
</tr>
<tr>
<td>7918</td>
<td>-18V</td>
</tr>
<tr>
<td>7924</td>
<td>-24V</td>
</tr>
</tbody>
</table>

---

**Table (5.1) 78XX series**

<table>
<thead>
<tr>
<th>Type Number</th>
<th>Output Voltage</th>
</tr>
</thead>
<tbody>
<tr>
<td>7806</td>
<td>+6V</td>
</tr>
<tr>
<td>7808</td>
<td>+8V</td>
</tr>
<tr>
<td>7809</td>
<td>+9V</td>
</tr>
<tr>
<td>7812</td>
<td>+12V</td>
</tr>
<tr>
<td>7815</td>
<td>+15V</td>
</tr>
<tr>
<td>7818</td>
<td>+18V</td>
</tr>
<tr>
<td>7824</td>
<td>+24V</td>
</tr>
</tbody>
</table>
Adjustable Positive Linear Voltage Regulators

The **LM317** is an excellent example of the three-terminal positive regulator with an adjustable output voltage. Notice that there is an input, an output and an adjustable terminal. The external fixed resistor \( R_1 \) and the external variable resistor \( R_2 \) provide the output voltage adjustment. \( V_{\text{out}} \) can be varied from 1.2V to 37V depending on the resistor values. The **LM317** can provide over 1.5A of output current to a load.

The **LM317** is operated as a “floating” regulator because the adjustment terminal is not connected to ground, but floats to whatever voltage is across \( R_2 \). This allows the output voltage to be much higher than that of a fixed-voltage regulator.

A constant 1.25V reference voltage (\( V_{\text{REF}} \)), is maintained by the regulator between the output terminal and the adjustment terminal. This constant reference voltage produces a constant current (\( I_{\text{REF}} \)) through \( R_1 \) regardless of the value of \( R_2 \). \( I_{\text{REF}} \) also flows through \( R_2 \),

\[
I_{\text{REF}} = \frac{V_{\text{REF}}}{R_1} = \frac{1.25V}{R_1}
\]

\[
V_{\text{OUT}} = V_{R1} + V_{R2} = I_{\text{REF}}R_1 + I_{\text{REF}}R_2
\]

\[
= I_{\text{REF}}(R_1 + R_2)
\]

\[
= \frac{V_{\text{REF}}}{R_1}(R_1 + R_2)
\]
Adjustable Negative Linear Voltage Regulators

\[ V_{\text{OUT}} = V_{\text{REF}} \left(1 + \frac{R_2}{R}\right) \]

The **LM 337** is the negative output counterpart of the **LM 317** and is a good example of the type of IC regulator. Like the **LM 317**, the **LM 337** requires two external resistors for output voltage adjustments as shown in **Fig: (5.4.b)**. The output voltage can be adjusted from \(-1.2 \text{V}\) to \(-37 \text{V}\), depending on the external resistors values. The electrical characteristics of the **LM 317** and **LM 337** are shown in **Table (5.4)**.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Conditions</th>
<th>LM317/LM337</th>
<th>Units</th>
</tr>
</thead>
<tbody>
<tr>
<td>Line Regulation</td>
<td>(T_A = 25^\circ \text{C}, 3 \text{V} \leq V_{\text{in}} - V_{\text{out}} \leq 40 \text{V})</td>
<td>0.04</td>
<td>%/V</td>
</tr>
<tr>
<td>Load Regulation</td>
<td>(T_A = 25^\circ \text{C}, 10 \text{mA} \leq I_{\text{out}} \leq I_{\text{max}})</td>
<td>(V_{\text{out}} \leq 5 \text{V})</td>
<td>25.00</td>
</tr>
<tr>
<td></td>
<td></td>
<td>(V_{\text{out}} \geq 5 \text{V})</td>
<td>0.4</td>
</tr>
<tr>
<td>Thermal Regulation</td>
<td>(T_A = 25^\circ \text{C}, 20 \text{ms Pulse})</td>
<td>0.07</td>
<td>%/W</td>
</tr>
<tr>
<td>Adj: Pin Current</td>
<td></td>
<td>100.0</td>
<td>(\mu\text{A})</td>
</tr>
<tr>
<td>Reference Voltage</td>
<td></td>
<td>1.25</td>
<td>V</td>
</tr>
<tr>
<td>Temperature Stability</td>
<td>(T_{\text{min}} \leq T_J \leq T_{\text{max}})</td>
<td>1.00</td>
<td>%</td>
</tr>
<tr>
<td>Ripple Rejection Ratio</td>
<td>(V_{\text{out}} = 10 \text{V}, F = 120\text{Hz}, C_{\text{adj}} = 10 \mu\text{F})</td>
<td>80.00</td>
<td>\text{Db}</td>
</tr>
<tr>
<td>Current Limit (Max)</td>
<td>((V_{\text{IN}} - V_{\text{OUT}} \leq 15 \text{V}))</td>
<td>1.00</td>
<td>\text{A}</td>
</tr>
<tr>
<td>Current Limit (Min)</td>
<td>((V_{\text{IN}} - V_{\text{OUT}} = 40 \text{V}))</td>
<td>0.40</td>
<td>\text{A}</td>
</tr>
</tbody>
</table>

Table (5.4) Electrical Characteristics of LM 317 & LM 337
7.4 The Regulated Dual Voltage DC Power Supply

By combining the step down transformer, rectifier, filters and voltage regulators together, we get a regulated dual voltage dc power supply circuit as shown in Fig: (6.1).

Working Principle

This is a simple circuit, which gives regulated ±1.2V to ±15V supply. ICs LM 317T and LM 337T are used here as positive and negative regulators respectively.

The LM 317T regulator has internal feedback regulating and current passing elements. It incorporates various protection circuits such as current limit (which limits package power dissipation to 15 watts for the TO-220 package) and thermal shutdown. Thus these two ICs form an independently adjustable bipolar power supply.

Capacitors although not always necessary are sometimes used on the input and output as indicated in Fig: (6.1). The output capacitors $C_7$ and $C_8$ acts basically as line filter to improve transient response. The input capacitors $C_3$ and $C_4$ are used to prevent unwanted oscillations when the regulator is some distance from the power supply filter such that the line has a significant inductance. $D_5$ and $D_6$ prevent short-circuit for input and output terminals.

The TO-220 packages will easily furnish one ampere each if the heat sinks are properly mounted. Variable resistors $V_{R1}$ and $V_{R2}$ are adjusted for each regulator to give a regulated output approximately between ±1.2V to ±15V. Capacitors $C_5$ and $C_6$ are used to improve AC ripple voltage rejection. However, if a short-circuit occurs across the regulator outputs, $C_5$ and $C_6$ will adjust the current in the terminals. The output can be calculated by the formula:

$$V_0 = 1.25 \, V \left(1 + \frac{V_{R1}}{R_1}\right)$$

$$I_{L(max)} = \frac{P_d}{V_{different}}$$

Circuit Diagram
Fig: (6.1) Circuit diagram of the regulated dual voltage DC power supply
7.5 Recommended Questions

1. Design an op-amp Schmitt trigger circuit to meet the following specifications.
2. \( V_{UTP} = 4V, V_{LTP} = -2V, V= \pm 12v. \) (July-2006)
3. Explain how an op-amp can be used as comparator.
4. Draw and explain inverting and non-inverting comparator circuits.
5. Explain op-amp is a Schmitt trigger.
6. Draw and explain the working of Schmitt trigger with different threshold levels LTP and UTP.
7. What is integrator? Derive the expression for an output of op-amp integrator.
8. Draw and explain the basic timing circuit of IC 555. Draw the necessary waveforms.
9. Draw and explain the function block diagram of IC 555.
10. Derive the expression for the pulse width of a monostable multivibrator using IC 555.
11. Derive the expression for the frequency of the output of an astable multivibrator.
12. List the features of IC 555.
13. Draw the circuit diagram of an integrator and explain its operation with a typical input pulse. (July-2008)
14. Draw and explain the circuit diagram of voltage controlled oscillator using the 555 timer. (July-2008)
15. Obtain an expression for the closed loop gain of non-inverting amplifier. (July-2007)
UNIT – 8: Operational Amplifier

8.1 Operational Amplifier (Op-Amp)

An operational amplifier ("op-amp") is a DC-coupled high-gain electronic voltage amplifier with a differential input and, usually, a single-ended output. An op-amp produces an output voltage that is typically hundreds of thousands times larger than the voltage difference between its input terminals.

Operational amplifiers are important building blocks for a wide range of electronic circuits. They had their origins in analog computers where they were used in many linear, non-linear and frequency-dependent circuits.

The circuit symbol for an op-amp is shown below:

![Op Amp Block Diagram](www.vtucs.com)

The input stage is a differential amplifier. The differential amplifier used as an input stage provides differential inputs and a frequency response down to DC. Special techniques are used to provide the high input impedance necessary for the operational amplifier.

The second stage is a high-gain voltage amplifier. This stage may be made from several transistors to provide high gain. A typical operational amplifier could have a voltage gain of 200,000. Most of this gain comes from the voltage amplifier stage.

The final stage of the OP AMP is an output amplifier. The output amplifier provides low output impedance. The actual circuit used could be an emitter follower. The output stage should allow the operational amplifier to deliver several milliamperes to a load.
Notice that the operational amplifier has a positive power supply (+$V_{CC}$) and a negative power supply (-$V_{EE}$). This arrangement enables the operational amplifier to produce either a positive or a negative output.

The two input terminals are labeled "inverting input" (-) and "non-inverting input" (+). The operational amplifier can be used with three different input conditions (modes). With differential inputs (first mode), both input terminals are used and two input signals which are 180 degrees out of phase with each other are used. This produces an output signal that is in phase with the signal on the non-inverting input. If the non-inverting input is grounded and a signal is applied to the inverting input (second mode), the output signal will be 180 degrees out of phase with the input signal (and one-half the amplitude of the first mode output). If the inverting input is grounded and a signal is applied to the non-inverting input (third mode), the output signal will be in phase with the input signal (and one-half the amplitude of the first mode output)

**Differential Amplifier Input Stage**

The differential amplifier configuration is also called as long-tail pair as the two transistors share a common-emitter resistor. The current through this resistor is called the tail current. The base terminal of transistor $TR_1$ is the non-inverting input and base terminal of transistor $TR_2$ is the inverting input. The output is in-phase with the signal applied at non-inverting input and out-of-phase with the signal applied at inverting input. If two different signals are applied to inverting and non-inverting inputs, the output is given by $A_d x (V_1-V_2)$
8.2 Ideal Op-Amp

The ideal opamp model was derived to simplify circuit calculations. The ideal opamp model makes three assumptions. These are as follows:

1. Input impedance, $Z_i = \infty$
2. Output impedance, $Z_o = 0$
3. Open-loop gain, $A_d = \infty$

From the above three assumptions, other assumptions can be derived. These include the following:

1. Since $Z_i = \infty$, $I_i = I_M = 0$.
2. Since $Z_o = 0$, $V_o = A_d \times V_d$.
3. Common mode gain = 0
4. Bandwidth = $\infty$
5. Slew Rate = $\infty$
6. Offset Drift = 0

Performance parameters

1. **Bandwidth**: Bandwidth of an opamp tells us about the range of frequencies it can amplify for a given amplifier gain.
2. **Slew rate**: It is the rate of change output response to the rate of change in input. It is one of the most important parameters of opamp. It gives us an idea as to how well the opamp output follows a rapidly changing waveform at the input. It is defined as the rate of change of output voltage with time. Slew rate limits the large signal bandwidth. Peak-to-peak output voltage swing for a sinusoidal signal ($V_{p-p}$), slew rate and bandwidth are inter-related by the following equation:

   \[
   \text{Bandwidth} = \frac{\text{Slewrate}}{\pi \times V_{p-p}}
   \]

3. **Open-loop gain**: Open-loop gain is the ratio of single-ended output to the differential input.
4. **Common Mode Rejection Ratio**: Common Mode Rejection Ratio (CMRR) is a measure of the ability of the opamp to suppress common mode signals. It is the ratio of the desired differential gain ($A_d$) to the undesired common mode gain ($A_c$).

   \[
   \text{CMRR} = 20 \log \left( \frac{A_d}{A_c} \right) \text{dB}
   \]
5. **Power Supply Rejection Ratio**: Power Supply Rejection Ratio (PSRR) is defined as the ratio of change in the power supply voltage to corresponding change in the output voltage. PSRR should be zero for an Ideal opamp.

6. **Input Impedance**: Input Impedance ($Z_i$) is the impedance looking into the input terminals of the opamp and is mostly expressed in terms of resistance. Input impedance is the ratio of input voltage to input current and is assumed to be infinite to prevent any current flowing from the source supply into the amplifiers input circuitry ($I_{in} = 0$).

7. **Output Impedance**: Output Impedance ($Z_o$) is defined as the impedance between the output terminal of the opamp and ground. The output impedance of the ideal operational amplifier is assumed to be zero acting as a perfect internal voltage source with no internal resistance so that it can supply as much current as necessary to the load. This internal resistance is effectively in series with the load thereby reducing the output voltage available to the load.

8. **Settling Time**: Settling Time is a parameter specified in the case of high speed opamps of the opamps with a high value of gain-bandwidth product.

9. **Offset Voltage ($V_{io}$)**: The amplifiers output will be zero when the voltage difference between the inverting and the non-inverting inputs is zero, the same or when both inputs are grounded. Real op-amps have some amount of output offset voltage.
8.3 Applications of Opamp

Peak Detector Circuit

Peak detector circuit produces a voltage at the output equal to peak amplitude (positive or negative) of the input signal. It is a clipper circuit with a parallel resistor-capacitor connected at its output.

The clipper here reproduces the positive half cycles. During this period, the diode D1 is forward-biased. The capacitor rapidly charges to the positive peak from the output of the opamp. The capacitor can now discharge only through the resistor (R) connected across it. The value of the resistor is much larger than the forward-biased diode’s ON resistance. The buffer circuit connected ahead of the capacitor prevents any discharge of the capacitor due to loading effects of the following circuit. The circuit can be made to respond to the negative peaks by reversing the polarity of the diode.

Input / Output Waveform of Peak Detector Circuit
Absolute Value Circuit

It is the configuration of opamp that produces at its output a voltage equal to the absolute value of the input voltage. The circuit shown above is the dual half wave rectifier circuit.

When the applied input is of positive polarity (+V), diode D1 is forward biased and diode D2 is reverse biased. The output (V_o) in this case is equal to +V.

When the applied input is of negative polarity (-V), diode D1 is reverse biased and diode D2 is forward biased.

By applying Kirchoff’s Current Law (KCL) at the inverting terminal of the opamp, we can determine voltage (V_x) to be equal to \( \frac{2}{3} V \). Also, \( V_x \) is related to \( V_o \) by \( V_x = \left( \frac{2}{3} \right) V_o \). This implies that \( V_o = V \).

Thus the output always equals the absolute value of the input signal.

Comparator

Non-inverting comparator with positive reference and negative reference

A comparator circuit is a two input, one-output building block that produces a high or low output depending upon the relative magnitudes of the two inputs. An opamp can be very conveniently used as a comparator when used without negative feedback. Because of very large value of open-loop voltage gain,
it produces either positively saturated or negatively saturated output voltage depending upon whether the amplitude of the voltage applied at the non-inverting terminal is more or less positive than the voltage applied at the inverting input terminal.

In general, reference voltage voltage may be a positive or a negative voltage. In the above figure, non-inverting comparator with a positive reference voltage, $V_{REF}$ is given by

$$+V_{CC} \times \left[ \frac{R_2}{R_1 + R_2} \right]$$

In the above figure, in the case of non-inverting comparator with a negative reference voltage, $V_{REF}$ is given by

$$-V_{CC} \times \left[ \frac{R_2}{R_1 + R_2} \right]$$

**Zero Crossing Detector**
Non-inverting zero-crossing detector  

Inverting zero-crossing detector

One of the inputs of the comparator is generally applied a reference voltage and the other input is fed with the input voltage that needs to be compared with the reference voltage. In special case where the reference voltage is zero, the circuit is referred to as zero-crossing detector. The above figure shows inverting and non-inverting zero crossing detector circuits with their transfer characteristics and their input / output waveforms.

In non-inverting zero-crossing detector, input more positive than zero leads to a positively saturated output voltage. Diodes D1 and D2 connected at the input are to protect the sensitive input circuits inside the opamp from excessively large input voltages. In inverting zero-crossing detector, input voltage slightly more positive than zero produces a negatively saturated output voltage. One common application of zero-crossing detector is to convert sine wave signal to a square wave signal.

Comparator with Hysteresis
The above circuit diagram shows the inverting and non-inverting comparator with hysteresis. The circuit functions as follows.

Let us assume that the output is in positive saturation (+V_{SAT}). Voltage at non-inverting input in this case is

\[ +V_{SAT} \times \frac{R_2}{R_1 + R_2} \]

Due to this small positive voltage at the non-inverting input, the output is reinforced to stay in positive saturation. Now, the input signal needs to be more positive than this voltage for the output to go to negative saturation. Once the output goes to negative saturation (-V_{SAT}), voltage fed back to non-inverting input becomes

\[ -V_{SAT} \times \frac{R_2}{R_1 + R_2} \]

A negative voltage at the non-inverting input reinforces the output to stay in negative saturation. In this manner, the circuit offers a hysteresis of

\[ 2V_{SAT} \times \frac{R_2}{R_1 + R_2} \]

Non-inverting comparator with hysteresis can be built by applying the input signal to the non-inverting input as shown in the figure above. Operation is similar to that of inverting comparator. Upper and lower trip points and hysteresis is given by

\[ UTP = +V_{SAT} \times \frac{R_2}{R_2} \]
In the case of a conventional comparator, the output changes state when the input voltage goes above or below the preset reference voltage. In a window comparator, there are two reference voltages called the lower and the upper trip points.

When the input voltage is less than the voltage reference corresponding to the lower trip point (LTP), output of opamp $A_1$ is at $+V_{SAT}$ and the opamp $A_2$ is at $-V_{SAT}$. Diodes $D_1$ and $D_2$ are respectively forward and reverse biased. Consequently, output across $R_L$ is at $+V_{SAT}$.

When the input voltage is greater than the reference voltage corresponding to the upper trip point (UTP), the output of opamp $A_1$ is $-V_{SAT}$ and that of opamp $A_2$ is at $+V_{SAT}$. Diodes $D_1$ and $D_2$ are respectively reverse and forward biased. Consequently, output across $R_L$ is at $+V_{SAT}$.

When the input voltage is greater than LTP voltage and lower than UTP voltage, the output of both opamps is at $-V_{SAT}$ with the result that both diodes $D_1$ and $D_2$ are reverse biased and the output across $R_L$ is zero.

Active Filters

Opamp circuits are used to build low-pass, high-pass, band-pass and band-reject active filters. Also filters are classified depending on their order like first-order and second-order. Order of an active filter is determined by number of RC sections used in the filter.

First-Order Filters
The simplest low-pass and high-pass active filters are constructed by connecting lag and lead type of RC sections, respectively, to the non-inverting input of the opamp wired as a voltage follower. The first order low-pass and high-pass filters are shown in the figure below.

![First-order low-pass active filter](image1)

![First-order high-pass active filter](image2)

In the case of low-pass filter, at low frequencies, reactance offered by the capacitor is much larger than the resistance value and therefore applied input signal appears at the output mostly unattenuated.

At high frequencies, the capacitive reactance becomes much smaller than the resistance value thus forcing the output to be near zero.

**The operation of high-pass filter can be explained as follows.**

At high frequencies, reactance offered by the capacitor is much larger than the resistance value and therefore applied input signal appears at the output mostly unattenuated.

At low frequencies, the capacitive reactance becomes much smaller than the resistance value thus forcing the output to be near zero.

![Low-pass filter with gain](image3)

![High-pass filter with gain](image4)

The cut-off frequency and voltage gain in both cases is given by

\[
 f_c = \frac{1}{2\pi RC}
\]
Inverting Low-pass filter with gain

Inverting High-pass filter with gain

The cut-off frequency and voltage gain in case of Inverting filters is given by

\[ f_c = \frac{1}{2\pi R_c C} \]

\[ A_v = \frac{R_v}{R_1} \]

**Second Order Filters**

Butterworth filter is the commonly used second order filter, it is also called as flat filter, offers a relatively flat pass and stop band response. The generalized form of second-order Butterworth filter is shown in the figure below.

1. If \( Z_1 = Z_2 = R \) and \( Z_3 = Z_4 = C \), we get a second-order low-pass filter
2. If \( Z_1 = Z_2 = C \) and \( Z_3 = Z_4 = R \), we get a second-order high-pass filter

Generalized form of second-order Butterworth filter

The cut-off frequency and pass band gain \( (A_v) \) is given by
Band-pass filters can be formed by cascading the high-pass and the low-pass filter sections in series. These filters are simple to design and offer large bandwidth.

At very low frequencies, $C_1$ and $C_2$ offer very high reactance. As a result, the input signal is prevented from reaching the output.

At very high frequencies, the output is shorted to the inverting input, which converts the circuit to an inverting amplifier with zero gain. Again, there is no output.

\[
f_c = \frac{1}{2\pi RC} \quad A_v = 1 + \frac{R_2}{R_1}
\]

Narrow band-pass filter

At some intermediate band of frequencies, the gain provided by the circuit offsets the loss due to potential divider $R_1$-$R_3$. The resonant frequency is given by

\[
f_R = \frac{2Q}{2\pi R_2 C}
\]

Where $Q$ is the quality factor, for $C_1 = C_2 = C$, the quality factor and voltage gain is given by

\[
Q = \left( \frac{R_2 R_2}{2R_5} \right)^{1/2} \quad A_v = \frac{Q}{2\pi R_2 f_R C}
\]

Band-reject filters can be implemented by summing together the outputs of the low-pass and high-pass filters. These filters are simple to design and have a broad reject frequency range.
Second-order Band-Reject filter

It uses a twin-T network that is connected in series with the non-inverting input of the opamp. Very low frequency signals find their way to the output via the low-pass filter formed by \( R_1 - R_2 - C_3 \). Very high frequency signals reach the output through the high-pass filter formed by \( C_1 - C_2 - R_3 \). Intermediate band of frequencies pass through both the filters, net signal reaching the non-inverting input and hence the output is zero. Component values are chosen by

\[
R_1 = R_2 = R, R_2 = \frac{R}{2}, \quad C_1 = C_2 = C, C_3 = 2C, \quad 0 \leq R_4 \leq (R_1 + R_2)
\]

\[
f_{R} = \frac{1}{2\pi RC}
\]

**Phase Shifters**

Phase Shifters can be used to shift the phase of the input signal over a wide range by varying \( R_p \) with 0\(^\circ\) and -180\(^\circ\) being the extremes. The output lags the input by an angle \( \theta \)

**Lag-type Phase Shifter**

\[
\theta = -2 \tan^{-1}(\omega R_p C_p)
\]

**Lead-type Phase Shifter**

\[
\theta = 2 \tan^{-1}(\omega R_p C_p)
\]
For \( R_p \ll \frac{1}{\omega C_p} \), \( \theta = 0^\circ \)  

For \( R_p \gg \frac{1}{\omega C_p} \), \( \theta = -180^\circ \)  

For \( R_p = \frac{1}{\omega C_p} \), \( \theta = -90^\circ \)  

For \( R_p = \frac{1}{\omega C_p} \), \( \theta = 90^\circ \)

**8.4 Instrumentation Amplifier**

Instrumentation amplifier is a differential amplifier that has been optimized for DC performance. It is having a high differential gain, high CMRR, high input impedance, low input offsets and low temperature drifts.

![Instrumentation amplifier diagram](image)

When common-mode signal is applied to \( V_i \), i.e., same positive voltage applied to both the non-inverting inputs, voltages appearing at the output of opamps A1 and A2 and also at R1-R2 and R3-R4 junctions are equal. A1 and A2 acts like voltage followers. In other words common-mode gain ACM of the preamplifier stage is unity.

On the other hand, when a differential signal is applied to the input, signals appearing at two R1-R2 and R3-R4 junctions are equal and opposite creating a virtual ground at point A. the differential gain of this stage is therefore \( 1 + \frac{R_2}{R_1} \)

The differential gain and common-mode gain of the amplifier is given by

\[
A_v = 1 + \frac{2R_2}{R_G}
\]

\[
A_{CM} = \pm \frac{2\Delta R}{R}
\]

Non-Linear Amplifier
Non-Linear amplifier is the circuit, where the gain value is a non-linear function of the amplitude of the signal applied at the input. i.e., the gain may be very large for weak input signals and very small for large input signals, which implies that for a very large change in the amplitude of input signal, resultant change in amplitude of output signal is very small.

For small values of input signal, diodes act as open circuit and the gain is high due to minimum feedback. When the amplitude of input signal is large, diodes offer very small resistance and thus gain is low.
8.5 Relaxation Oscillator

Relaxation oscillator is an oscillator circuit that produces a non-sinusoidal output whose time period is dependent on the charging time of a capacitor connected as a part of the oscillator circuit.

Let us assume that the output is initially in positive saturation. As a result, voltage at non-inverting input of opamp is \( +V_{SAT} \times \frac{R_1}{(R_1 + R_2)} \). This forces the output to stay in positive saturation as the capacitor C is initially in fully discharged state. Capacitor C starts charging towards +VSAT through R.

The moment the capacitor voltage exceeds the voltage appearing at the non-inverting input, the output switches to −VSAT. The voltage appearing at non-inverting input also changes to \( -V_{SAT} \times \frac{R_1}{(R_1 + R_2)} \).

The capacitor starts discharging after reaching zero, it begins to discharge towards −VSAT.

Again, as soon as it becomes more negative than the negative threshold appearing at non-inverting input of the opamp, the output switches back to +VSAT. The cycle repeats thereafter. The output is a rectangular wave.

The expression for time period of output waveform can be derived from the exponential charging and discharging process and is given by

\[
T = 2RC \ln \left( \frac{1 + \beta}{1 - \beta} \right)
\]

Where \( \beta = \frac{R_1}{(R_1 + R_2)} \)
8.6 Current-to-Voltage Converter

Current-to-Voltage converter is nothing but a transimpedance amplifier. An ideal transimpedance amplifier makes a perfect current-to-voltage converter as it has zero input impedance and zero output impedance.

\[ V_o = I_i \times R \times \left( \frac{A_{OL}}{1 + A_{OL}} \right) \]

For \( A_{OL} >> 1 \), \( V_o = I_i \times R \)

Where \( Z_{in} = \frac{R}{1 + A_{OL}} \) and \( Z_o = \frac{R_o}{1 + A_{OL}} \)

8.7 Voltage-to-Current Converter

Voltage-to-Current converter is a transconductance amplifier. An ideal transconductance amplifier makes a perfect voltage-controlled current source or a voltage-to-current converter.

\[ I_o = \frac{V_i}{R_2 + \left( \frac{R_1 + R_2}{A_{OL}} \right)} \]

For \( AOL >> 1 \), \( I_o = \frac{V_i}{R_2} \)
8.8 Exercise Problems

Problem 1

From the circuit shown below determine the quiescent DC voltage at the collector terminal of each transistor assuming $V_{BE}$ of the non-inverting transistors to be negligible. What will be the quiescent DC value if $V_{BE}$ is taken at 0.7V?

Solution:-

1. Assuming $V_{BE}$ to be negligible, the tail current

$$I_T = \frac{-V_{EE}}{R_E} = \frac{12}{10 \times 10^3} = 1.2mA$$

2. Emitter current of each transistor

$$I_E = \frac{1.2 \times 10^{-3}}{2} = 0.6mA$$

3. Therefore, $I_C = I_E = 0.6mA$

4. Quiescent DC voltage at the collector of each transistor
5. If $V_{BE} = 0.7V$, Tail current

$$I_T = \frac{-V_{EE} - V_{BE}}{R_E} = \frac{12 - 0.7}{10 \times 10^3} = 1.13mA$$

6. Thus Emitter and Collector currents of each transistor is

$$I_E = I_C = \frac{I_T}{2} = \frac{1.13 \times 10^{-3}}{2} = 0.565 mA$$

7. Quiescent DC voltage at collector of each transistor is

$$V_{CEQ} = V_{CC} - I_C R_C = 12 - 0.565 \times 10^{-3} \times 10 \times 10^3 = 6.35V$$

Problem 2

From the circuit shown below determine the cut-off frequency and the gain value at four times the cut-off frequency.

Solution:

1. Cut-off frequency,

$$f_c = \frac{1}{2\pi RC} = \frac{\frac{1}{2\pi} \times 10^{10} \times 1}{1000 \times 10^{-12}} = \frac{10^5}{2\pi} Hz = 15.915 kHz$$

2. Gain, $A_v = 1 + \frac{R_2}{R_1} = 1 + \frac{100 \times 10^3}{(10 \times 10^3)} = 11 = 20.827 dB$

3. Gain at cut-off point = 20.827 - 3 = 17.827dB

4. Gain at frequency four times the cut-off frequency will be 12 dB below the value of mid-band gain.

5. Therefore, gain at four times the cut-off frequency = 20.827 - 12 = 8.827 dB
Problem 3

A second-order low-pass filter built around a single opamp is shown in the figure below. Calculate the values of R1, R2, C1, C2 and R3 if the filter had a cut-off frequency of 10 kHz, Q-factor of 0.707 and input impedance not less than 10kΩ.

Solution:-

1. Q-factor is given by \( Q = \left( \frac{1}{2} \right) \times \sqrt{\frac{C_1}{C_2}} \)
2. For Q = 0.707, \( C_1 = 2C_2 \).
3. For input impedance of 10kΩ, \( R_1 = R_2 = 10kΩ \), also \( R_3 = R_1 + R_2 \). \( \therefore R_3 = 20kΩ \)
4. \( f_c = \frac{1}{2\pi R \sqrt{C_1C_2}} \), \( 10 \times 10^3 = \frac{1}{2\pi \times 10 \times 10^3 \times C_2 \times \sqrt{2}} \). \( \therefore C_2 = 0.0011\mu F \)
5. \( C_1 = 2C_2 = 0.0022\mu F \)

8.9 Recommended Questions

1. Explain with the block diagrams, the basic types of voltage regulator circuits.
2. Compare series regulator and shunt regulator.
3. Draw neat diagram of a series regulator with fold back protection and explain its operation.
4. Explain the working principle of basic switching regulator.
5. Draw the circuit and explain the operation of step up switching regulator.
6. Draw and explain block diagram of 3 terminal IC regulator.
7. Explain the working of SAR ADC. (July 2007)
8. Explain the working R-2R ladder DAC. (July 2007)
9. Explain the applications of astable multivibrator as: i) Square wave generator ii) To achieve variable duty cycle control. (July-2007)
10. Define the terms ‘Load Regulation’ and ‘Line Regulation’ with respect to a power supply. (July-2008)
11. Describe with circuit diagram, the operation of a shunt regulation power supply. (July-2008)
12. What are the advantages and disadvantages of shunt regulator? What are the advantages of a series regulator? (July-2008)

13. Explain the working of D/A converter [Binary weighted resistors] with neat sketch.

14. Draw simple sketches of IC linear regulators and IC switching regulators and compare their characteristics. (Jan-2009)